



PHOTON 2015  
BINP, Novosibirsk

*New Physics  
Opportunities in Photon Induced  
interactions at the LHC*

*Victor P. Goncalves*

*High and Medium Energy Group*

*Physics and Mathematics Institute*

*UFPel - Brazil*

**Photon 2015 – BINP - Novosibirski  
15 June 2015**



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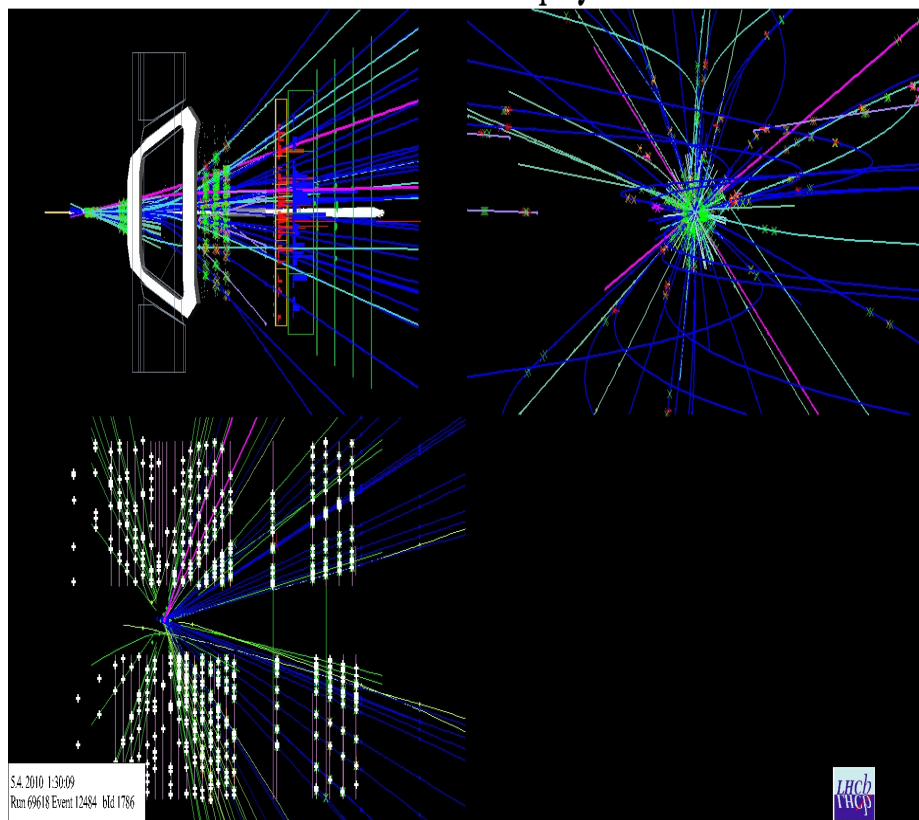
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# Motivation

Typical *pp* events:

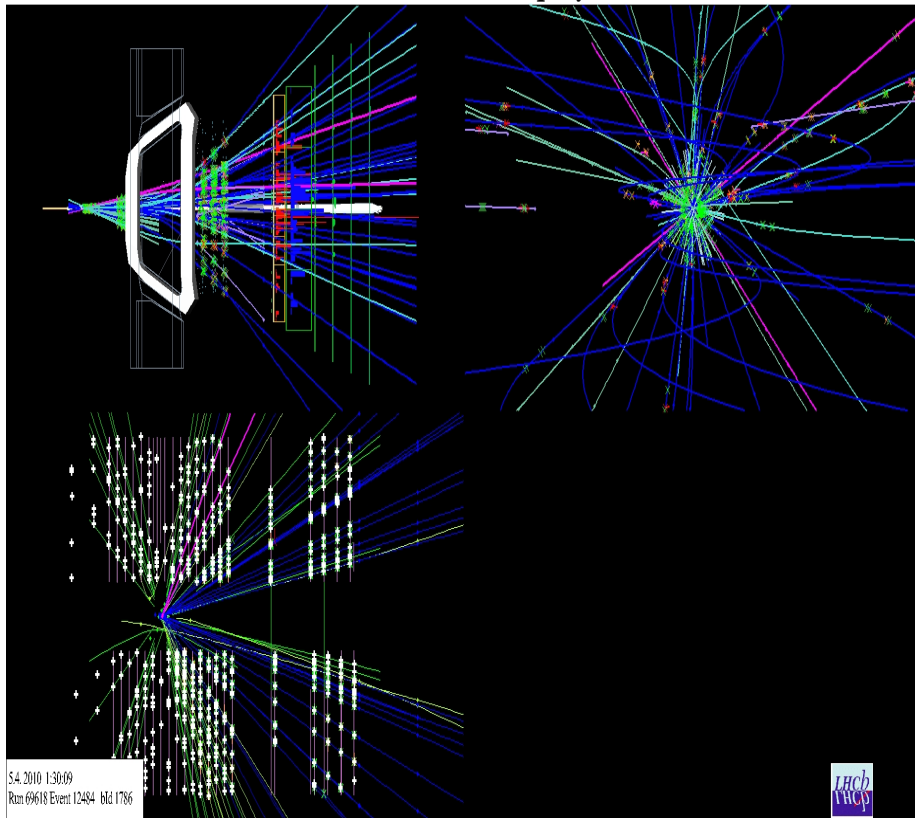
LHCb Event Display



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LHCb Event Display

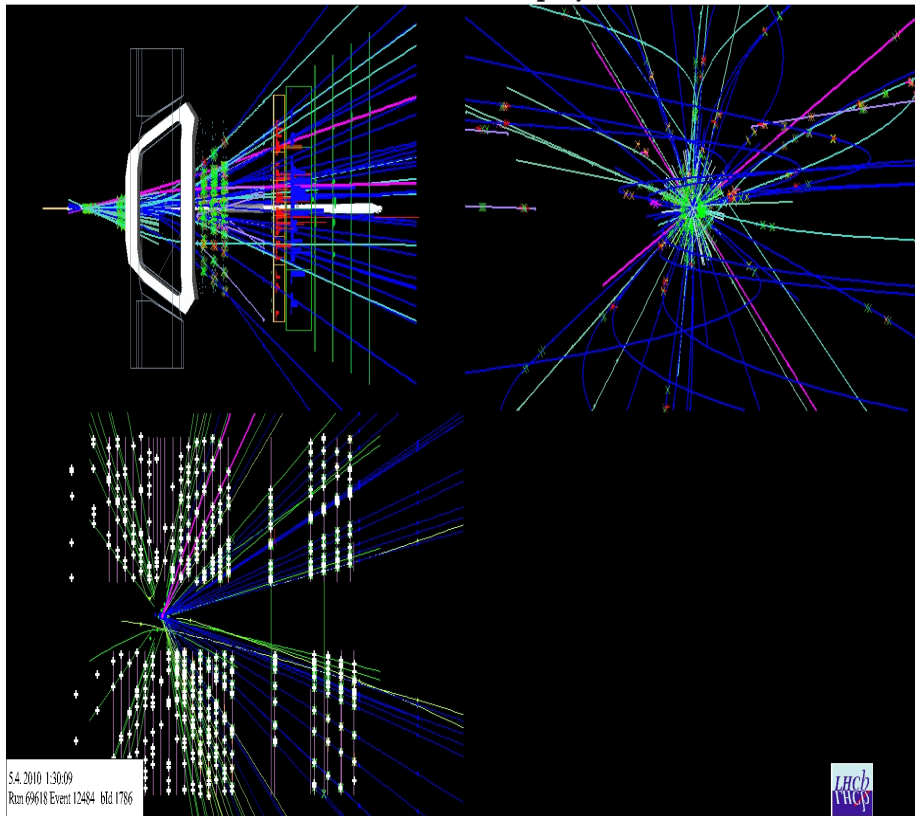


Many tracks + high  $p_T$  particles

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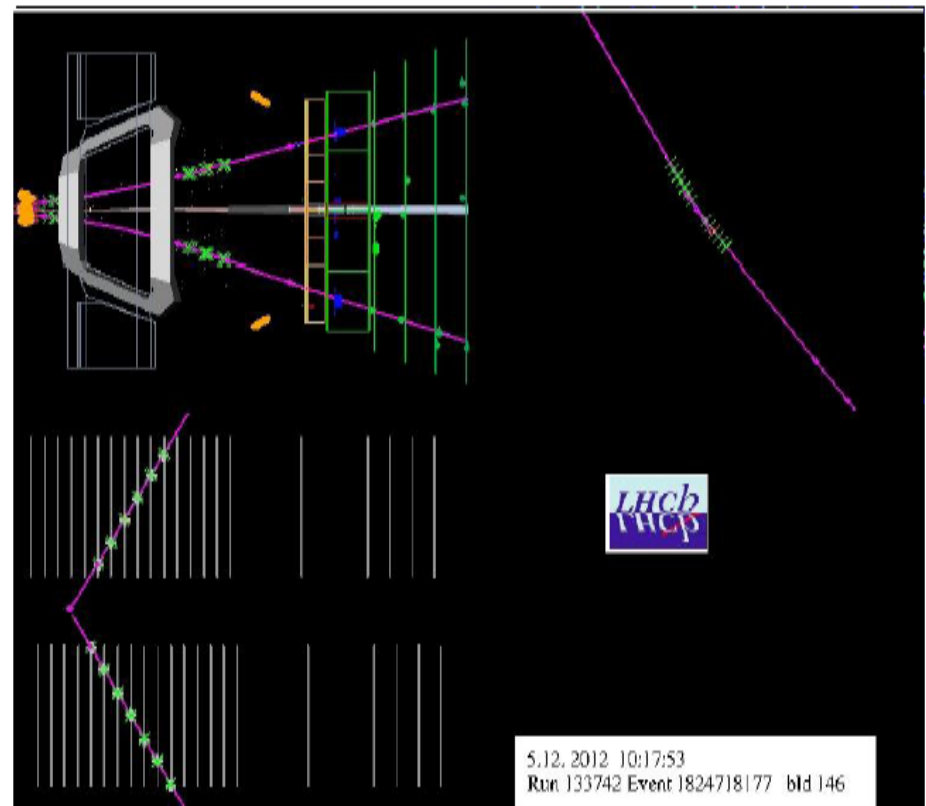
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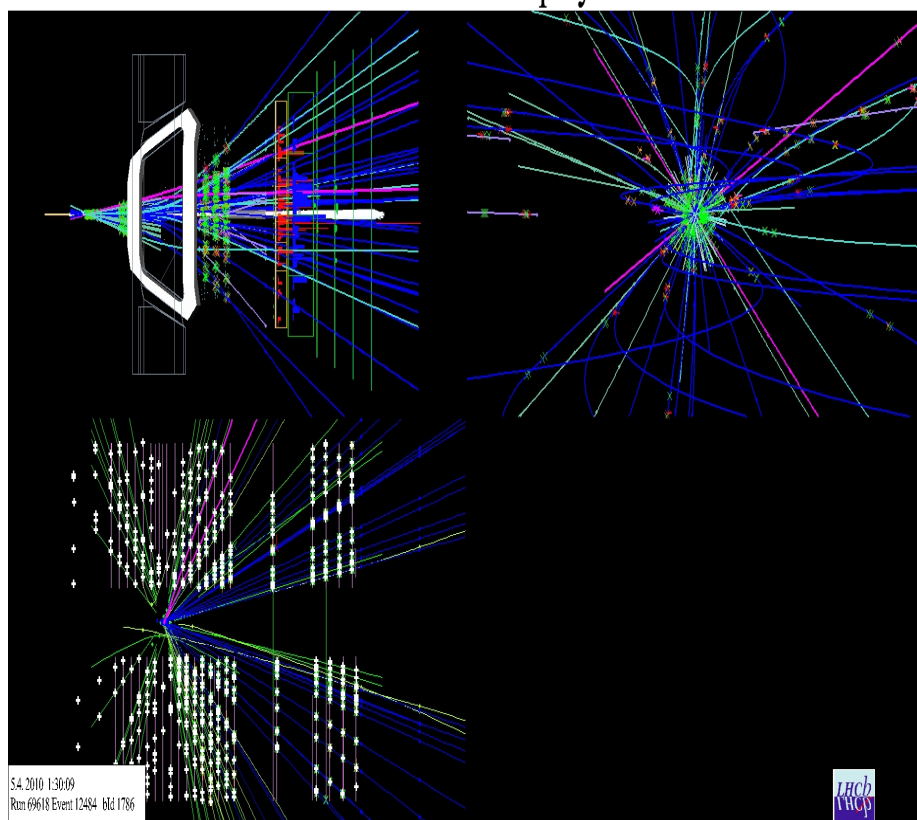




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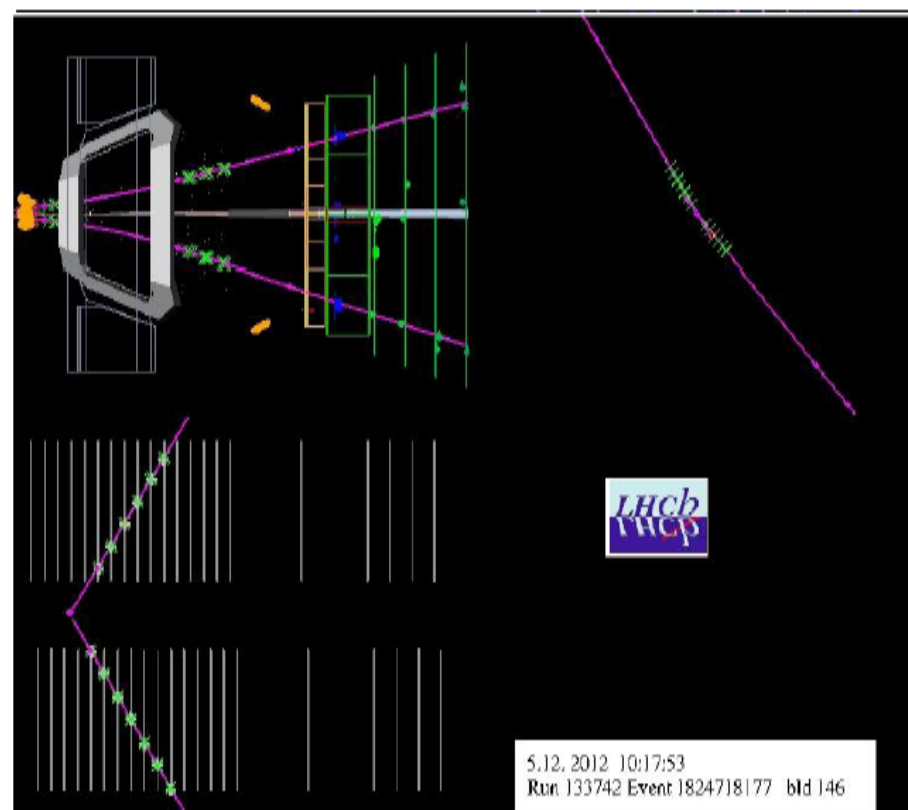
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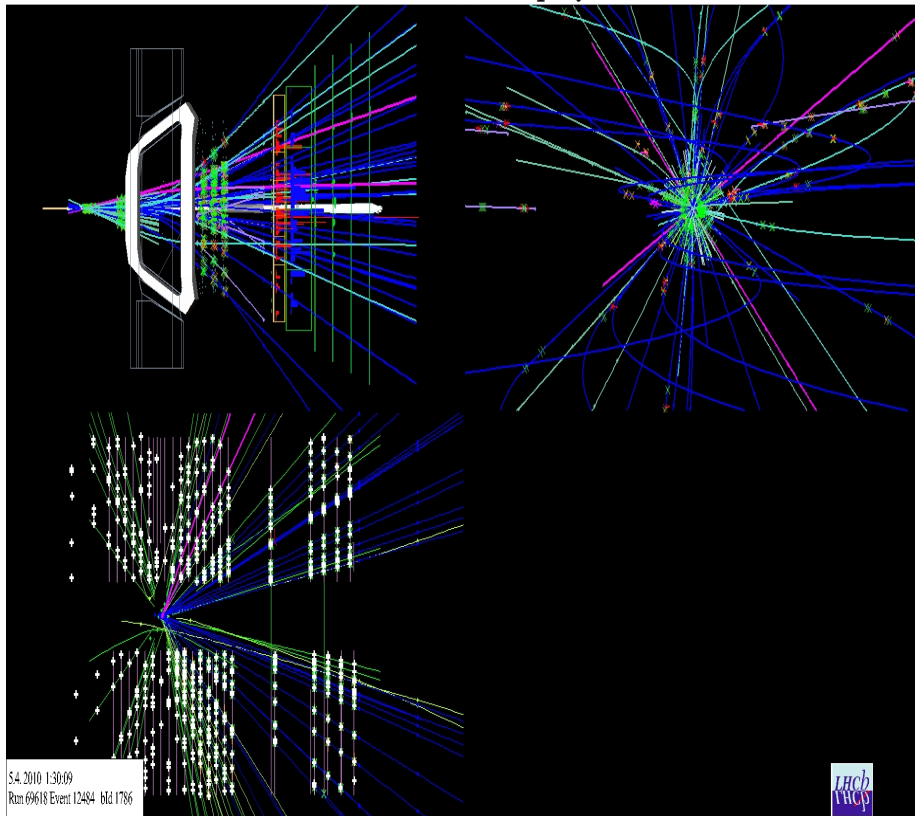


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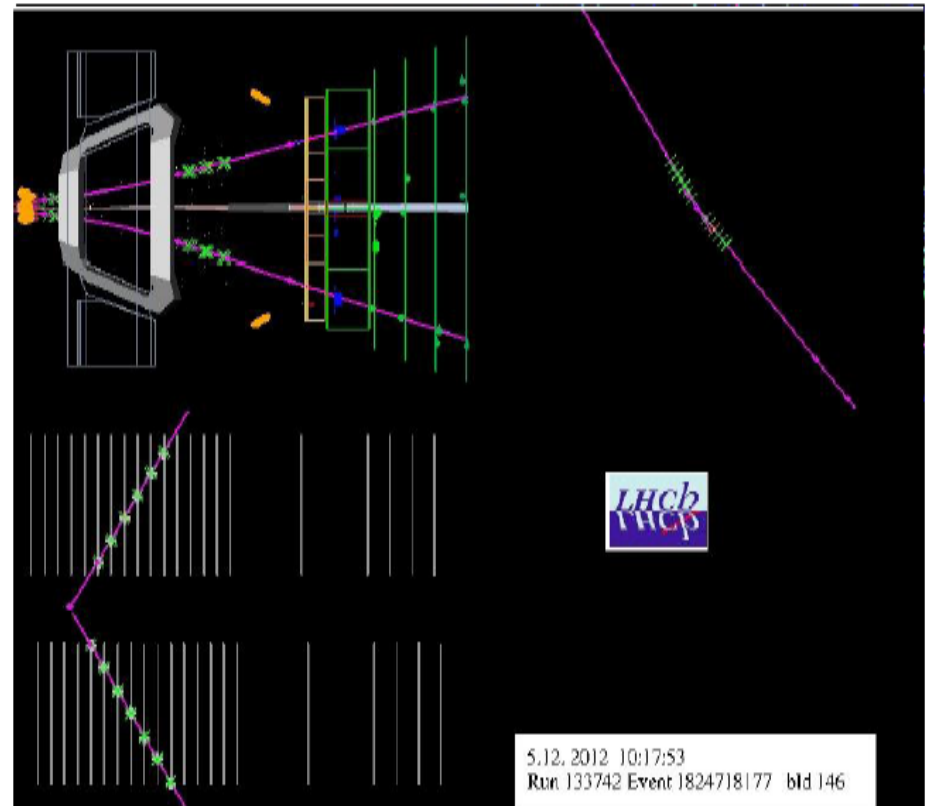
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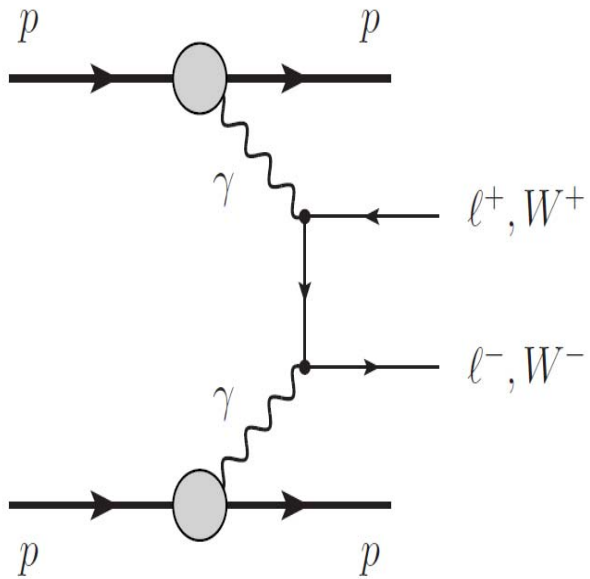
**Exclusive events:**



**Central Exclusive Process (CEP)**

# Central Exclusive Processes

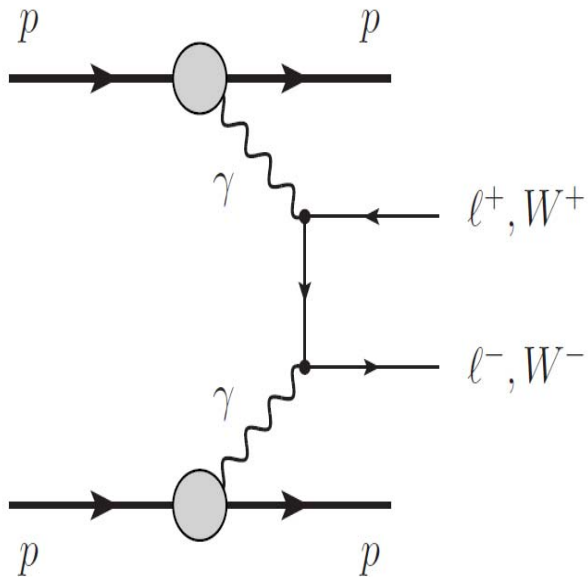
Two-photon interactions:





# Central Exclusive Processes

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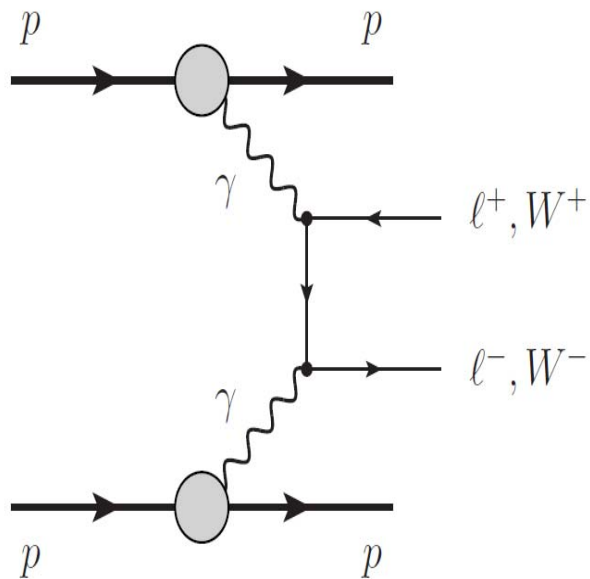
• Very clean processes: Central production with forward hadrons

• Accessible measurements:

1. Luminosity via dilepton production ( $\gamma\gamma \rightarrow \mu^+\mu^-$ );
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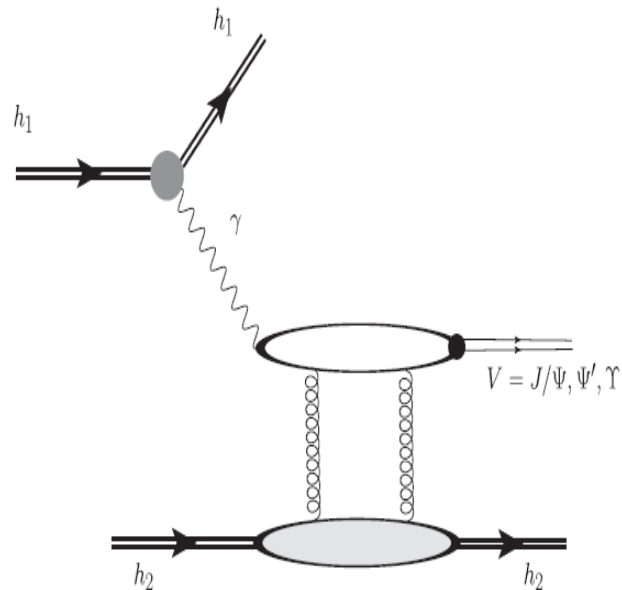


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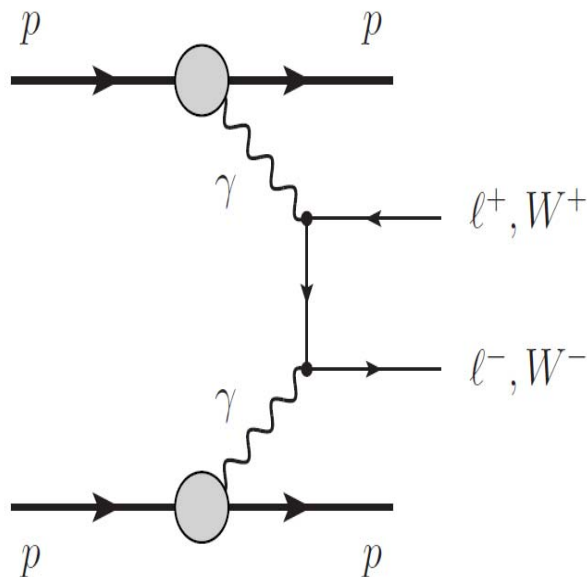
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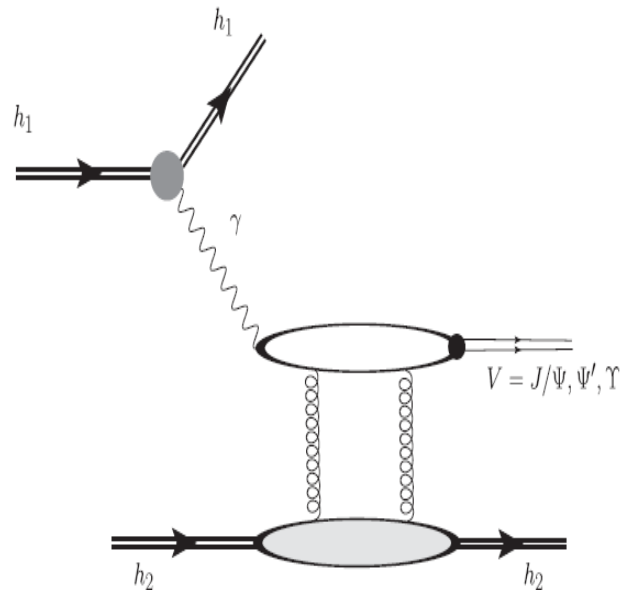


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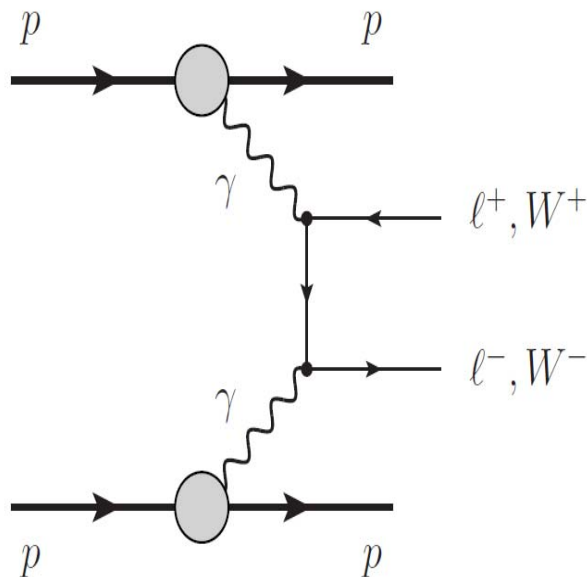
## Photon - Pomeron:



- Allow us to study the QCD dynamics at small- $x$ .
- Sensitive to the description of diffraction.
- Determination of the gluon distribution and the magnitude of the shadowing effects.
- Search for saturation effects.
- Search for Odderon, Charmoniumlike exotic states, ...

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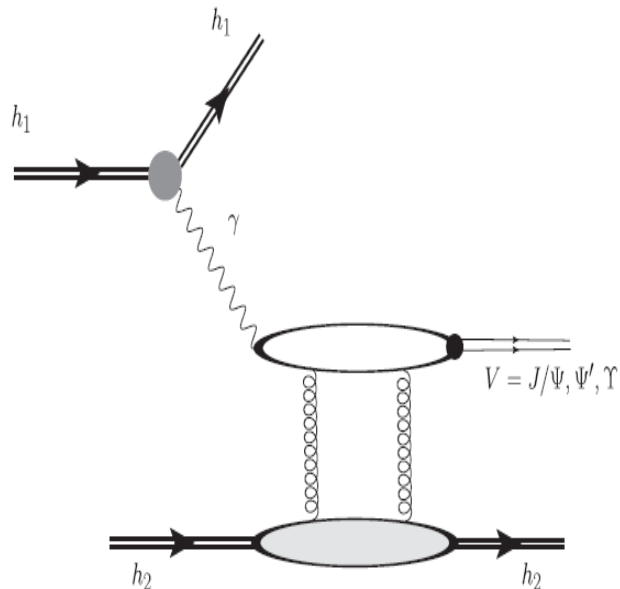


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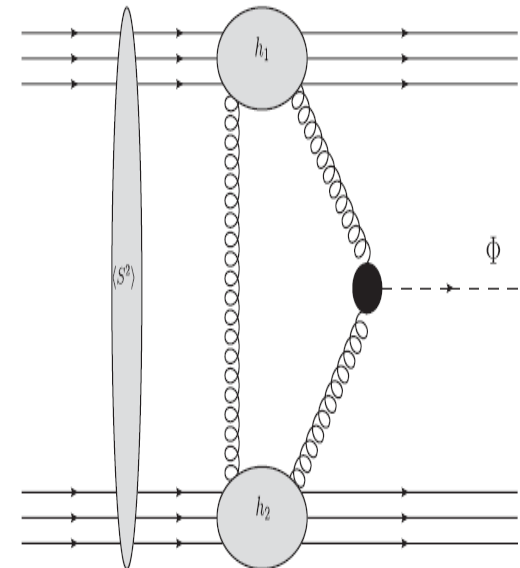
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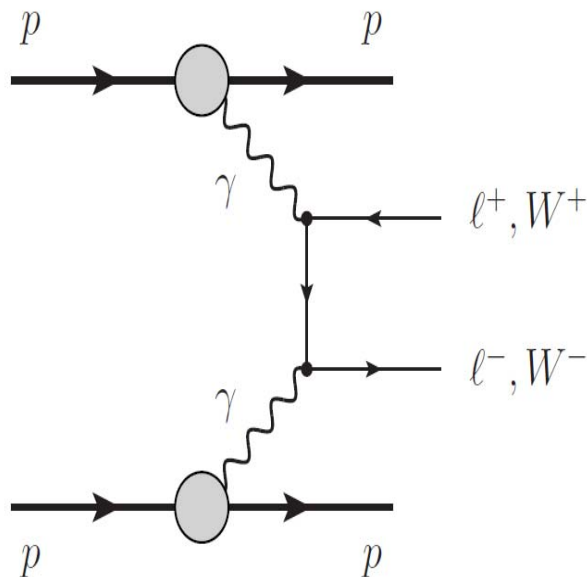
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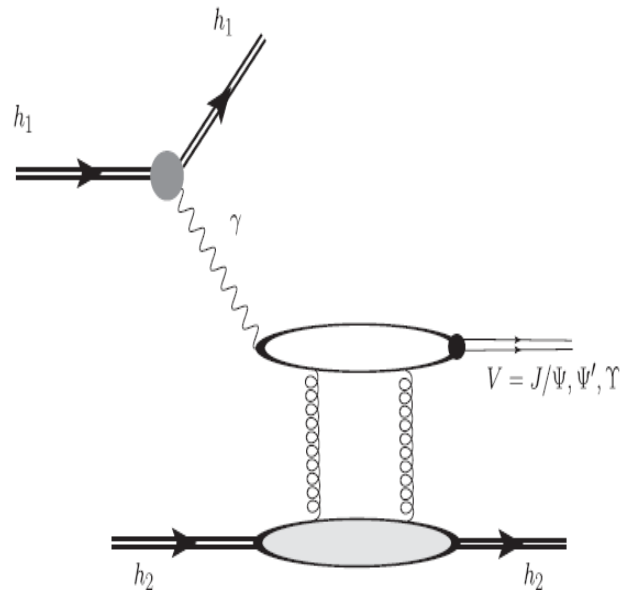
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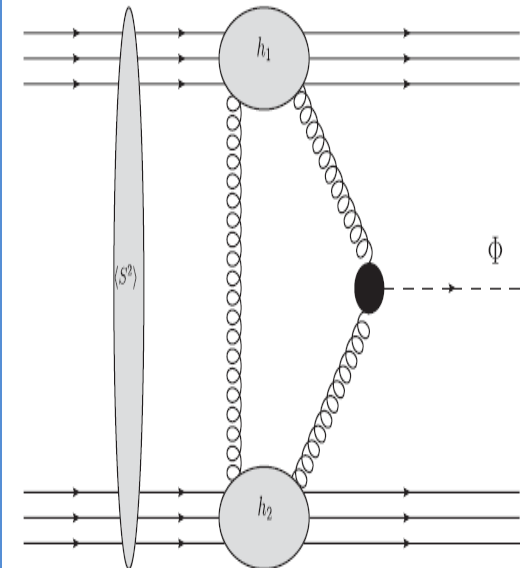
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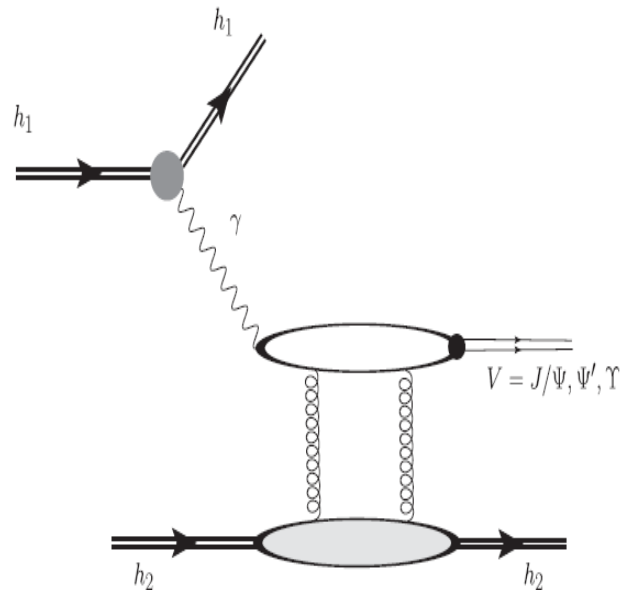
## Pomeron - Pomeron:



- Spin - parity analyser: only a subset of resonant states can be produced. In particular  $0^{++}$  but not, for example,  $1^{++}$ .
- Sensitive to the description of diffraction.
- Very sensitive to beyond Standard Model Physics.

# Central Exclusive Processes

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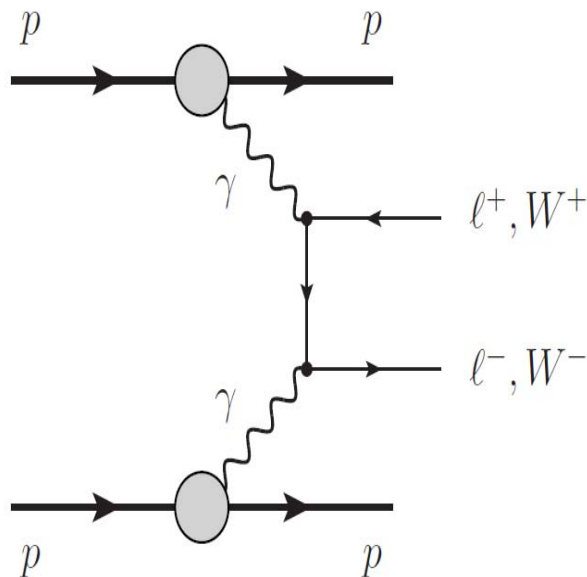


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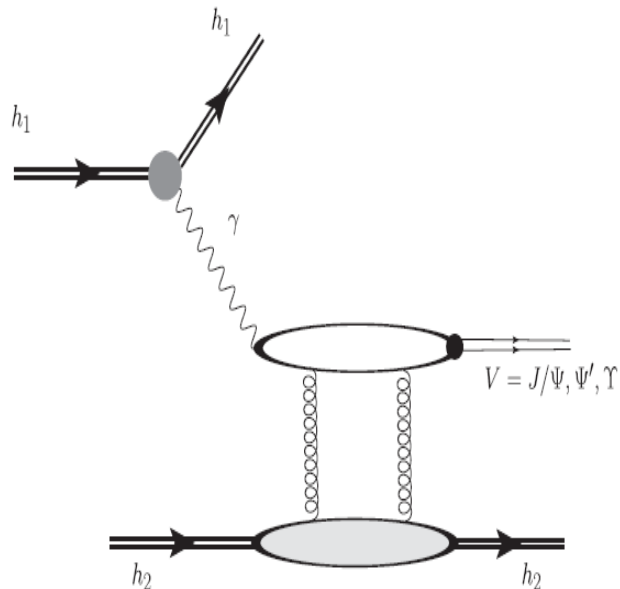


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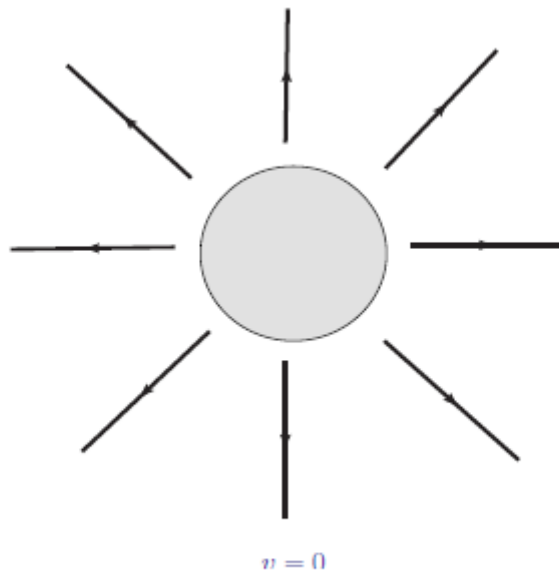
# **Photon – Induced Interactions:**

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## **The Equivalent Photon Approximation**

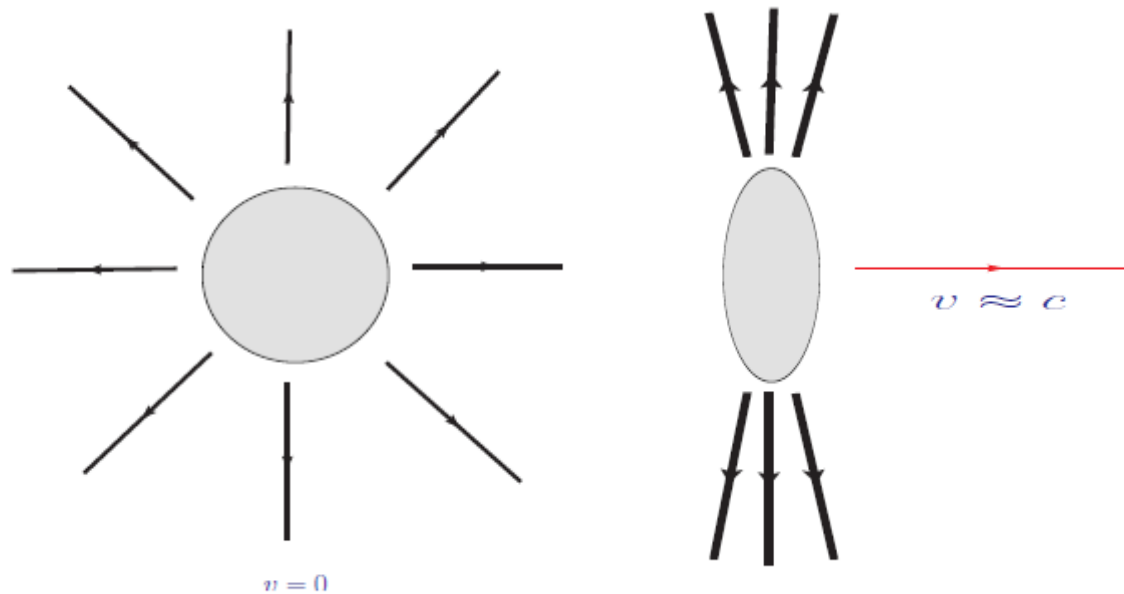
# Photon – Induced Interactions: The Equivalent Photon Approximation

- Consider a charged nucleus at rest. The associated electromagnetic field can be represented by:



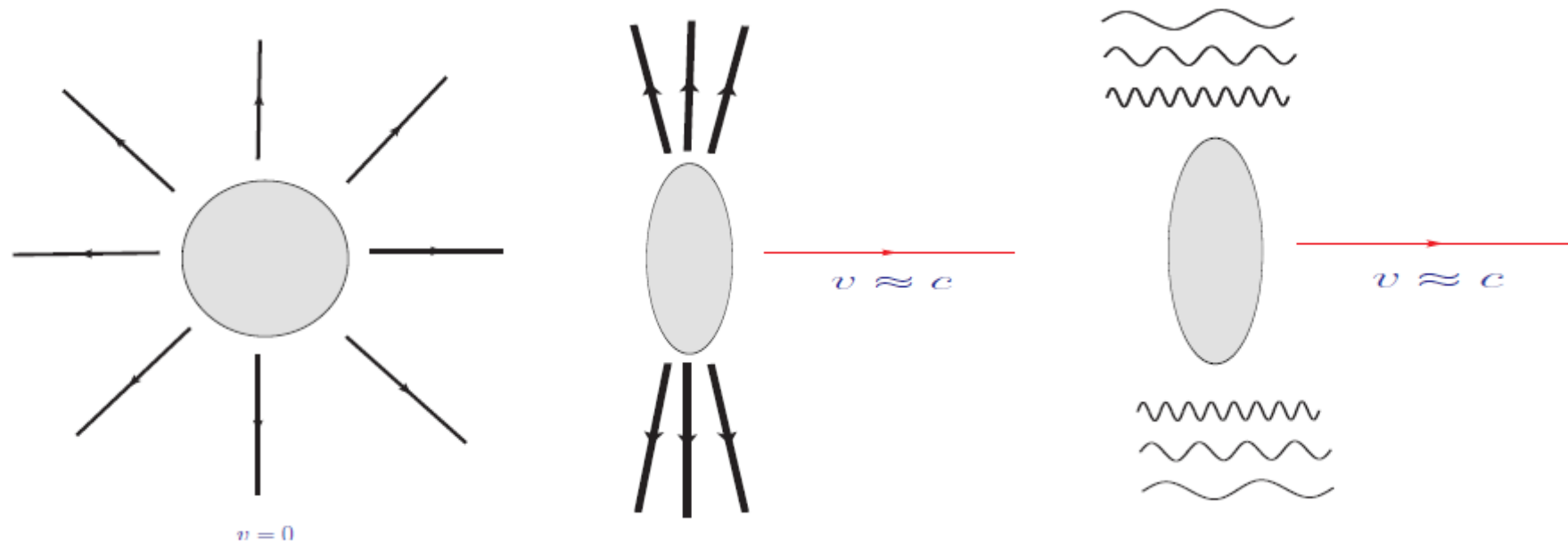
# Photon – Induced Interactions: The Equivalent Photon Approximation

- As a charged nucleus moves with nearly the speed of light, the electromagnetic field becomes transverse to its velocity.



# Photon – Induced Interactions: The Equivalent Photon Approximation

- Since the electric and magnetic field associated to the nucleus take on the same absolute value, this transverse electromagnetic field can be simulated by an equivalent swarm of photons <sup>a</sup>.



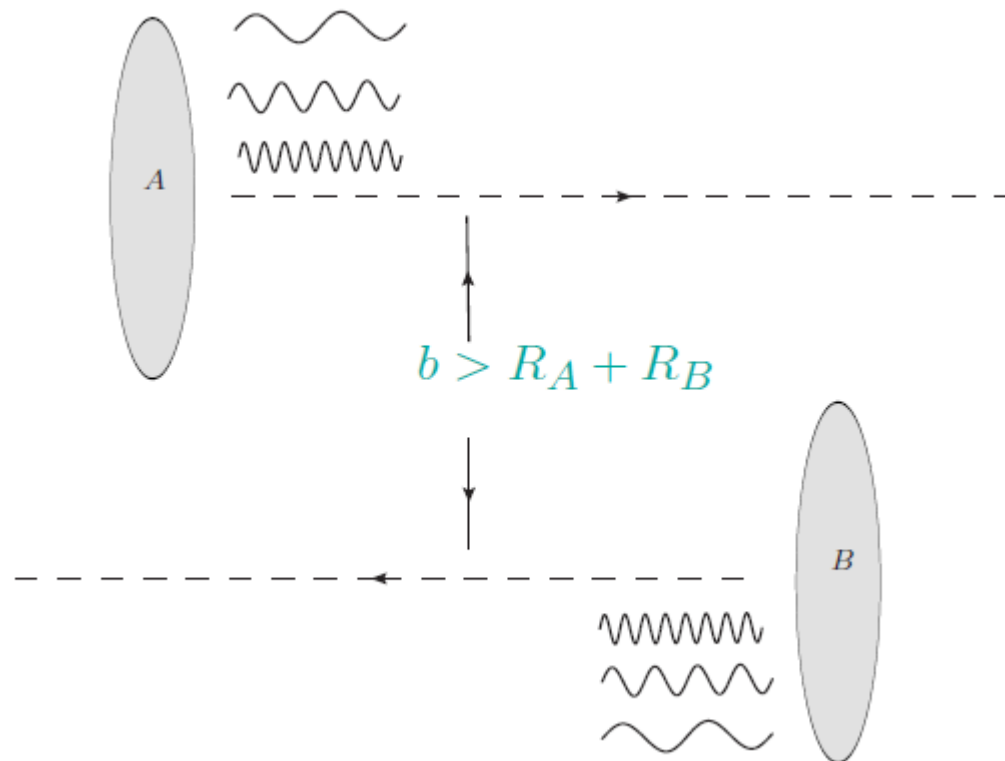
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<sup>a</sup>E. Fermi (1924), E. J. Williams (1933), C. F. Von Weizacker (1934)



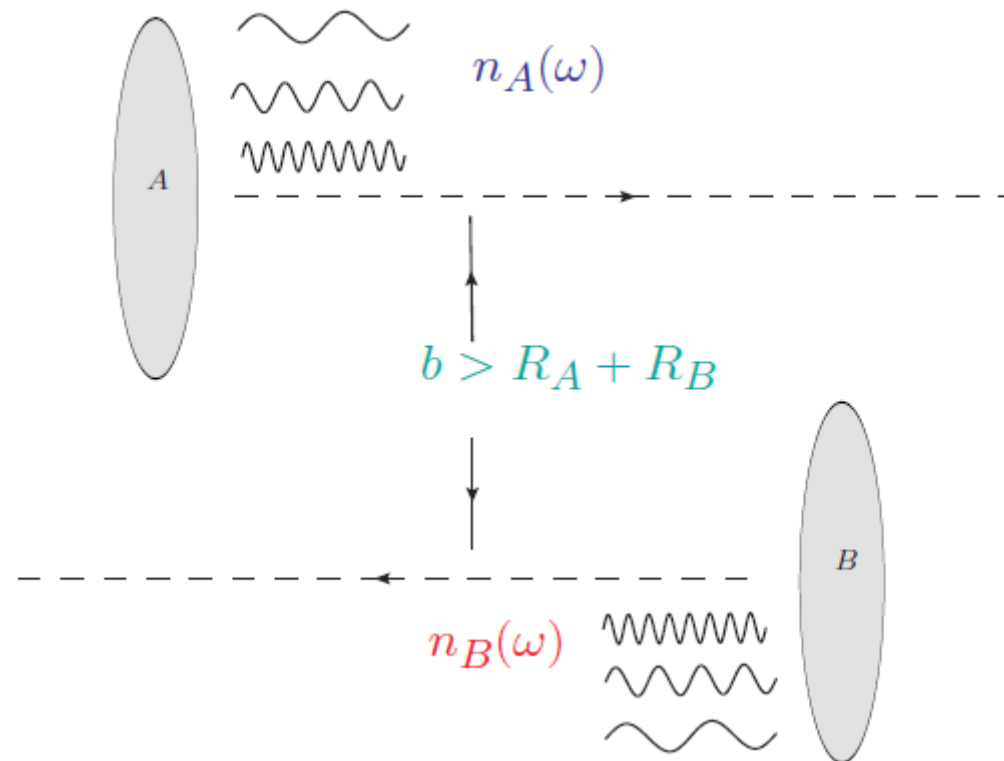
# Photon – Induced Interactions: The Equivalent Photon Approximation

- Thus the collision of two charged nuclei at large impact parameter can be described as the collision of two equivalent swarms of photons.



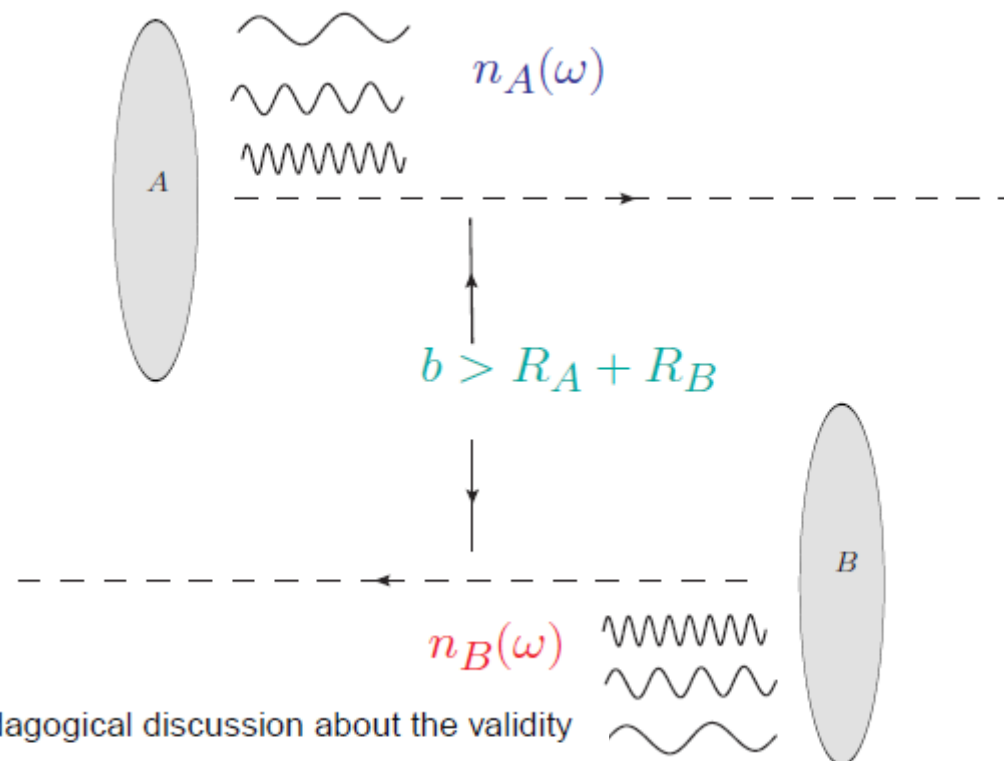
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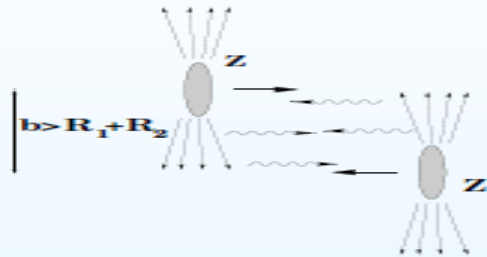
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- For a detailed and pedagogical discussion about the validity of the EPA see Budnev, Ginzburg, Meledin and Serbo, Phys. Rep. 15, 181 (1975).

# Photon – Induced Interactions:

## Center of mass energies



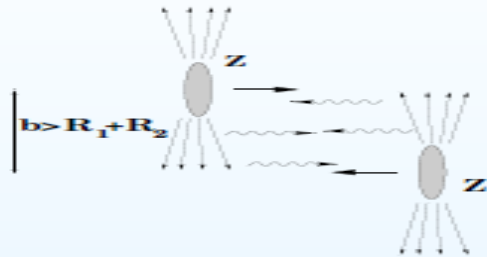
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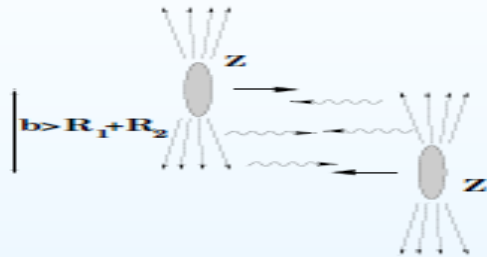


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# Photon – Induced Interactions: Center of mass energies



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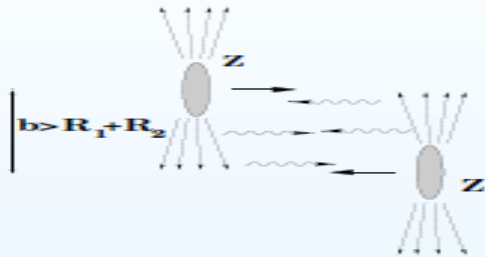
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- Photoproduction in  $pp$  collisions at LHC probes energies one order of magnitude larger than HERA.



# Photon – Induced Interactions: Center of mass energies

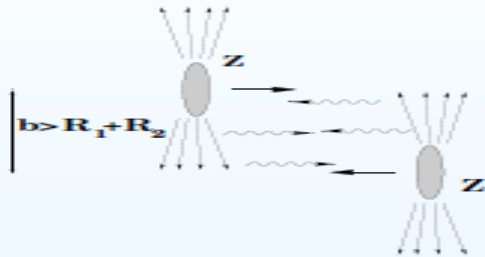


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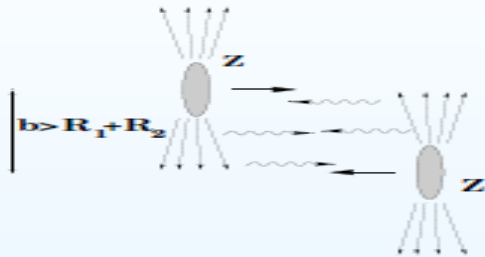
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- Photoproduction in  $pA$  and  $AA$  collisions probes an unexplored regime of center of mass energies.

# Photon – Induced Interactions: Center of mass energies



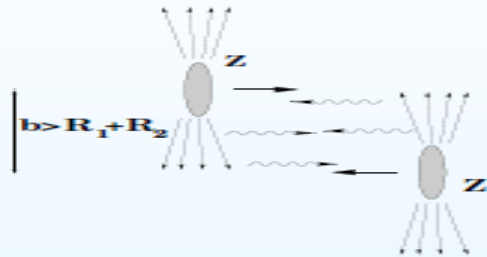
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- $\gamma\gamma$  interactions with center of mass energies larger than those obtained at LEP - CERN.

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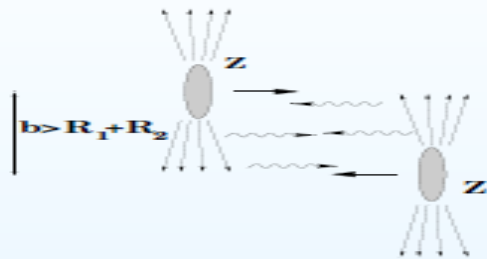


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- $\gamma\gamma$  interactions with center of mass energies larger than those expected in the future ILC.

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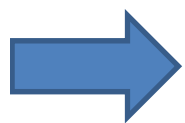
## Center of mass energies



Center of mass energies

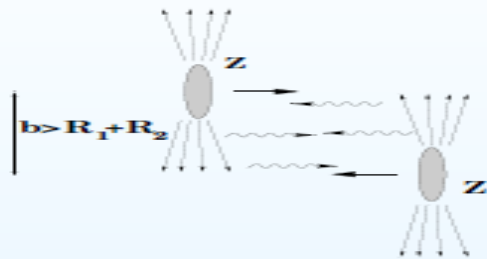
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The LHC is the world's most powerful collider not only for protons and lead ions but also for  $\gamma\gamma$  and  $\gamma h$  collisions.

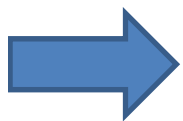
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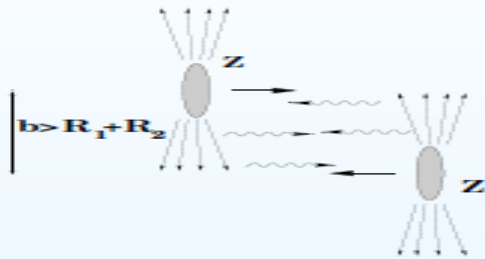


- Photon - induced interactions at LHC allows to study Quantum Chromodynamics in an unexplored regime of center of mass energies.



# Photon – Hadron Interactions at the LHC:

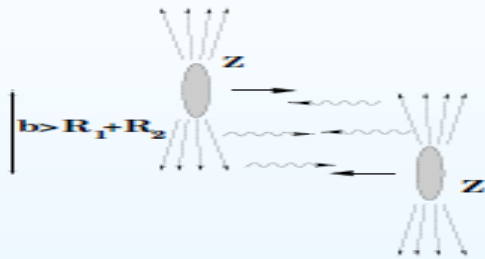
## Typical processes



$\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$

# Photon – Hadron Interactions at the LHC:

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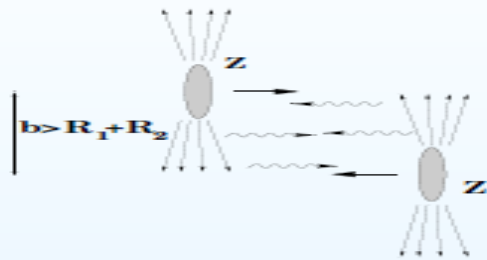
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The final state is characterized by **one rapidity gap** due to the dissociation of the hadron target ( $pp \rightarrow p \otimes XY$ ).

# Photon – Hadron Interactions at the LHC:

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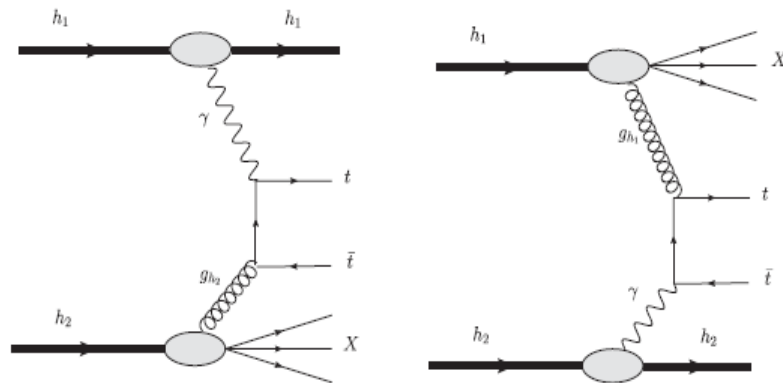


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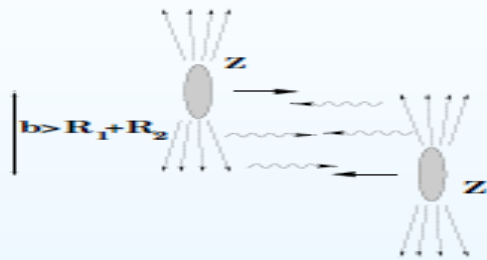
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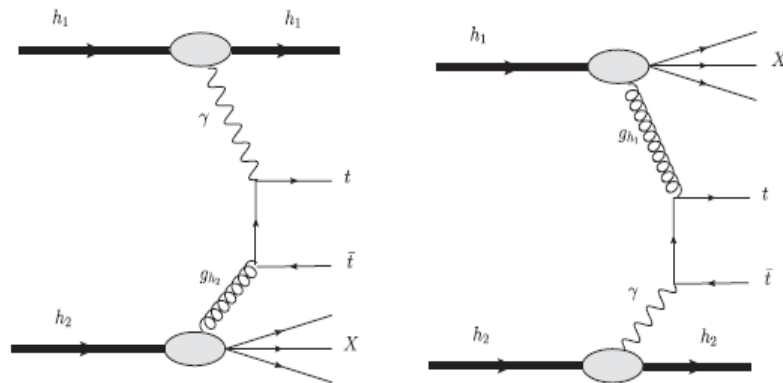


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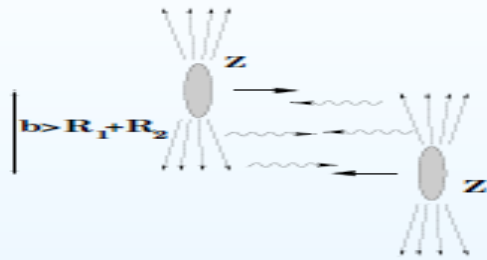
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## Typical processes

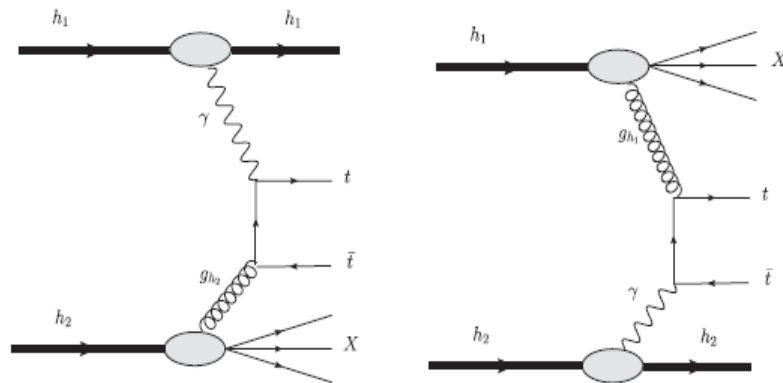


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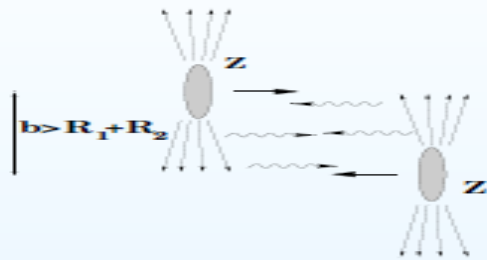
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TABLE I. The integrated cross section (events rate) for the photoproduction of top quarks in  $pp$ ,  $pPb$ , and  $PbPb$  collisions at LHC energies.

$pp$	MRST	CT10
$\sqrt{s} = 8$ TeV	0.739 pb (73900)	0.764 pb (76400)
$\sqrt{s} = 14$ TeV	2.50 pb (250000)	2.53 pb (253000)
$pPb$	MRST	MRST + EPS09
$\sqrt{s} = 5.5$ TeV	0.036 nb (5.4/3600)	0.038 nb (5.7/3800)
$\sqrt{s} = 8.8$ TeV	0.159 nb (23.85/15900)	0.165 nb (24.75/16500)
$PbPb$	MRST	MRST + EPS09
$\sqrt{s} = 5.5$ TeV	0.42 nb (0.18)	0.40 nb (0.17)

# Photon – Hadron Interactions at the LHC:

## Typical processes



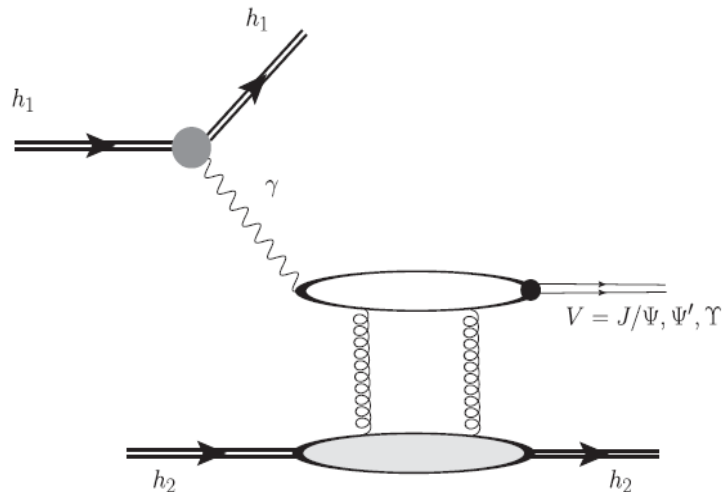
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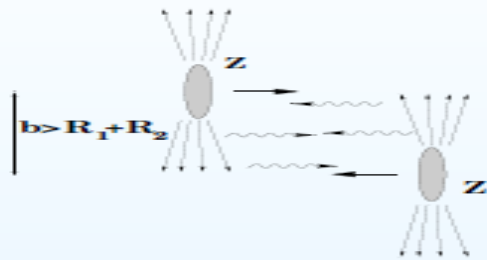
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# Photon – Hadron Interactions at the LHC:

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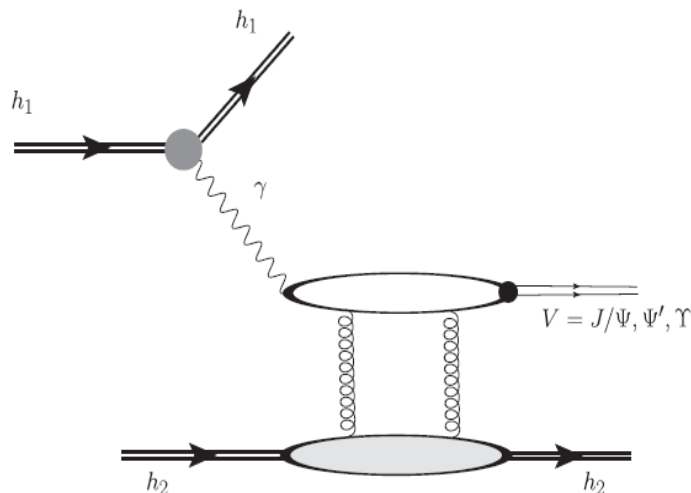
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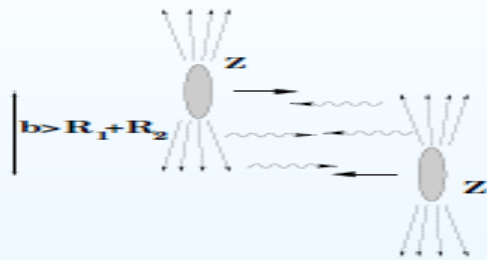
$(pp \rightarrow p \otimes X \otimes p)$ .



- Cross section is proportional to the **square** of the proton/nuclear gluon distribution.

# Photon – Hadron Interactions at the LHC:

## Typical processes



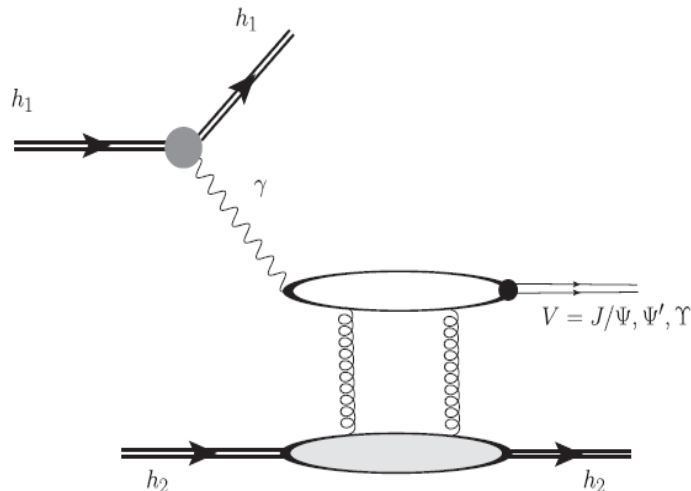
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- Cross section is proportional to the **square** of the proton/nuclear gluon distribution.

- Diffractive vector meson photoproduction in UPHIC is a **probe** of the gluon distribution <sup>a</sup>

<sup>a</sup>VPG, Bertulani, PRC65, 054905 (2002)



# **Photon – Hadron Interactions at the LHC:** **Constraining the nuclear gluon distribution**

# Photon – Hadron Interactions at the LHC: Constraining the nuclear gluon distribution

$$R_g \equiv \frac{xg_A(x, Q^2)}{A \cdot xg_p(x, Q^2)}$$

# Photon – Hadron Interactions at the LHC: Constraining the nuclear gluon distribution

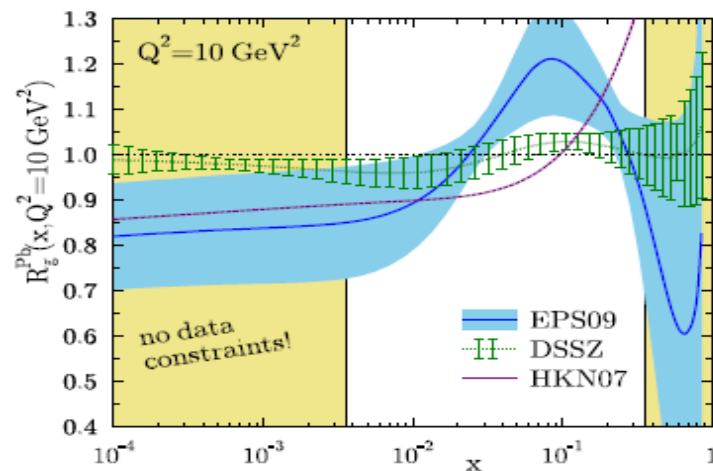
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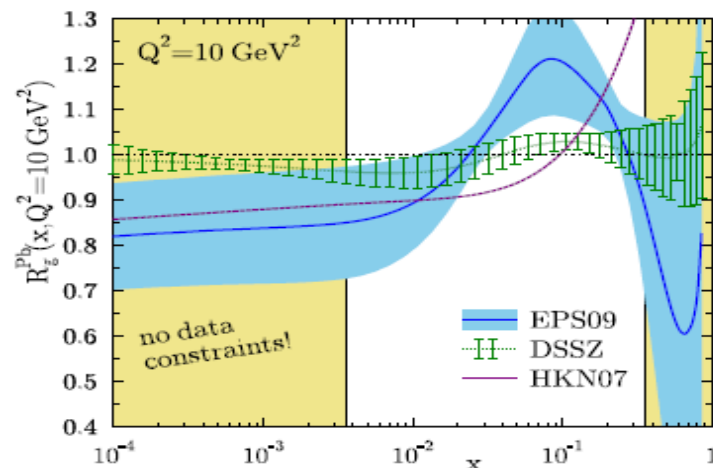
Eskola, Puukkunen, arXiv:1401.2345

- The current  $eA$  experimental data does not constrain the small- $x$  behaviour.
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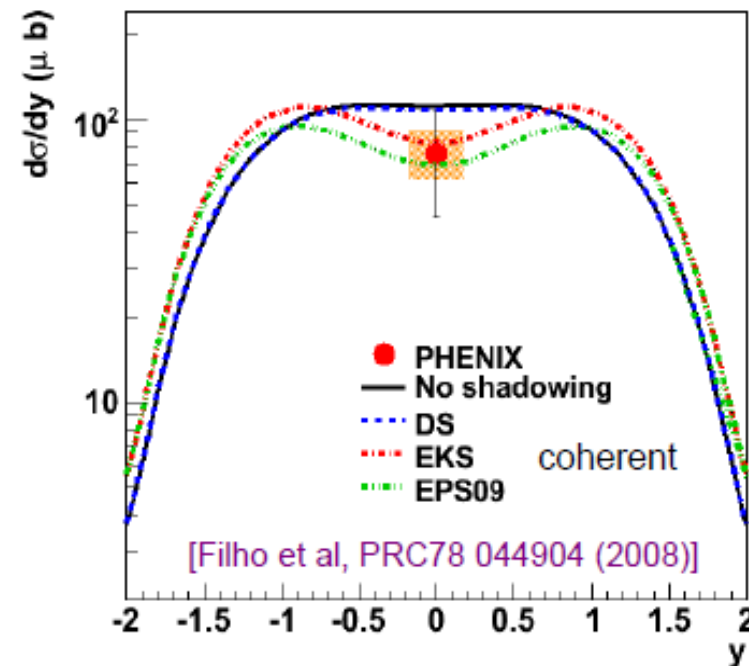
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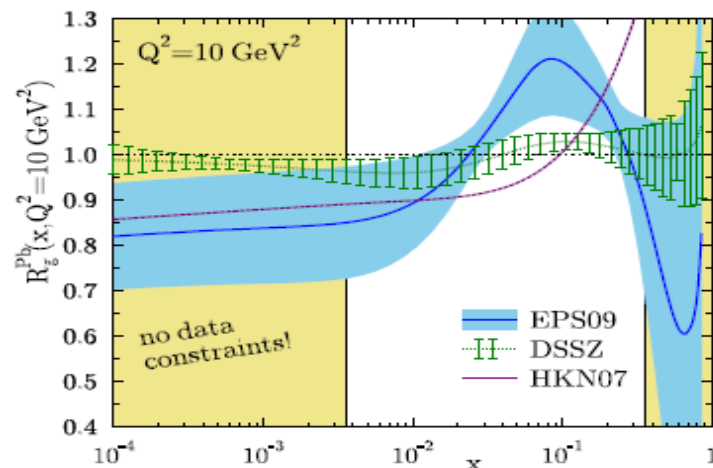


[Filho et al, PRC78 044904 (2008)]

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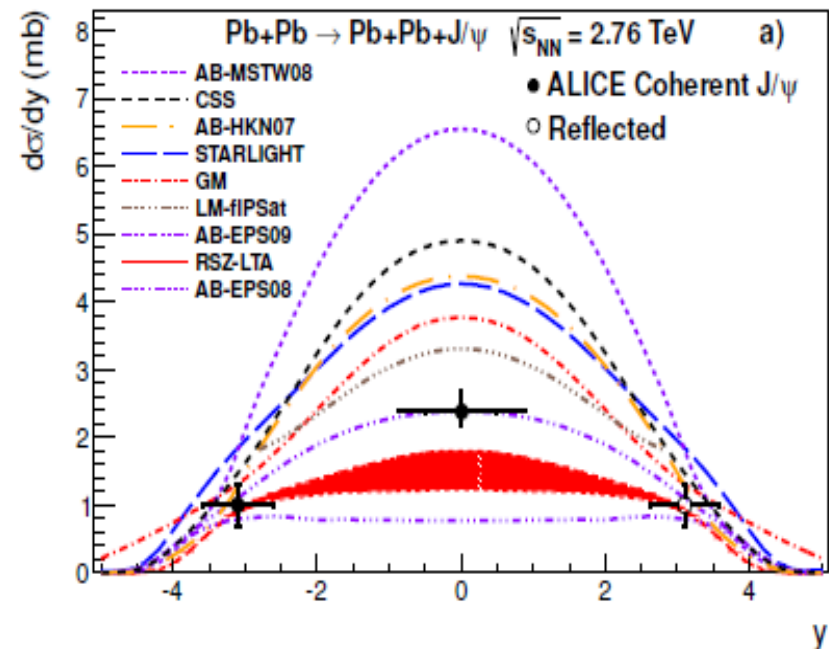
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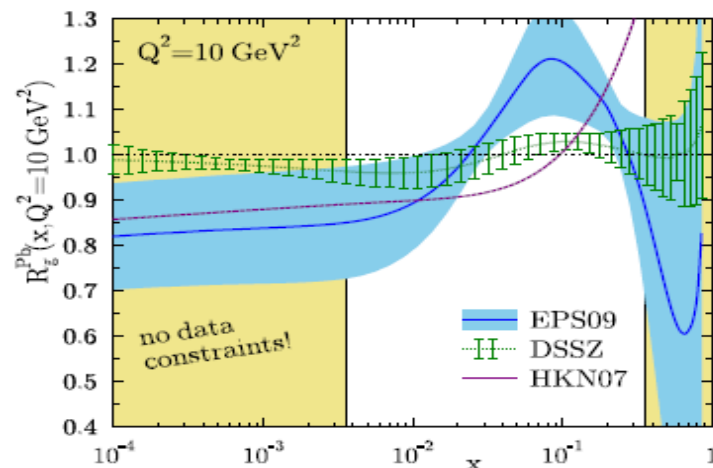
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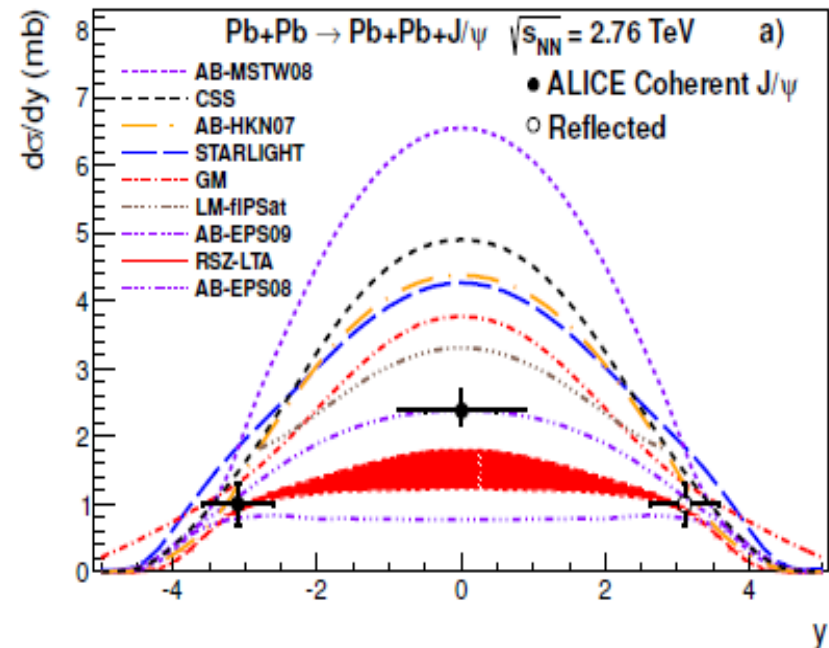
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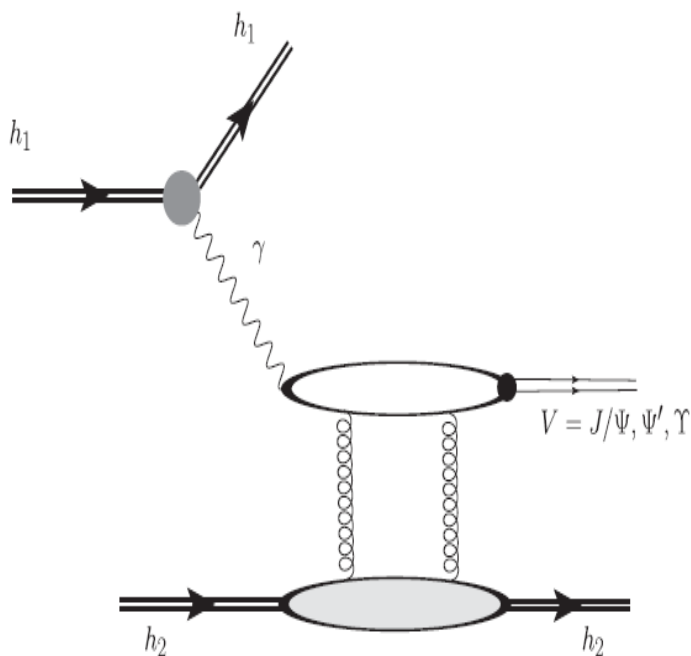
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- ALICE gives the first evidence of large nuclear shadowing effect at small- $x$ .

# Photon – Hadron Interactions at the LHC: Probing the QCD dynamics

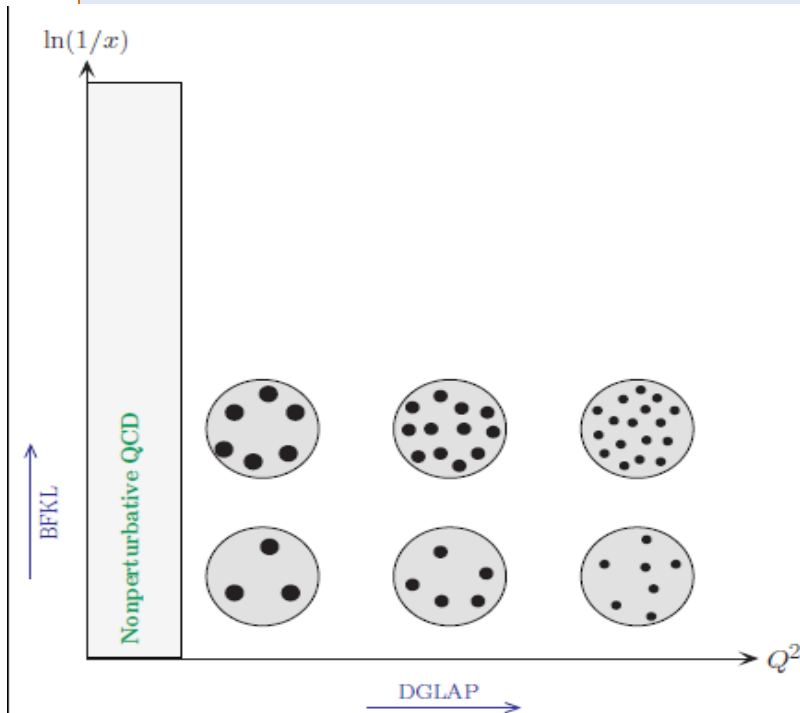


- Diffractive vector meson photoproduction in photon - induced interactions is a **probe** of the nonlinear effects in the QCD dynamics at high energies and the vector meson wave function <sup>a</sup>.

<sup>a</sup>VPG, Machado, EPJC 40, 519 (2005)



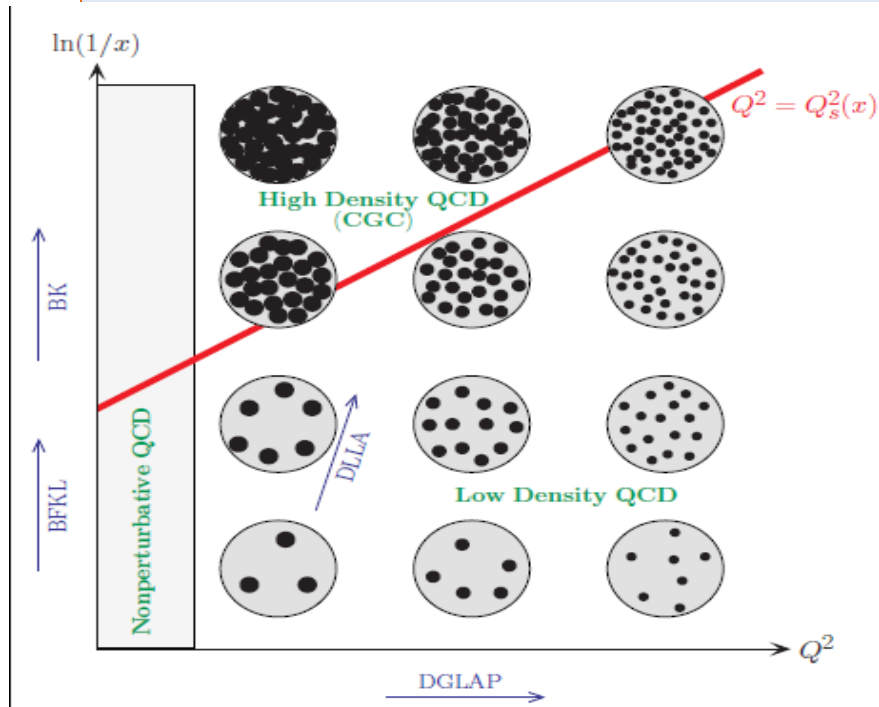
# Photon – Hadron Interactions at the LHC: Probing the QCD dynamics



Linear QCD evolution equations predict a power growth of gluon distribution as  $x \rightarrow 0$  (violates unitarity).

# Photon – Hadron Interactions at the LHC:

## Probing the QCD dynamics

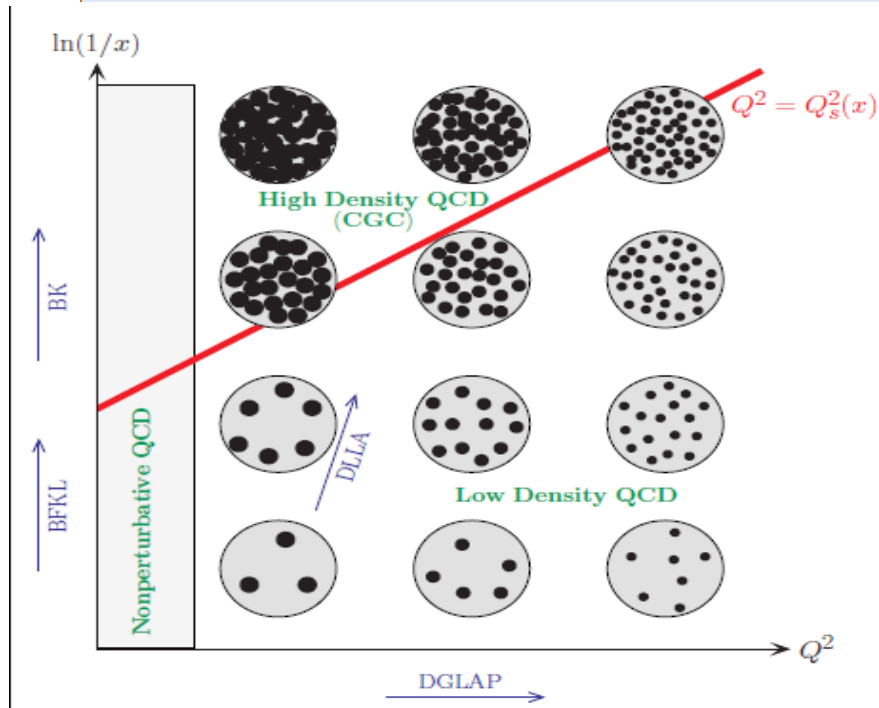


Open questions: Are present at RHIC/LHC/Cosmic Rays? What is the magnitude of these effects? What is the more adequate description of the nonlinear effects?

- Linear QCD evolution equations predict a power growth of gluon distribution as  $x \rightarrow 0$  (violates unitarity).
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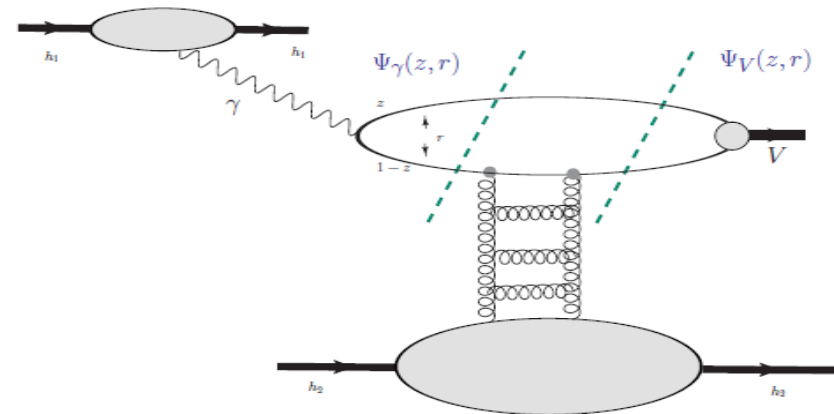
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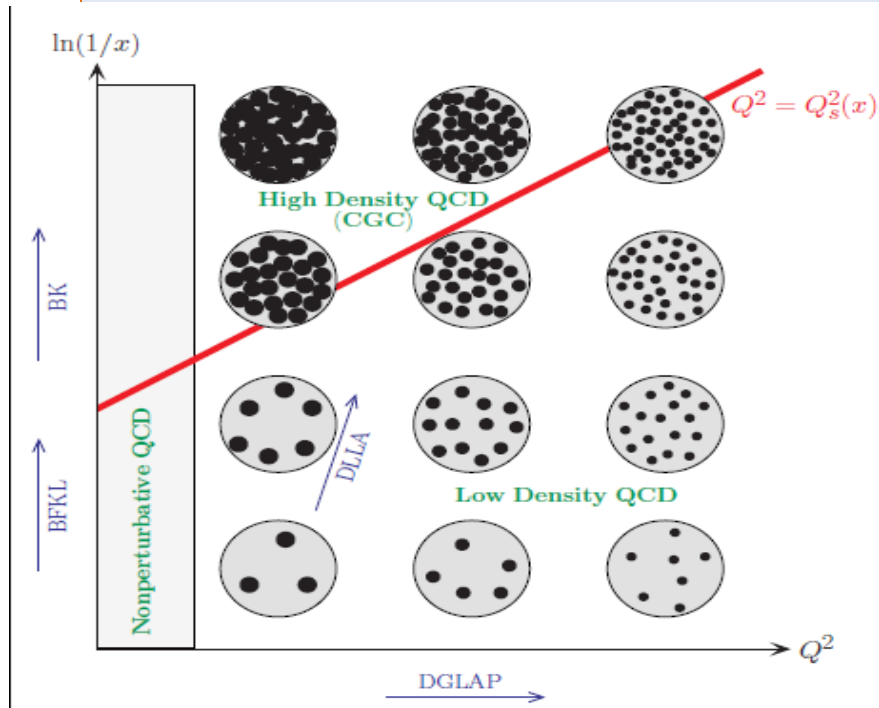


$$\mathcal{A}^{\gamma p \rightarrow J/\Psi p}(x, \Delta) = i \int dz d^2r d^2b e^{-i[b-(1-z)r] \cdot \Delta} (\Psi_{J/\Psi}^* \Psi) 2\mathcal{N}_p(x, r, b)$$

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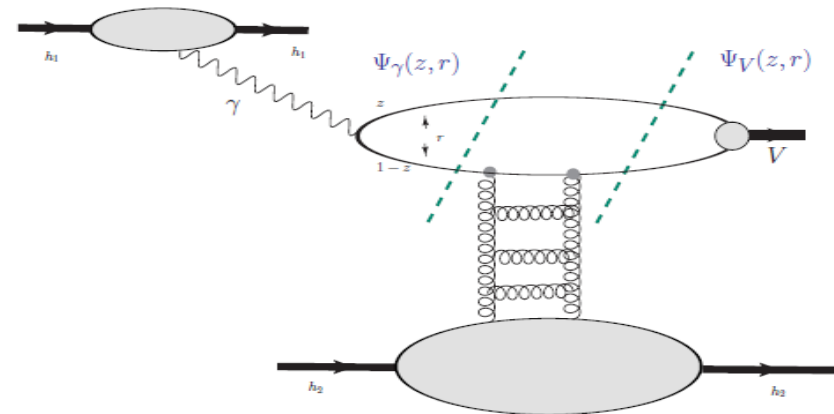
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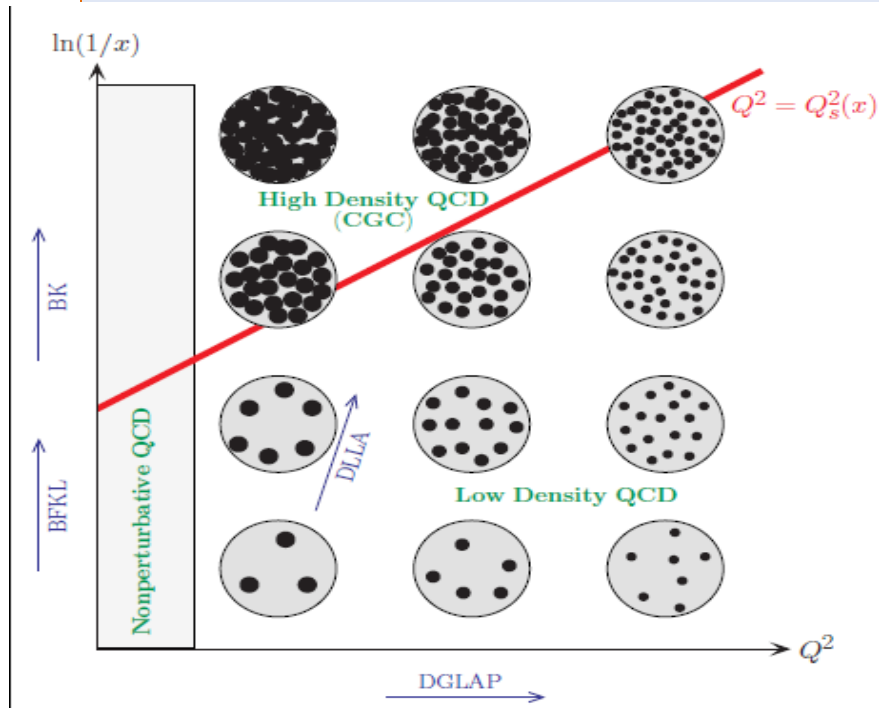


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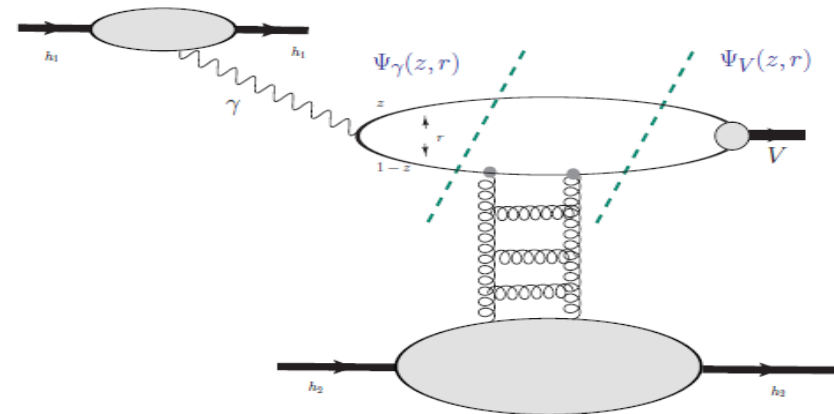
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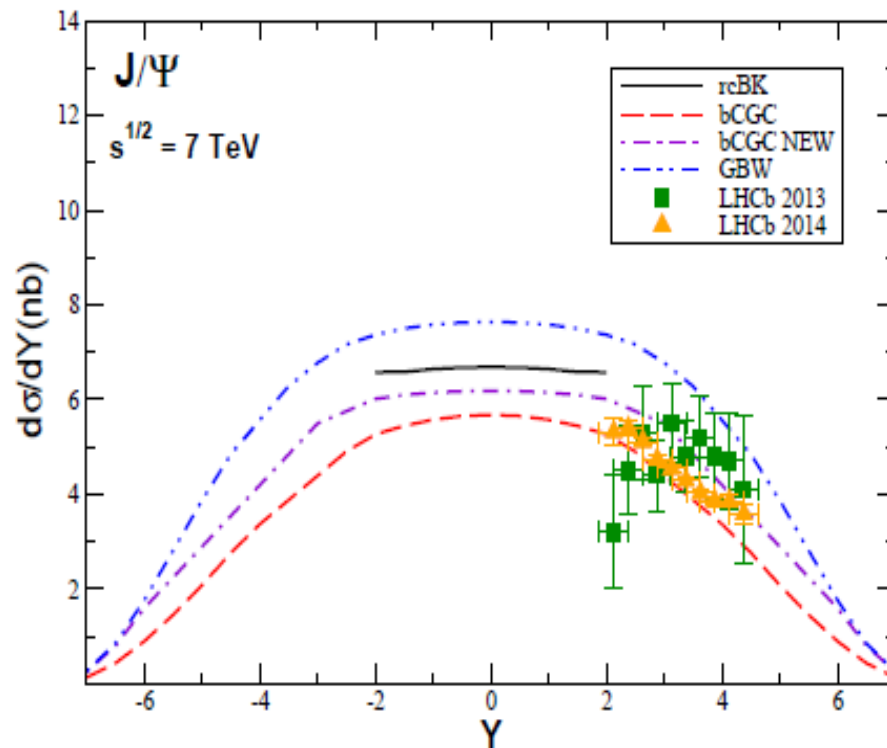
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<sup>a</sup>VPG, Moreira, Navarra, PRC90, 15203 (2014)

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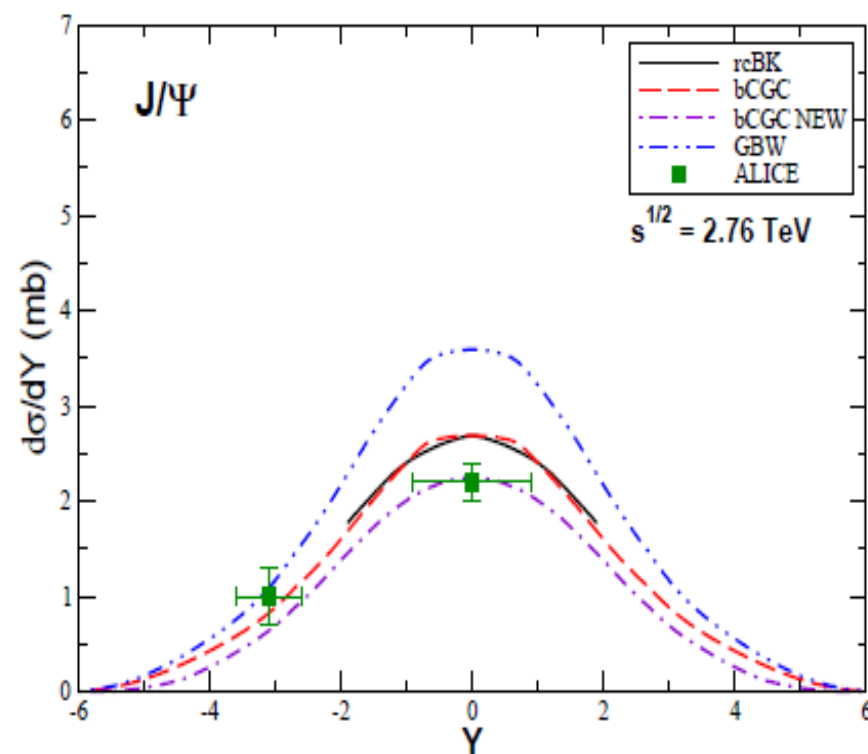
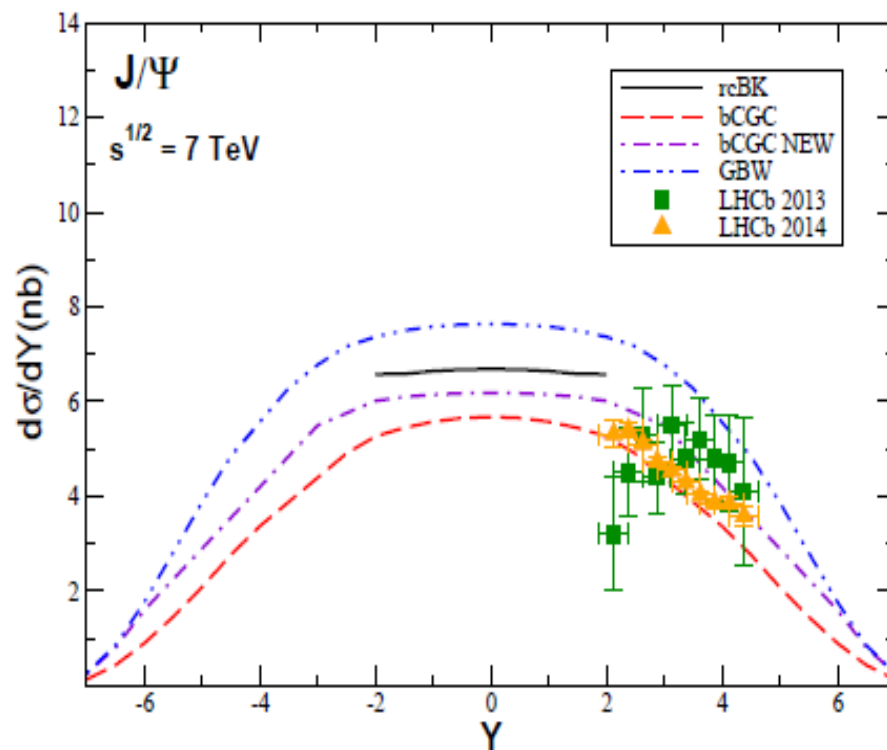


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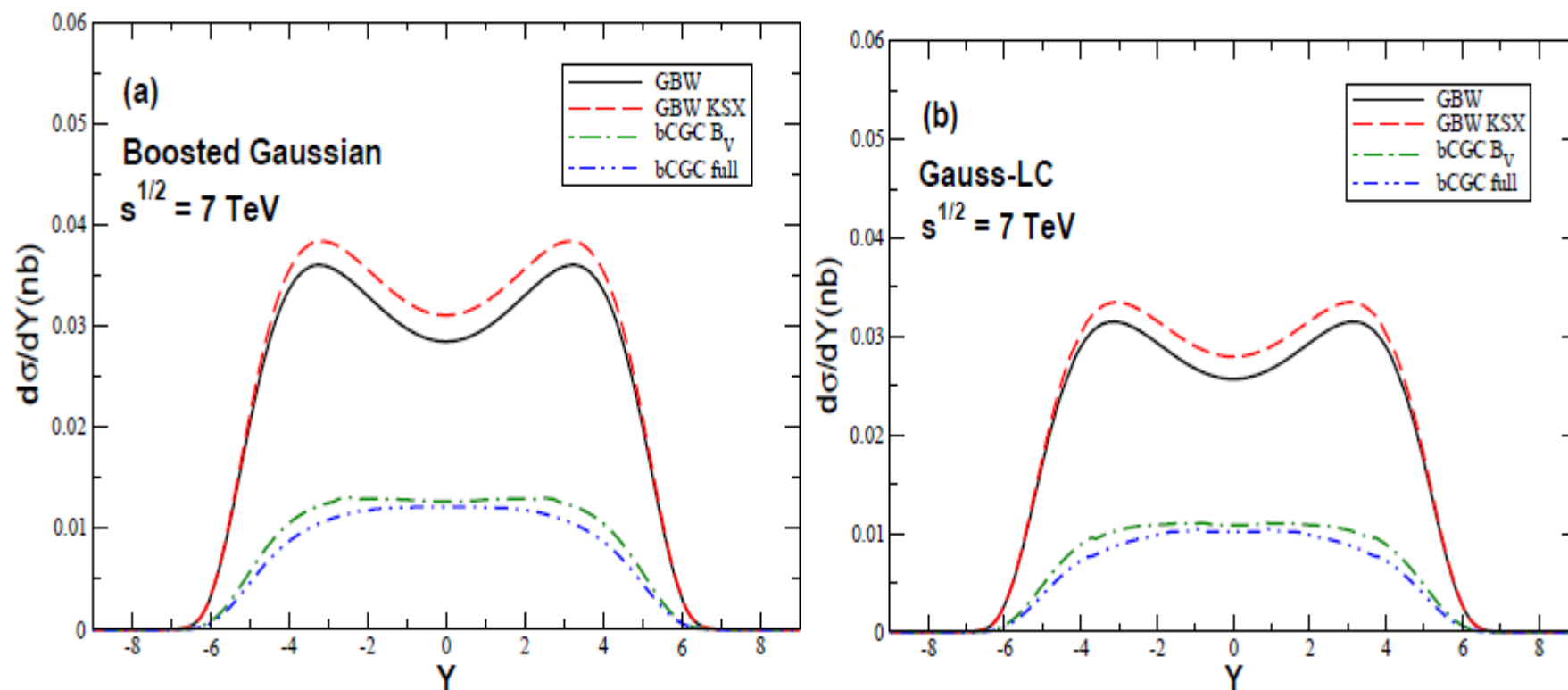
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<sup>*b*</sup>VPG, Moreira, Navarra, PLB 472, 172 (2015) )

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## Probing the QCD dynamics

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<sup>c</sup>VPG, Machado, PRC80, 054901 (2009); PRC84, 011902 (2011)

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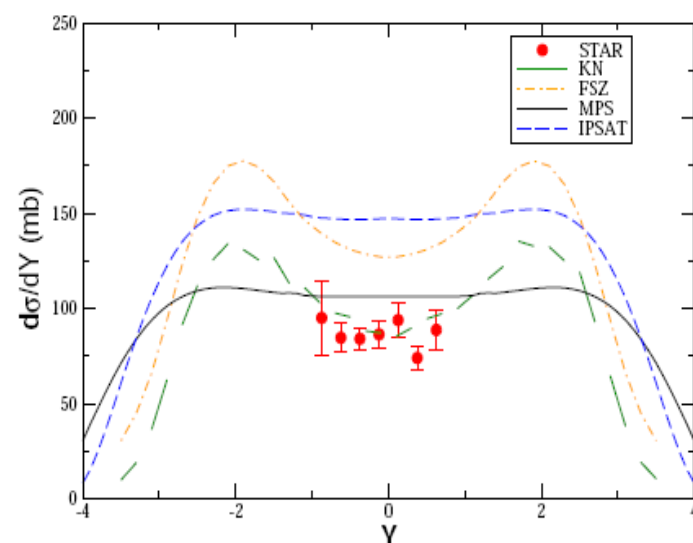


FIG. 1: (Color online) Predictions for the rapidity distribution of  $\rho^0$  photoproduction at RHIC energy considering distinct theoretical approaches. Data from STAR Collaboration [24].

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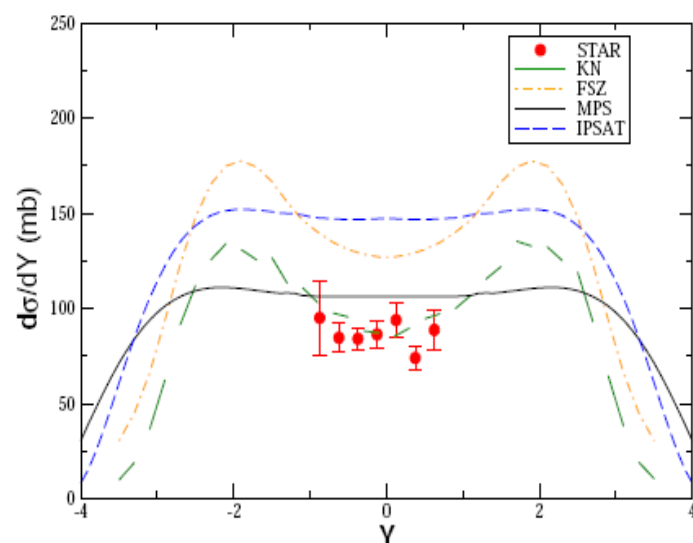
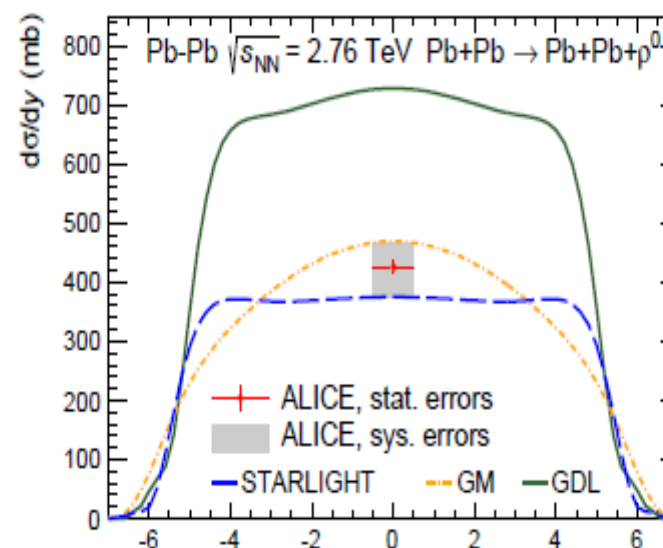


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arXiv:1503.09177v1 [nucl-ex]

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# **Photon – Hadron Interactions at the LHC:**

## **Probing the Odderon**

# Photon – Hadron Interactions at the LHC:

## Probing the Odderon

- In pQCD the **Pomeron** corresponds to a  $C$  - even parity ( $C$  being the charge conjugation) compound state of two  $t$ -channel reggeized gluons, which determines the high energy behaviour of the total hadronic cross sections.

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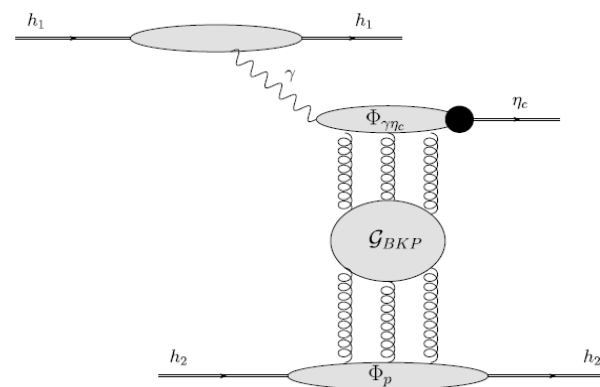
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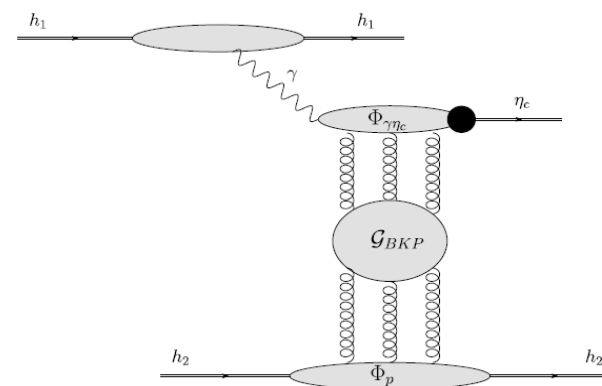
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- A natural prediction of the QCD is the presence of a  $C$  - odd compound state of three reggeized gluons, the so-called **Odderon**, which dominates the hadronic cross section difference between the direct and crossed channel processes at very high energies.
- Open question:** Does the Odderon exist?

- Alternative:** Consider exclusive processes in which the Odderon is the only contribution !
- Diffractive  $\eta_c$  photoproduction in hadronic collisions <sup>a</sup>.



- Basic idea:** As the photon carries negative  $C$  parity, its transformation into a diffractive final state system of positive  $C$  parity requires the  $t$  - channel exchange of an object of negative  $C$  parity.

<sup>a</sup>VPG, NPA 902, 32 (2013)

# Photon – Hadron Interactions at the LHC:

## Probing the Odderon

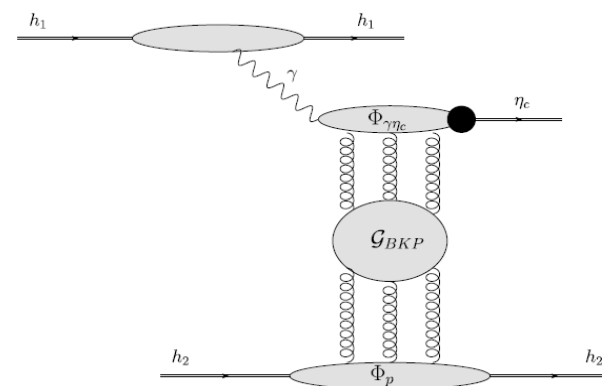
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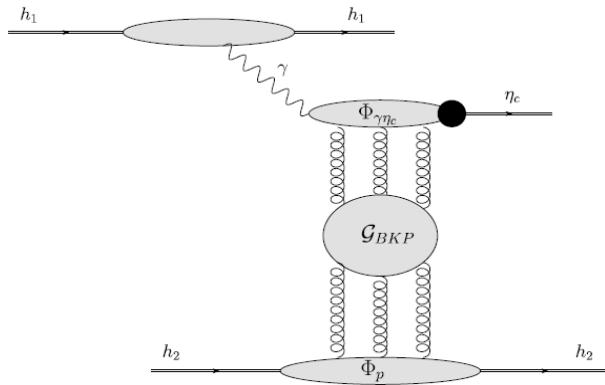
- Pomeron exchange cannot contribute to this process.

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## Probing the Odderon

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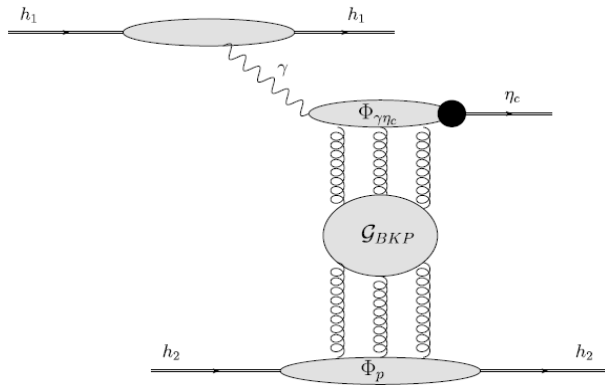
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### ● Predictions:

Table 1

Cross sections (event rates/year) for the diffractive  $\eta_c$  photoproduction in  $pp$  collisions at LHC energies.

$\sqrt{s_{NN}}$	CKMS	BBCV
8 TeV	0.55 pb (55 000)	10.10 pb ( $1 \times 10^6$ )
14 TeV	0.65 pb (65 000)	13.90 pb ( $1.4 \times 10^6$ )

Table 2

Cross sections (event rates/year) for the diffractive  $\eta_c$  photoproduction in PbPb collisions at LHC energies.

$\sqrt{s_{NN}}$	CKMS	BBCV
2.76 TeV	0.30 $\mu$ b (126)	14.25 $\mu$ b (5985)
5.5 TeV	0.40 $\mu$ b (168)	23.59 $\mu$ b (9912)

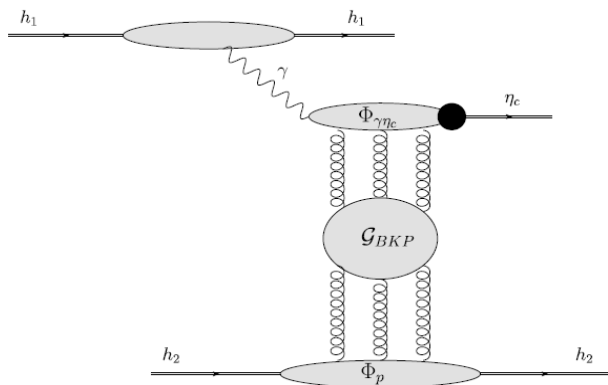
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- Predictions for the AFTER@LHC experiment <sup>b</sup>.

$h_1 h_2$	CKMS	BBCV
$pp$ ( $\sqrt{s} = 115$ GeV)	0.05 pb (1000.0)	0.30 pb (6000.0)
$Pbp$ ( $\sqrt{s} = 72$ GeV)	28.1 pb (31.0)	356.6 pb (393.0)
$PbPb$ ( $\sqrt{s} = 72$ GeV)	5870.0 pb (41.0)	74366.0 pb (520.0)

TABLE I: Cross sections (event rates/year) for the exclusive  $\eta_c$  photoproduction in  $pp/Pbp/PbPb$  collisions at AFTER@LHC experiment.

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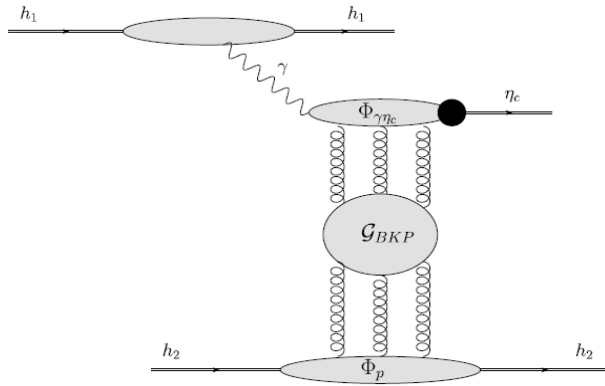
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### Background is only present in $pp$ collisions, which makes the observation of the exclusive $\eta_c$ production in $Pbp$ and $PbPb$ collisions a signature of the Odderon.

<sup>a</sup>VPG, NPA 902, 32 (2013)

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# Photon – Hadron Interactions at the LHC:

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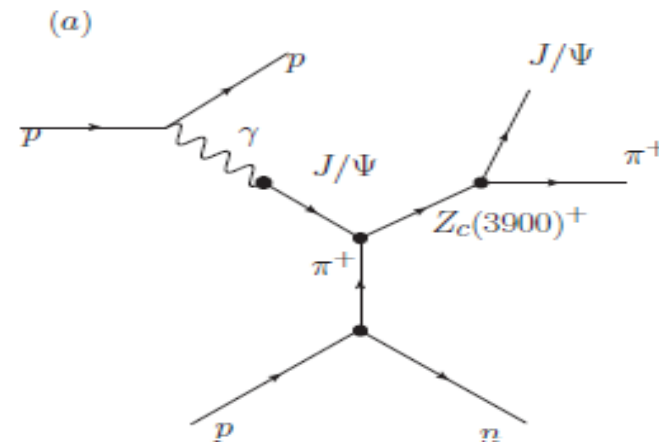
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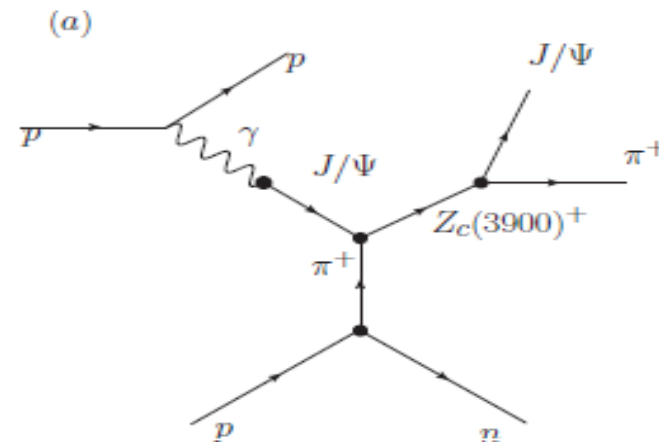


<sup>a</sup>VPG, Silva, PRD 89, 114005 (2014)

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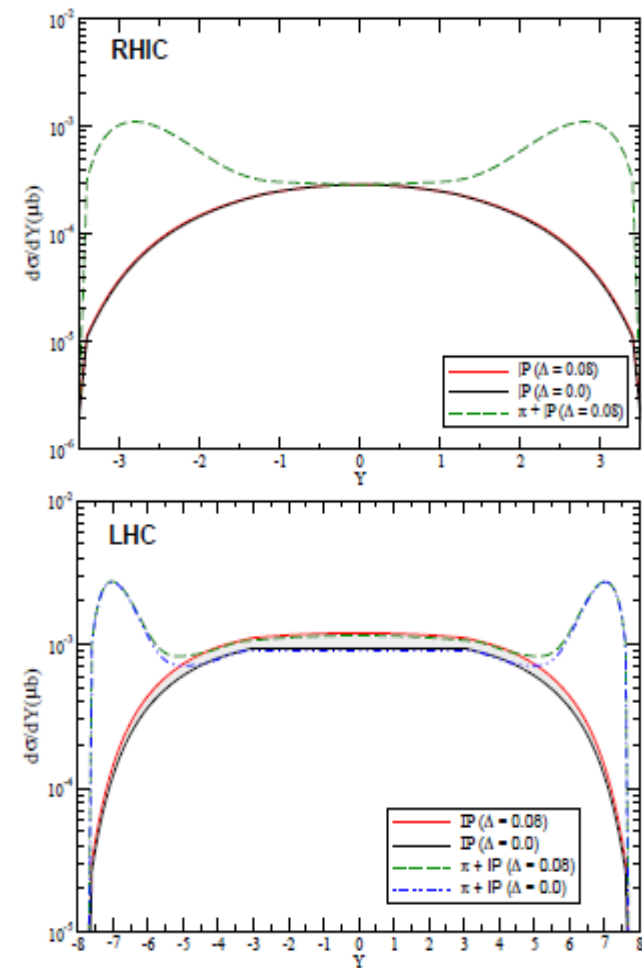
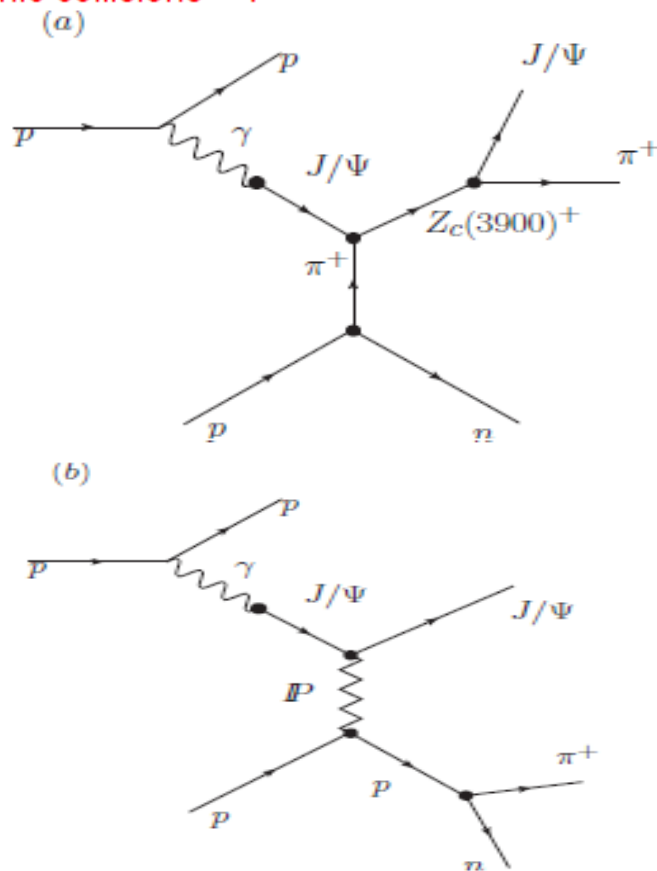


- **Basic idea:** the photon stemming from one of the incident protons fluctuates into a  $J/\Psi$  which interacts with the other proton through the  $\pi$  - exchange producing a  $n$  and a  $Z_c^+(3900)$  state which decays in the  $J/\Psi + \pi$  system.

<sup>a</sup>VPG, Silva, PRD 89, 114005 (2014)

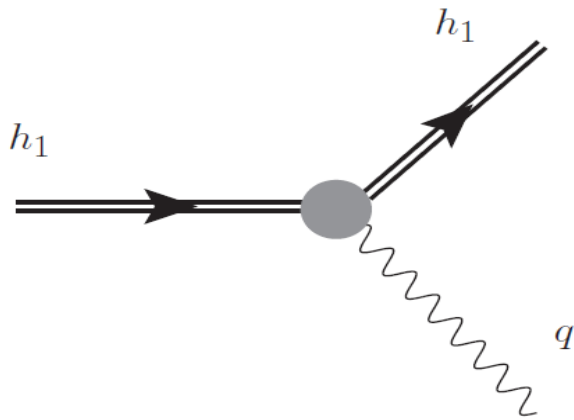
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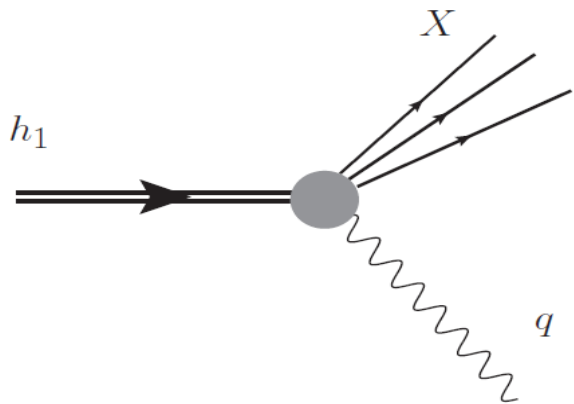
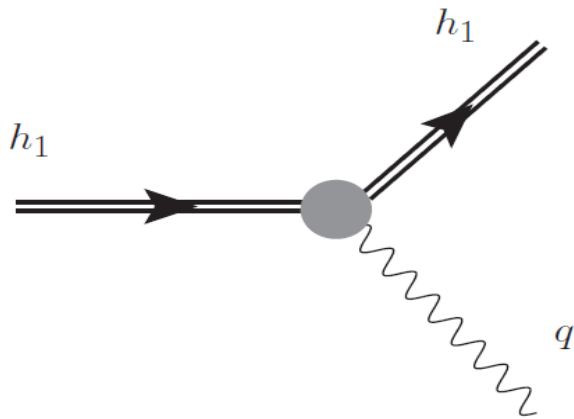
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# Photon – Hadron Interactions at the LHC: Probing the Photon Distribution of the Proton

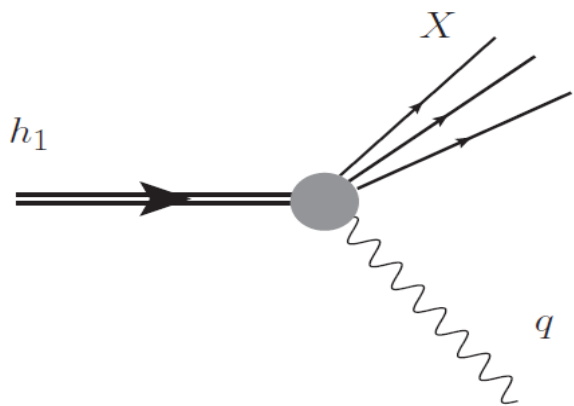
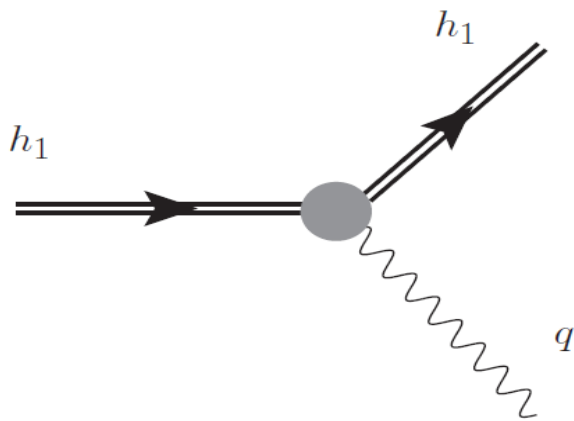




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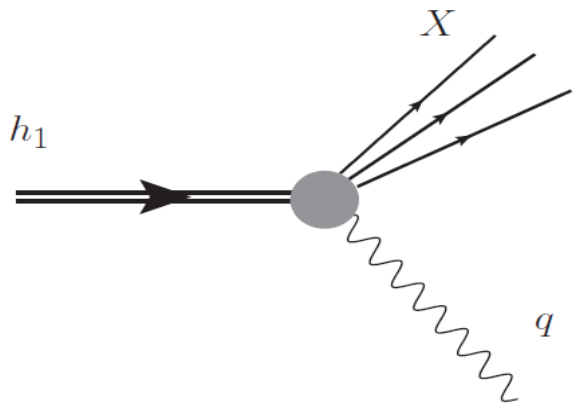
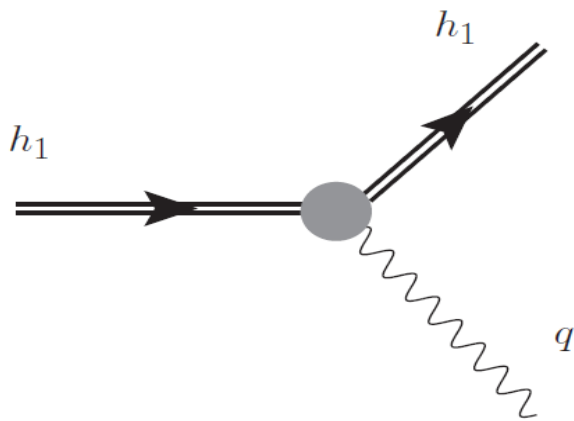
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$$\gamma(x, \mu^2) = \gamma_{\text{el}}(x) + \gamma_{\text{inel}}(x, \mu^2)$$

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- Several groups (MRST, NNPDF, CTEQ) have proposed to treat the photon as one of the point-like partons inside the nucleon and to account for this QED effect explicitly in the global analysis.

# Photon – Hadron Interactions at the LHC: Probing the Photon Distribution of the Proton

- The inclusion of the QED effects implies that the quark, gluon and photon distributions satisfy QED-modified DGLAP equations, which are given at leading order in both  $\alpha_s$  and  $\alpha$  by:

$$\begin{aligned}\frac{\partial q_i(x, \mu^2)}{\partial \log \mu^2} &= \frac{\alpha_s}{2\pi} \int_x^1 \frac{dy}{y} \left\{ P_{qq}(y) q_i\left(\frac{x}{y}, \mu^2\right) + P_{qg}(y) g\left(\frac{x}{y}, \mu^2\right) \right\} \\ &+ \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} \left\{ \tilde{P}_{qq}(y) e_i^2 q_i\left(\frac{x}{y}, \mu^2\right) + P_{q\gamma}(y) e_i^2 \gamma\left(\frac{x}{y}, \mu^2\right) \right\} \\ \frac{\partial g(x, \mu^2)}{\partial \log \mu^2} &= \frac{\alpha_s}{2\pi} \int_x^1 \frac{dy}{y} \left\{ P_{gq}(y) \sum_j q_j\left(\frac{x}{y}, \mu^2\right) + P_{gg}(y) g\left(\frac{x}{y}, \mu^2\right) \right\} \\ \frac{\partial \gamma(x, \mu^2)}{\partial \log \mu^2} &= \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} \left\{ P_{\gamma q}(y) \sum_j e_j^2 q_j\left(\frac{x}{y}, \mu^2\right) + P_{\gamma\gamma}(y) \gamma\left(\frac{x}{y}, \mu^2\right) \right\},\end{aligned}$$

where

$$\begin{aligned}\tilde{P}_{qq} &= C_F^{-1} P_{qq}, & P_{\gamma q} &= C_F^{-1} P_{gq}, \\ P_{q\gamma} &= T_R^{-1} P_{qg}, & P_{\gamma\gamma} &= -\frac{2}{3} \sum_i e_i^2 \delta(1-y)\end{aligned}$$

and momentum is conserved:

$$\int_0^1 dx \, x \left\{ \sum_i q_i(x, \mu^2) + g(x, \mu^2) + \gamma(x, \mu^2) \right\} = 1.$$

# Photon – Hadron Interactions at the LHC:

## Probing the Photon Distribution of the Proton



Initial condition  $\gamma(x, Q_0^2)$ :

MRST: Naive model

NNPDF: Freely parametrized

CTEQ: Similar to that proposed by MRST, but with arbitrary normalization.

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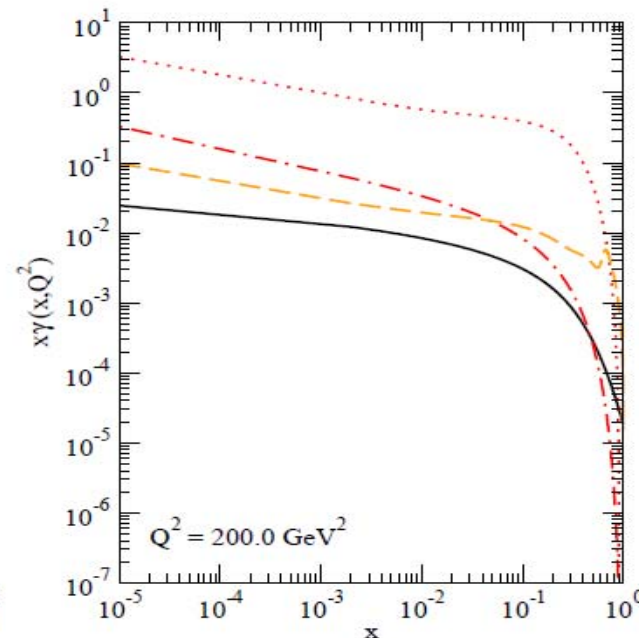
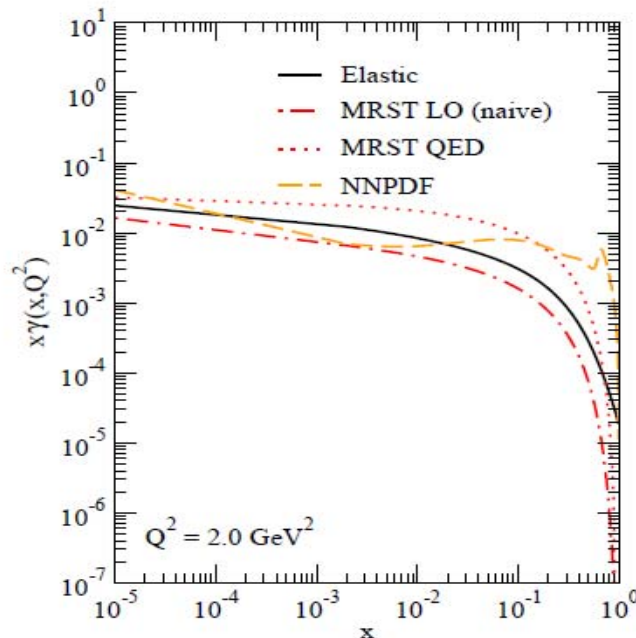


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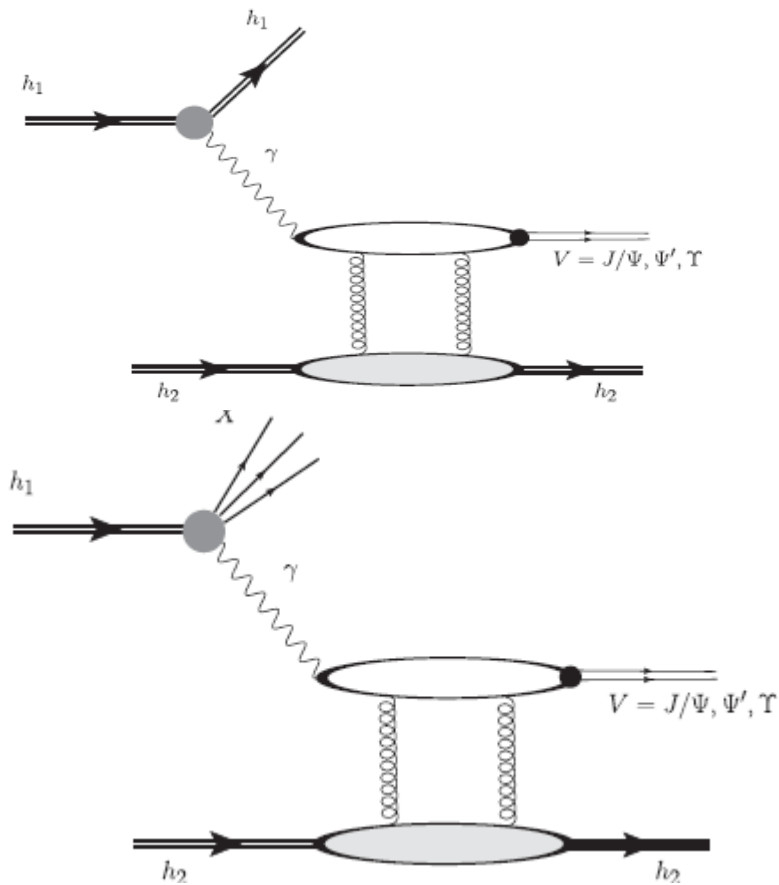
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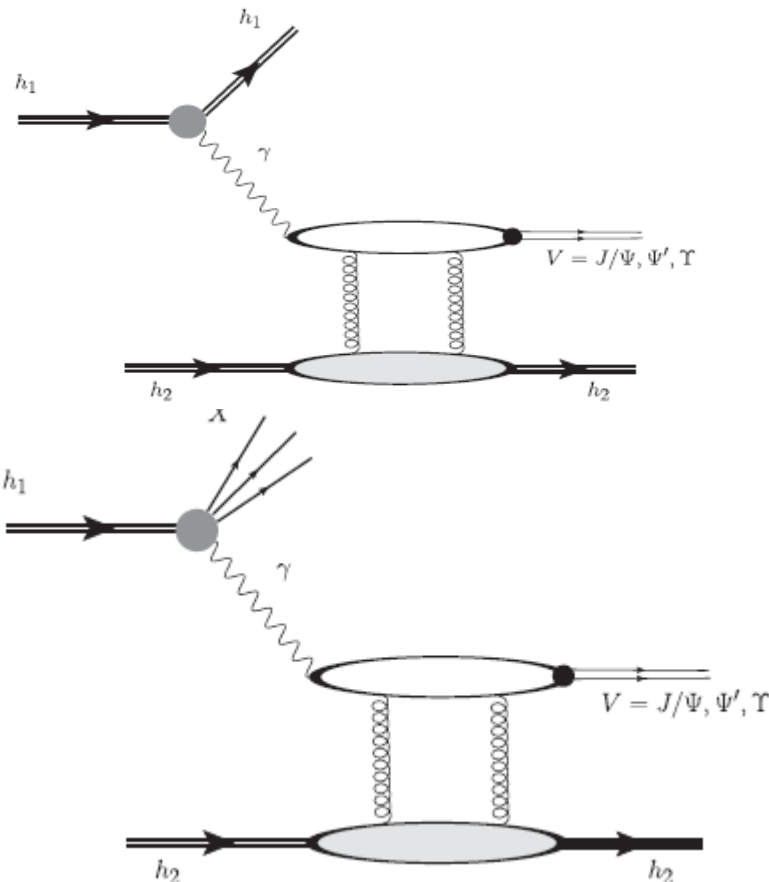
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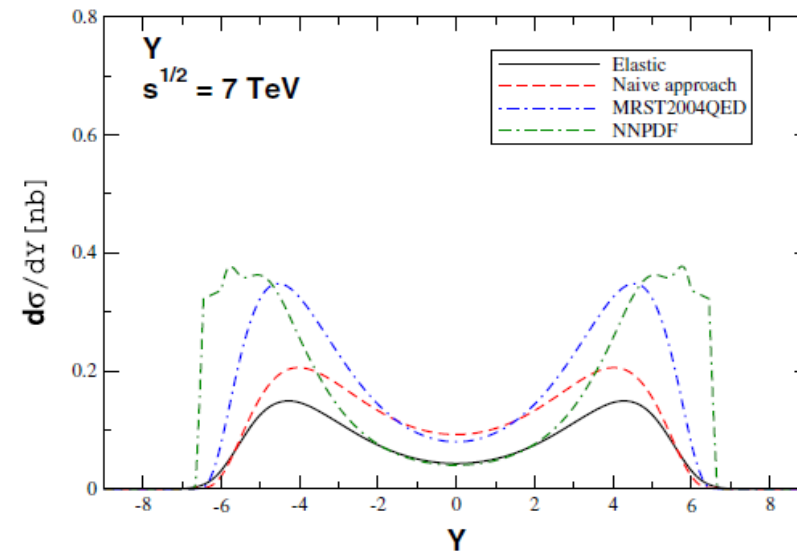
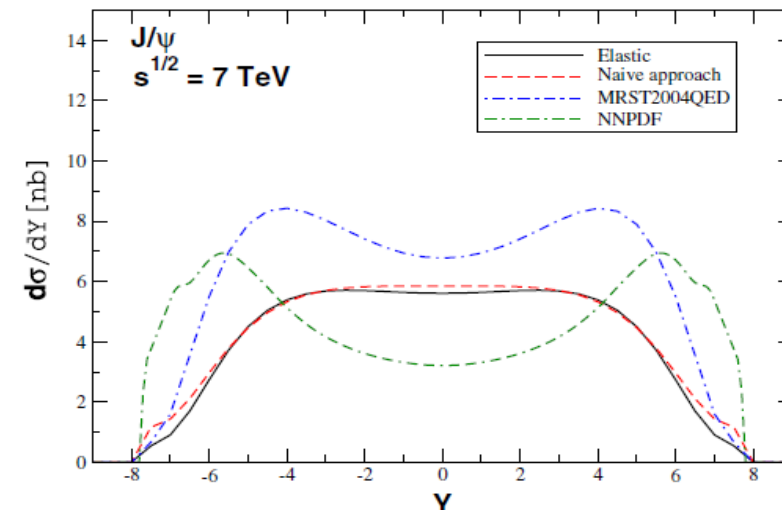
<sup>a</sup>VPG, da Silveira, PRD 91, 054013 (2015)

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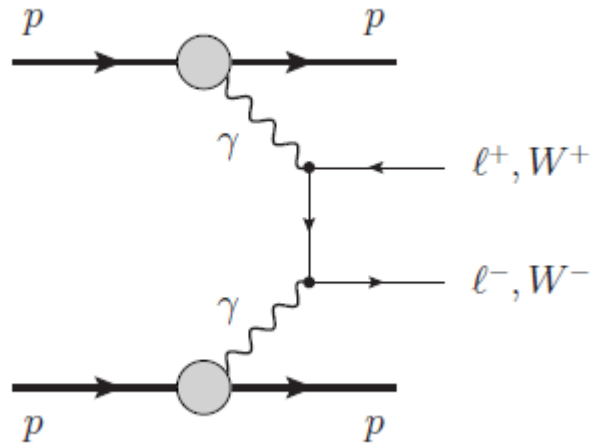


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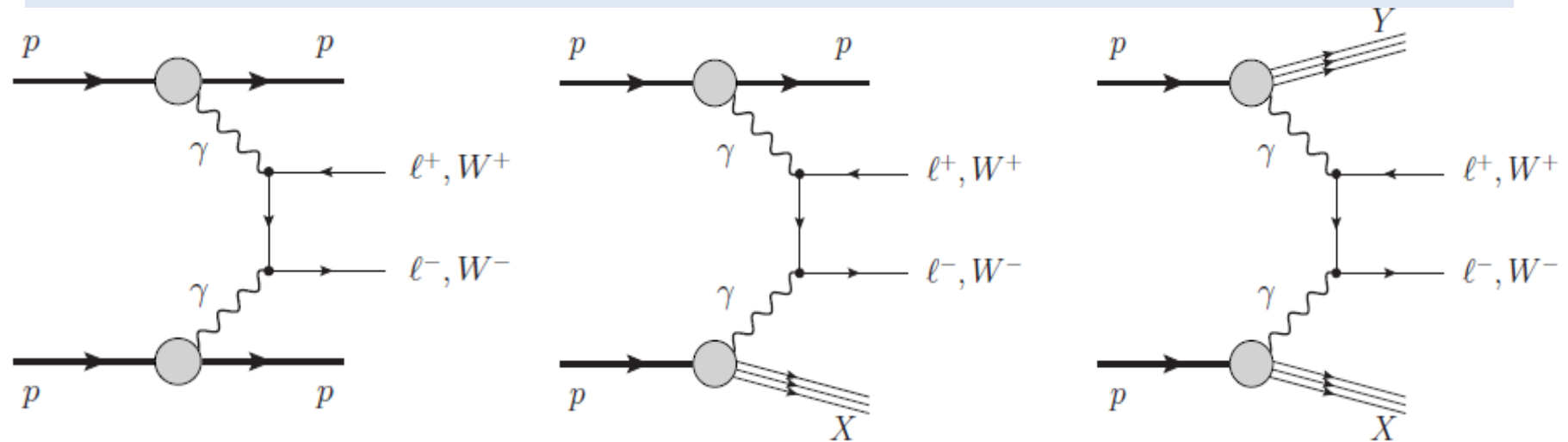




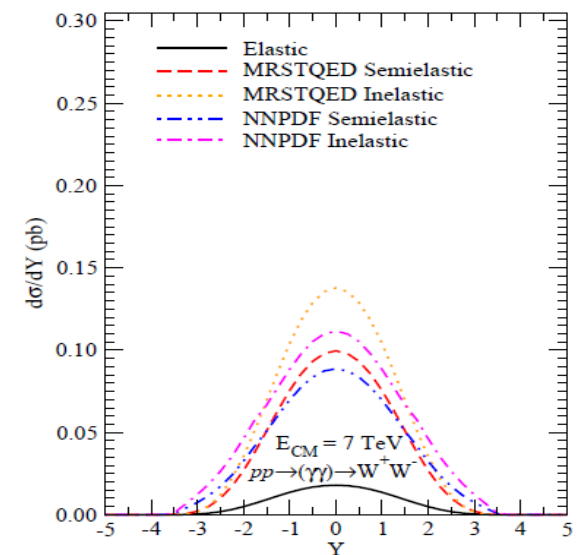
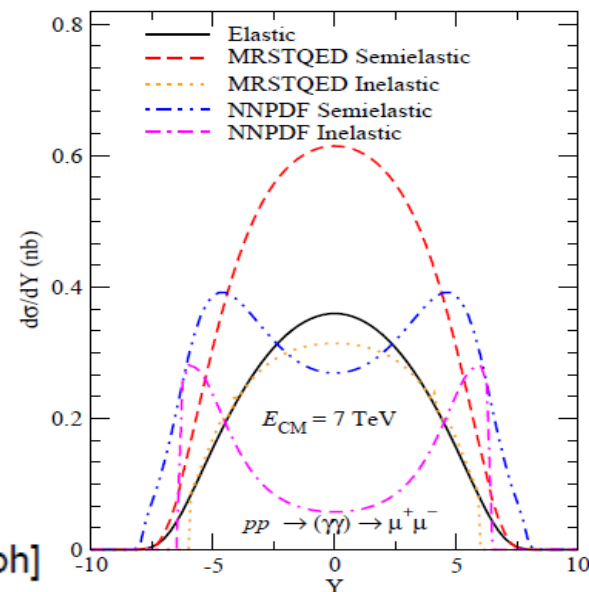
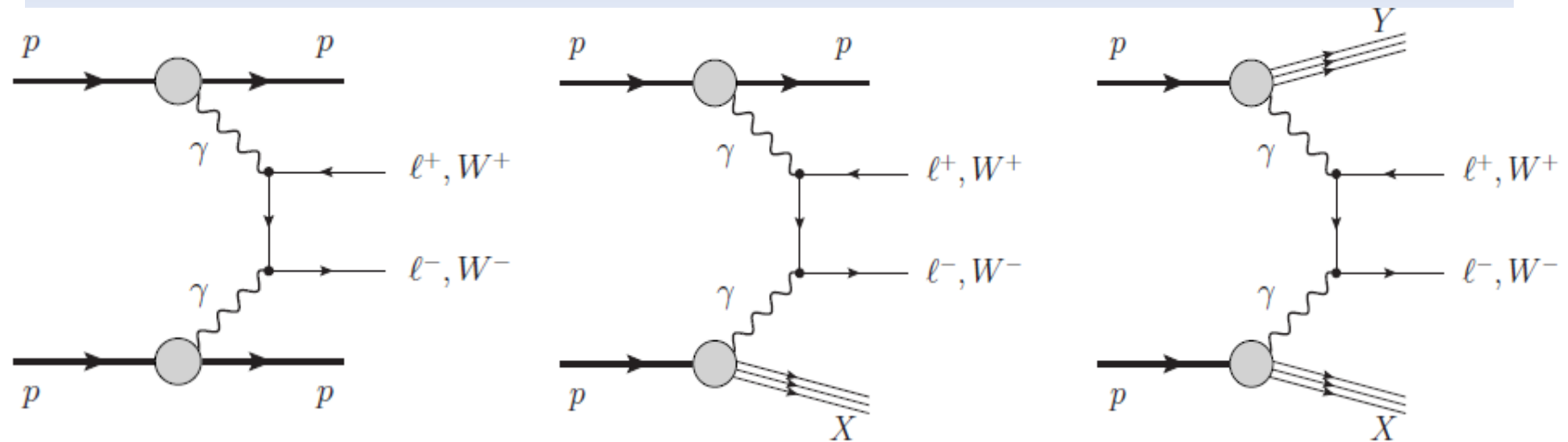
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**Thank you for your attention !**

**Extras**



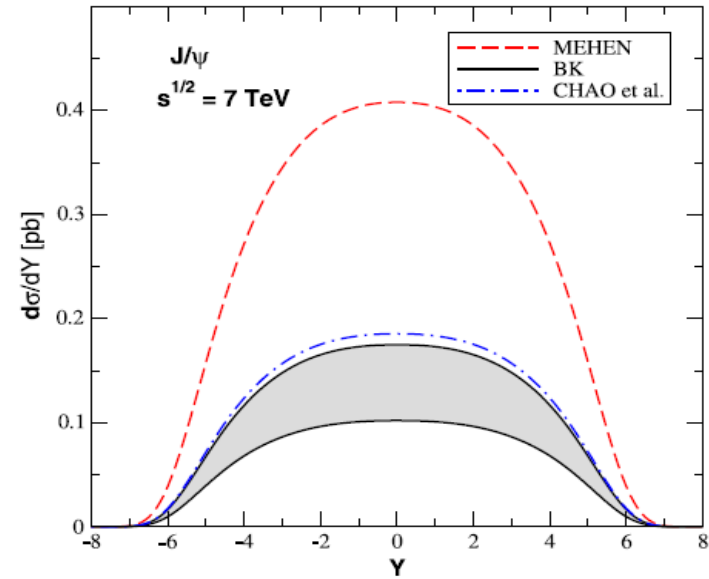
# **Probing Quarkonium Production in photon – induced interactions**

# Quarkonium + Photon production in photon – induced interactions (\*)

$$\begin{aligned} & \frac{d\sigma[p + p \rightarrow p \otimes H + \gamma + X]}{dY} \\ &= \omega \frac{dN_{\gamma/h_1}(\omega)}{d\omega} \sigma_{\gamma h_2 \rightarrow H + \gamma + X}(\omega) \\ &+ \omega \frac{dN_{\gamma/h_2}(\omega)}{d\omega} \sigma_{\gamma h_1 \rightarrow H + \gamma + X}(\omega), \end{aligned}$$

$$\begin{aligned} & \sigma(\gamma + p \rightarrow H + \gamma + X) \\ &= \int dz dp_{\perp}^2 \frac{xg(x, Q^2)}{z(1-z)} \frac{d\sigma}{dt}(\gamma + g \rightarrow H + \gamma) \end{aligned}$$

$$\begin{aligned} & \frac{d\sigma}{dt}(\gamma + g \rightarrow H + \gamma) \\ &= \frac{64\pi^2}{3} \frac{e_Q^4 \alpha^2 \alpha_s m_Q}{s^2} \left( \frac{s^2 s_1^2 + t^2 t_1^2 + u^2 u_1^2}{s_1^2 t_1^2 u_1^2} \right) \langle O_8^V(^3S_1) \rangle \end{aligned}$$



**Table 1** The total cross section for the  $H + \gamma$  photoproduction in coherent hadron–hadrons collisions at LHC energies

$J/\psi + \gamma$	MEHEN	BK
LHC (7 TeV)	3.62 pb	$1.23 \pm 0.50$ pb
LHC (14 TeV)	5.60 pb	$1.90 \pm 0.32$ pb
$\Upsilon + \gamma$	BFL	BSV
LHC (14 TeV)	5.46 fb	$1.45 \pm 0.13$ fb

(\*) VPG, M. M. Machado, EPJC 72, 2231 (2012)

# Inelastic Quarkonium production in photon – induced interactions (\*)

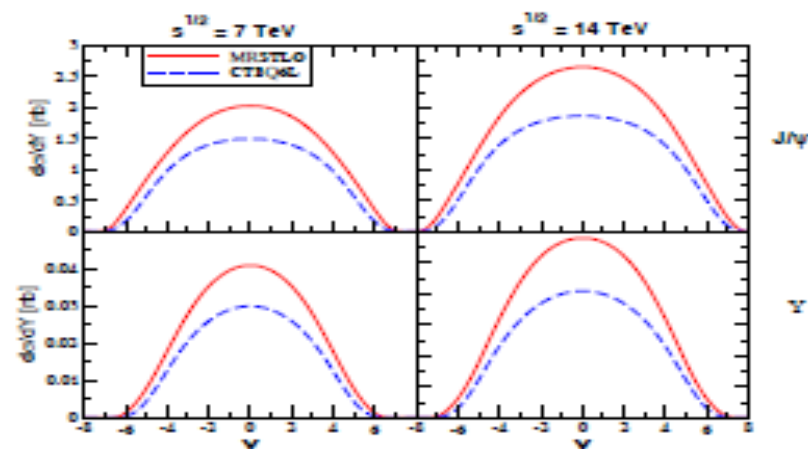
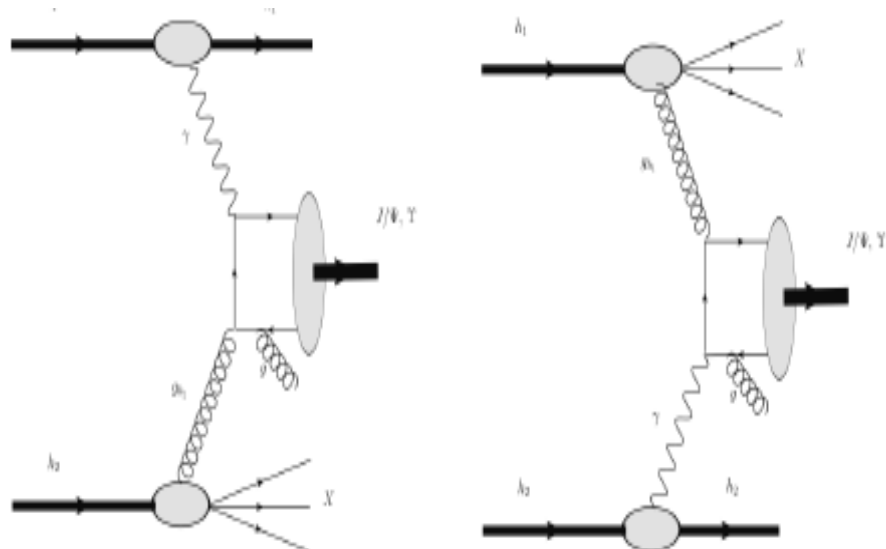


Fig. 4. Rapidity distribution for the  $J/\psi$  and  $\gamma$  production in coherent  $pp$  collisions at  $\sqrt{s} = 7$  TeV (left panels) and 14 TeV (right panels) considering two different parametrizations for

Table 1. The total cross section (event rates) for the inelastic quarkonium photoproduction in coherent  $pp$  collisions at LHC energies.

$J/\psi$	MRSTLO	CTEQ6L
$\sqrt{s} = 7$ TeV	18.0 nb ( $1.8 \times 10^9$ )	13.0 nb ( $1.3 \times 10^9$ )
$\sqrt{s} = 14$ TeV	25.0 nb ( $2.5 \times 10^9$ )	18.0 nb ( $1.8 \times 10^9$ )
$\gamma$	MRSTLO	CTEQ6L
$\sqrt{s} = 7$ TeV	0.30 nb ( $30 \times 10^6$ )	0.21 nb ( $21 \times 10^6$ )
$\sqrt{s} = 14$ TeV	0.47 nb ( $47 \times 10^6$ )	0.33 nb ( $33 \times 10^6$ )

(\*) VPG, M. M. Machado, EPJA 50, 72 (2014)