# $X Y Z$ at Belle 

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#### Abstract

Recent results of exotic states with heavy quarks, denoted as $X Y Z$ states, are presented. The results include searches for the $Y(4260)$ in $B$ meson decays, a spin partner of $Y(4630)$, and the strange Pentaquark $P_{s}^{+}$. An addition, the measurements of the absolute branching fraction of $X(3872)$ and search for double $Z_{c}$ production is dicussed.


## 1 Introduction

In recent years, numerous narrow states with heavy quarks $(c, b)$ have been observed, which do not fit the expected properties for charmonium or bottomonium states from the Cornell potential [1]. For detailed information about the states, their discovery and their properties, we refer to the recent reviews [2]. Throughout this paper below, charge conjugated states and decay modes are implied.

## $2 Y(4260)$ in $B$ meson decays

The vector state $Y(4260)$ has been observed in the ISR (initial state radiation) process $e^{+} e^{-} \rightarrow \gamma Y(4260) \rightarrow J / \psi \pi^{+} \pi^{-}$by BaBar [3], Belle [4], CLEO-c [5] and in direct production in $e^{+} e^{-} \rightarrow Y(4260)$ at BESIII [6]. Up to date, in total four states $Y(4008), Y(4260), Y(4350)$ and $Y(4660)$ have been observed by several experiments, all of them carrying the quantum numbers $J^{P C}=1^{--}$. In addition, there are the well known conventional vector charmonium states $\psi(4040), \psi(4140), \psi(4400)$ [7] in the same mass regime, leading to a overpopulation of states and thus pointing to the exotic nature of the $Y$ states.

In literature, the $Y(4260)$ has been discussed as a $D \bar{D}_{1}(2460)$ molecule and as a $P$-wave tetraquark [2]. In particular, the decay to $D^{(*)} \bar{D}^{(*)}$ has not been seen yet [7], although the phasespace is large. This observation can be explained in the context of a $\left[J / \psi f_{0}(980)\right]$ hadrocharmonium state (with dominant decays to "core" and "cloud" components) or of a [ $(c \bar{c})_{8} g$ ] hybrid (with open charm decays being blocked by the gluon fluxtube) [2].

Searches for the $Y(4260)$ in $B$ meson decays have started more than a decade ago by BABAR in the decay $B^{+} \rightarrow K^{+} Y(4260)$ with $Y(4260) \rightarrow J / \psi \pi^{+} \pi^{-}$. In a data set of $211 \mathrm{fb}^{-1}$, $128 \pm 42$ signal events were observed [8], corresponding to a significance of $3.1 \sigma$, and being translated to a range of a branching fraction of $1.2<\mathcal{B}<2.9 \times 10^{-5}$.

Recently, the D0 experiment reported the surprising observation with $167 \pm 41$ signal events ( $4.3 \sigma$ significance) in a data set of $10.4 \mathrm{fb}^{-1}$ [9]. However, different from $B A B A R$,

[^0]$b$-flavoured hadrons ( $B, B_{s}, B *, \Lambda_{b}, \ldots$ ) were selected inclusively instead of $B$ mesons, by requiring a $\left[J / \psi \pi^{+} \pi^{-}\right]$system with a secondary vertex.

A search in Belle was performed [10], using the complete Belle $\Upsilon(4 S)$ data set of $711 \mathrm{fb}^{-1}$ and corresponding to a factor 3.3 more data than in case of $B A B A R$. Scaling the signal and background yields would give $S / \sqrt{S+B}=5.7$, which would be a clear observation. In addition, neutral B decays with $B^{0} \rightarrow K_{s}^{0} Y(4260)$ were added to the analysis. Fig. 1 shows the $J / \psi \pi^{+} \pi^{-}$invariant mass distributions for the three different mass regions around the $\psi^{\prime}$, the $X(3872)$ and the $Y(4260)$. The first two served as control signals. For the $Y(4260)$, no signal was observed, with the measured branching fraction reaching the $10^{-6}$ level. The derived upper limit of $1.4 \times 10^{-5}$ for the charged $B$ meson decays is consistent with the BABAR measurement.

## 3 Search for a spin partner of the $Y(4630)$

The $Y(4630)$ was discovered in ISR as well and thus also represents a vector state with $J^{P C}=1^{--}$. Along with the $Y(4660)$ mentioned above in Sec. 2, it represents the highest charmed XYZ state with charm observed so far. If being a pure $c \bar{c}$ state, the radius ${ }^{1}$ of can be calculated in the Cornell potential $r \simeq 2.2 \mathrm{fm}$, far beyond string breaking limit outside the confinement region. In addition, it is the only XYZ state observed so far decaying into baryons. It was recognized early, that the $Y(4660)$, decaying to $J / \psi \pi^{+} \pi^{-}$and already mentioned above, and the $Y(4630)$, decaying to $\Lambda_{c}^{+} \Lambda_{c}^{-}$, may represent same state [11]. In the hadrocharmonium model, the $\mathrm{X}(3872)$ can be interpreted as a $\left[J / \psi \rho^{0}\right]$ state, the $Y(4260)$ as a $\left[J / \psi f_{0}(980)\right]$ state and the $Y(4630)$ as a $\left[\psi^{\prime} f_{0}(980)\right]$ state. The latter one would imply the existence of a $\left[\eta_{c}(2 S)\right.$ $f_{0}(980)$ ] partner state, called the $Y_{\eta}$, with a mass of $4613 \pm 4 \mathrm{MeV}$ and a width $\simeq 30 \mathrm{MeV}$, whereas the width is dominated by the $\eta_{c}(2 S)$. With $L=0$, the quantum numbers would be $J^{P C}=0^{-+}$, and thus it can not be produced in ISR. Therefore a search in $B$ meson decays was performed with the complete Belle $\Upsilon(4 \mathrm{~S})$ data set of $711 \mathrm{fb}^{-1}$ [13]. No signal was observed yet, leading to an upper limit of $\mathcal{B}\left(B^{+} \rightarrow K^{+} Y_{\eta}\right) \times \mathcal{B}\left(Y_{\eta} \rightarrow \Lambda_{c}^{+} \Lambda_{c}^{-}\right)<2.0 \times 10^{-4}$.

## 4 Absolute branching fraction of the $X(3872)$

The $X$ (3872) has been observed in five different decays [7]. However, all branching fractions are only established as lower limits [7], because the total branching fraction $\mathcal{B}_{\text {total }}$, being required as normalisation, is not known. There is a way to measure $\mathcal{B}_{\text {total }}$ at $B$ factories, which operate with $e^{+} e^{-}$collisions at the $\Upsilon(4 \mathrm{~S})$ resonance at $\sqrt{s}=10.58 \mathrm{GeV}$. The particular situation kinematic situation at the $\Upsilon(4 \mathrm{~S})$ is, that the mass of the $\Upsilon(4 \mathrm{~S})$ ) coincides very closely with the sum of the masses of the $B$ meson and the $\bar{B}$ meson, the decay products of the $\Upsilon(4 \mathrm{~S})$. Therefore, the two $B$ mesons are almost at rest in the center-of-mass (cms) system. One of the two $B$ mesons can now be tagged, with a hierarchical full reconstruction of 1104 hadronic decays [15]. The other, remaining $B$ meson could be decay into a $K$ meson and a charmonium or charmonium-like, exotic state $X_{c \bar{c}}$. Due to the peculiar kinematic situation, the missing mass $m_{\text {miss }}$ of the $K$ meson can be utilized for the measurements of total cross sections of any $X_{c \bar{c}}$, and in particular the $\mathrm{X}(3872)$. The missing mass is calculated as $m_{\text {miss }}(K)=\sqrt{\left.p_{e^{+} e^{-}}^{*}-p_{\text {tag }}^{*}-p_{K}^{*}\right)^{2}} c^{2}$ with cms momenta of the $e^{+} e^{-}$collision system, the tagging side and the $K$ meson, respectively. The recently measured value is $\mathcal{B}_{\text {total }}=1.2 \pm 1.1 \pm 0.1 \times 10^{-4}$ [14], leading to an upper limit of $\mathcal{B}_{\text {total }}<2.6 \times 10^{-4}$. Interestingly, this value almost reaches the value of the highest known partial product branching fraction

[^1][7] for $B^{+} \rightarrow K^{+} X(3872)$ with $X(3872) \rightarrow D^{0} \bar{D}^{0} \pi^{0}$ of $\mathcal{B}_{\text {partial }}=(1.0 \pm 0.4) \times 10^{-4}$. Thus, the measurement of $\mathcal{B}_{\text {total }}$ is in close reach for the ongoing data taking at Belle II.

## 5 Double $Z_{c}$ production

The charged $Z_{c}$ states have been observed in $e^{+} e^{-} \rightarrow Z_{c}^{+}(3900) \pi^{-}$at BESIII at $\sqrt{s}=4.26 \mathrm{GeV}$ [21]. In addition, the charged $Z_{b}$ states have been observed in $e^{+} e^{-} \rightarrow Z_{b}^{+} \pi^{-}$at Belle at higher $\sqrt{s}=10.86 \mathrm{GeV}$ in the bottomonium mass regime. Thus, it is intuitive to perform a search for $e^{+} e^{-} \rightarrow Z_{c}^{+} Z_{c}^{()-}$pair production at such higher cms energies of $\sqrt{s}=10.52 \mathrm{GeV}, 10.58 \mathrm{GeV}$ and 10.86 GeV , using the complete Belle data set. This search was performed at Belle [23], adding even 102 Mill. $\Upsilon(1 S)$ and 158 Mill. $\Upsilon(2 S)$ decays at $\sqrt{s}=9.46 \mathrm{GeV}$ and 10.02 GeV , respectively. $Z_{c}^{+} Z_{c}^{(\text {(')- }}$ refer to any possible combination of the $Z_{c}(3900)$ and the $Z_{c}(4200)$ states. No signal was observed, resulting, as an example, in an upper limit on the product branching fraction $\sigma \times \mathcal{B}\left(Z_{c}^{+}(3900) \rightarrow J / \psi \pi^{+}\right)$of 2.3 fb for the $Z_{c}(3900) Z_{c}(3900)$ combination. The order of magnitude of these upper limits is about a factor of $10^{6}$ smaller than typical cross sections of $\sigma\left(e^{+} e^{-} \rightarrow \Upsilon(5 \mathrm{~S})\right)=0.3 \mathrm{nb}$ at $\sqrt{s}=10.86 \mathrm{GeV}$ and about a factor $10^{4}$ smaller than $\sigma\left(e^{+} e^{-} \rightarrow \mathrm{Y}(4260)\right) \simeq 65 \mathrm{pb}$ at $\sqrt{s}=4.26 \mathrm{GeV}$, with $\Upsilon(5 \mathrm{~S}) \rightarrow Z_{b}^{()+} \pi^{-}$and $Y(4260) \rightarrow Z_{c}^{+}(3900) \pi^{-}$.

## 6 Search for the $P_{s}^{+}$pentaquark

The charmed pentaquark candidate $P_{c}^{+} \rightarrow J / \psi p$ has been observed in $\Lambda_{b}$ decays at LHCb [16]. For Belle (II), the $\Lambda_{b}$ is outside the kinematically accessible range. However, $\Lambda_{c}$ baryons are produced copiously and provide a suitable data set for the search for the strange partner of the $P_{c}^{+}$, called $P_{s}^{+}$, in the decay $P_{s}^{+} \rightarrow \phi p$ in $\Lambda_{c} \rightarrow[\phi p] \pi^{0}$. The search was performed at Belle [20] in a data set of $915 \mathrm{fb}^{-1}$ with total $1.468 .435 \pm 4816 \Lambda_{c}$ candidates. Fig. 2 shows the background subtracted $m(\phi p)$ distribution. Each data point corresponds to a fitted yield of a 2D fit. For details about the background subtraction and the fit procedure see [20]. No significant signal is observed. The fit result of $77.6 \pm 28.1 P_{s}$ events at a mass of $2.025 \pm 0.005 \mathrm{GeV} / c^{2}$ and a width of $0.022 \pm 0.012 \mathrm{GeV} / c^{2}$ leads to an upper limit of the product branching fraction of

$$
\mathcal{B}\left(\Lambda_{c}^{+} \rightarrow P_{s}^{+} \phi\right) \times \mathcal{B}\left(P_{s}^{+} \rightarrow \phi p\right)<8.3 \times 10^{-5}
$$

The determined upper limit for $P_{s}^{+}$is still about a factor of 6 higher than measured branching fraction for the $P_{c}^{+}$at LHCb .

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Figure 1. Invariant $m\left(J / \psi \pi^{+} \pi^{-}\right)$mass distributions of $B^{+} \rightarrow K^{+} J / \psi \pi^{+} \pi^{-}$decays (top) and $B^{0} \rightarrow K_{s}^{0} J / \psi \pi^{+} \pi^{-}$decays (bottom) at Belle, using the complete recorded data set of $711 \mathrm{fb}^{-1}$. Left: zoom to mass region of $\psi^{\prime}$, center: zoom to mass region of $X(3872)$, and right: zoom to mass region of $Y(4260)$.


Figure 2. Invariant $m(\phi p)$ mass distributions in $\Lambda_{c} \rightarrow \phi p \pi^{0}$ decays in the search for the $P_{s}^{+}$pentaquark state by Belle. For details about the background subtraction and the fit procedure see [20]. The red dotted line shows the phasespace distribution, which after background subtraction exhibits a negative fluctuation.

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[^1]:    ${ }^{1}$ In the Cornell potential, the radius can be approximated by the turning point of the wavefunction.

