

Pole position of the $a_1(1260)$

Misha Mikhasenko

Joint Physics Analysis Center,
COMPASS @ CERN,
Universität Bonn, HISKP, Bonn, Germany

CHARM 2018
Akademgorodok, Novosibirsk

23/05/2018



Bundesministerium
für Bildung
und Forschung

Overview

1

Introduction

- Hadrons in QCD
- Analytical structure of the scattering amplitude

2

Three pions dynamics

- Constrains
- Data

3

Extraction of the resonance parameters

- Fit
- Analytical continuation

4

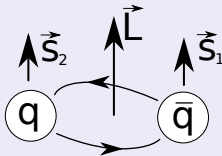
Remarks

- COMPASS analysis
- CLEO analysis
- $a_1(1420)$ phenomenon

Flavor and excitation

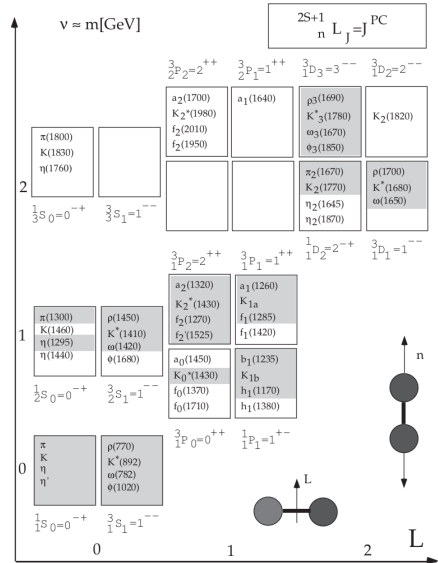
Quark model

- color-binding,



- Radial excitation (n),
- Orbital excitation (L),

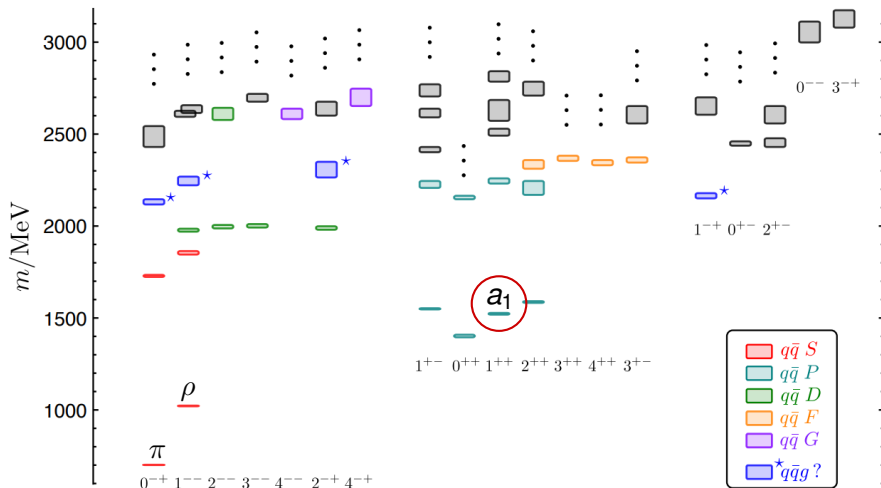
- many states ($J^{PC} = 0^{++}, 1^{--} \dots$) are coupled to $\pi\pi$.
- Some other to 3π system



[Amsler et al., Phys. Rept. 389, 61 (2004)]

Lattice QCD

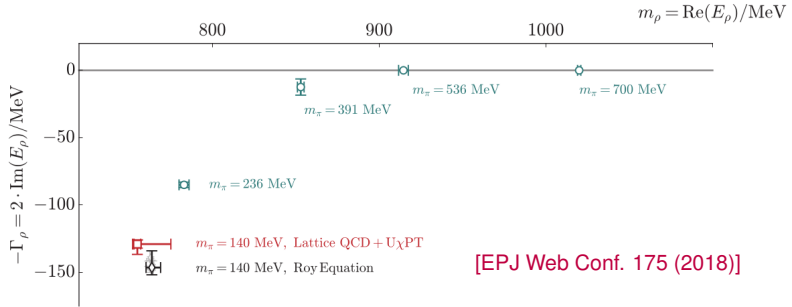
Lattice QCD spectrum matches experimental observations well but predicts more



[Dudek et. al, Phys.Rev. D82 (2010) 034508]

Resonances on the Lattice

$\pi\pi$ system in the box



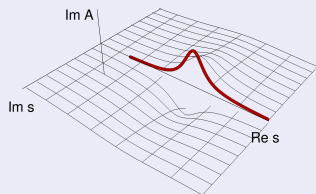
Tracking pole position

- For high m_π the ρ -becomes stable. Pole of ρ approaches real axis.

[Wilson, D. et al., PRD 92, (2015)]

Resonances = Poles at the Complex plane

Breit-Wigner amplitude

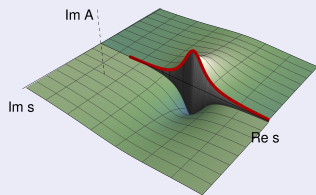


Features of the complex s plane:

- $s = E^2$ – the total inv.mass squared
- The Real axis $\rightarrow$ physical world
- The Imaginary axis $\rightarrow$ analytical continuation

Resonances = Poles at the Complex plane

Breit-Wigner amplitude

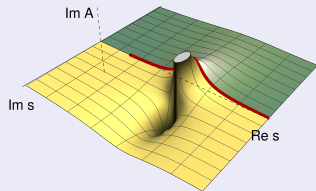


Features of the complex s plane:

- $s = E^2$ – the total inv.mass squared
- The Real axis $\rightarrow$ physical world
- The Imaginary axis $\rightarrow$ analytical continuation

Resonances = Poles at the Complex plane

Breit-Wigner amplitude



Features of the complex s plane:

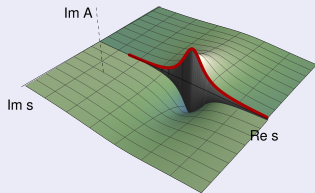
- $s = E^2$ – the total inv.mass squared
- The Real axis $\rightarrow$ physical world
- The Imaginary axis $\rightarrow$ analytical continuation

Unitarity constraints for **two-body** scattering

- $\hat{S}^\dagger \hat{S} = \hat{\mathbb{I}} \quad \hat{S} = \hat{\mathbb{I}} + i\hat{T} \quad \hat{T} - \hat{T}^\dagger = i\hat{T}^\dagger \hat{T}.$
- $T(s, t) = \langle p'_1 p'_2 | \hat{T} | p_1 p_2 \rangle \quad 2\text{Im} T(s, t) = \int d\Phi_2 T^*(s, t') T(s, t'')$
- Partial wave expansion $\rightarrow$
- The final form $2\text{Im} t_l(s) = t_l^*(s) \rho(s) t_l(s)$

Resonances = Poles at the Complex plane

Breit-Wigner amplitude



Features of the complex s plane:

- $s = E^2$ – the total inv.mass squared
- The Real axis $\rightarrow$ physical world
- The Imaginary axis $\rightarrow$ analytical continuation

Unitarity constraints for **two-body** scattering

- $\hat{S}^\dagger \hat{S} = \hat{\mathbb{I}} \quad \hat{S} = \hat{\mathbb{I}} + i\hat{T} \quad \hat{T} - \hat{T}^\dagger = i\hat{T}^\dagger \hat{T}.$
- $T(s, t) = \langle p'_1 p'_2 | \hat{T} | p_1 p_2 \rangle \quad 2\text{Im} T(s, t) = \int d\Phi_2 T^*(s, t') T(s, t'')$
- Partial wave expansion $\rightarrow$
- The final form $2\text{Im} t_l(s) = t_l^*(s) \rho(s) t_l(s)$

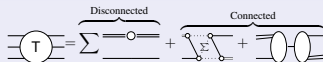
Quasi-two-body unitarity

[MM (JPAC) in preparation]

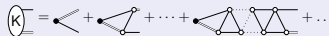
Three-body unitarity

[Eden, Landshoff et al.(2002)]

- Singularity splitting:



- Final state interaction:



$$s \left\{ \sigma \left\{ \begin{array}{c} \xi \\ \text{---} \end{array} \right. \left(\text{shaded circle} \right) \left. \begin{array}{c} \xi \\ \text{---} \end{array} \right\} \sigma' \right.$$

Quasi-two-body unitarity

[MM (JPAC) in preparation]

Three-body unitarity

[Eden, Landshoff et al.(2002)]

- Singularity splitting:

$$T = \sum_{\text{Disconnected}} + \sum_{\text{Connected}} \quad s \left\{ \sigma \left\{ \begin{array}{c} \xi \\ \text{---} \end{array} \right. \right. \left. \left. \begin{array}{c} \xi \\ \text{---} \end{array} \right\} \sigma' \right.$$

- Final state interaction:

$$K = \text{---} + \text{---} + \dots + \text{---} + \dots$$

$$T(\sigma', s, \sigma) = K_{\xi}(s, \sigma') t(s) K_{\xi}(s, \sigma)$$

Quasi-two-body unitarity

[MM (JPAC) in preparation]

Three-body unitarity

[Eden, Landshoff et al.(2002)]

- Singularity splitting:

$$T = \underbrace{\sum \text{Disconnected}}_{\text{Disconnected}} + \underbrace{\sum \text{Connected}}_{\text{Connected}} \quad s \left\{ \sigma \left\{ \begin{array}{c} \xi \\ \text{---} \end{array} \right. \right. \left. \left. \begin{array}{c} \xi \\ \text{---} \end{array} \right\} \sigma' \right.$$

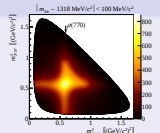
- Final state interaction:

$$K = \text{---} + \text{---} + \dots + \text{---} + \dots$$

$$T(\sigma', s, \sigma) = K_\xi(s, \sigma') t(s) K_\xi(s, \sigma) \quad 2\text{Im } t(s) = t^*(s) \rho(s) t(s),$$

Symmetrized quasi-two-body phase space factor

$$\rho(s) = \frac{1}{2} \int d\Phi_3 \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} - \begin{array}{c} \text{---} \\ \text{---} \end{array} \right|^2 = \int d\Phi_3 \left[\begin{array}{c} \text{---} \\ \text{---} \end{array} - \underbrace{\begin{array}{c} \text{---} \\ \text{---} \end{array}}_{\text{interference}} \right]$$



Quasi-two-body unitarity

[MM (JPAC) in preparation]

Three-body unitarity

[Eden, Landshoff et al.(2002)]

- Singularity splitting:

$$T = \underbrace{\sum \text{Disconnected}}_{\text{Disconnected}} + \underbrace{\sum \text{Connected}}_{\text{Connected}} \quad s \left\{ \sigma \left\{ \begin{array}{c} \xi \\ \text{---} \end{array} \right. \right. \left. \left. \begin{array}{c} \xi \\ \text{---} \end{array} \right\} \sigma' \right.$$

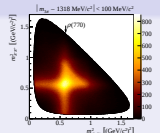
- Final state interaction:

$$K = \text{---} + \text{---} + \dots + \text{---} + \dots$$

$$T(\sigma', s, \sigma) = K_{\xi}(s, \sigma') t(s) K_{\xi}(s, \sigma) \quad 2\text{Im } t(s) = t^*(s) \rho(s) t(s),$$

Symmetrized quasi-two-body phase space factor

$$\rho(s) = \frac{1}{2} \int d\Phi_3 \left| \text{---} - \text{---} \right|^2 = \int d\Phi_3 \left[\text{---} - \underbrace{\text{---}}_{\text{interference}} \right]$$



The model: symmetrized

[Bowler, Phys.Lett.B182 (1986)]

$$t(s) = \frac{g^2}{m^2 - s - ig^2 \rho(s)/2}$$

Quasi-two-body unitarity

[MM (JPAC) in preparation]

Three-body unitarity

[Eden, Landshoff et al.(2002)]

- Singularity splitting:

$$T = \underbrace{\sum \text{Disconnected}}_{\text{Disconnected}} + \underbrace{\sum \text{Connected}}_{\text{Connected}} \quad s \left\{ \sigma \left\{ \begin{array}{c} \xi \\ \text{---} \end{array} \right. \right. \left. \left. \begin{array}{c} \xi \\ \text{---} \end{array} \right\} \sigma' \right.$$

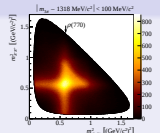
- Final state interaction:

$$K = \text{---} + \text{---} + \dots + \text{---} + \dots$$

$$T(\sigma', s, \sigma) = K_{\xi}(s, \sigma') t(s) K_{\xi}(s, \sigma) \quad 2\text{Im } t(s) = t^*(s) \rho(s) t(s),$$

Symmetrized quasi-two-body phase space factor

$$\rho(s) = \frac{1}{2} \int d\Phi_3 \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} - \begin{array}{c} \text{---} \\ \text{---} \end{array} \right|^2 = \int d\Phi_3 \left[\begin{array}{c} \text{---} \\ \text{---} \end{array} - \underbrace{\begin{array}{c} \text{---} \\ \text{---} \end{array}}_{\text{interference}} \right]$$



The model: symmetrized, dispersive

$$t(s) = \frac{g^2}{m^2 - s - ig^2 \tilde{\rho}(s)/2}, \quad \tilde{\rho}(s) = \frac{s}{\pi i} \int_{9m_{\pi}^2}^{\infty} \frac{\rho(s')}{s'(s' - s)} ds', \quad \text{Im } i\tilde{\rho} = i\rho.$$

1^{++} light meson spectrum

Axial vector states below 2 GeV

- Dominated by 3π scattering

- ▶ $\rho\pi \sim 60\% - 80\%$
- ▶ $\sigma\pi \sim 5\% - 10\%$
- ▶ $f_2\pi \sim < 5\%$

- $K\bar{K}\pi < 3\%$

1^{++} light meson spectrum

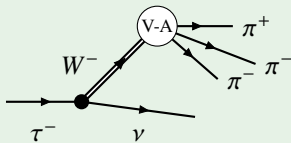
Axial vector states below 2 GeV

- Dominated by 3π scattering

- ▶ $\rho\pi \sim 60\% - 80\%$
- ▶ $\sigma\pi \sim 5\% - 10\%$
- ▶ $f_2\pi \sim < 5\%$

- $K\bar{K}\pi < 3\%$

$$\tau^- \rightarrow \pi^- \pi^+ \pi^- \nu$$



- V-A: Vector (1^{--}) or Axial (1^{++})

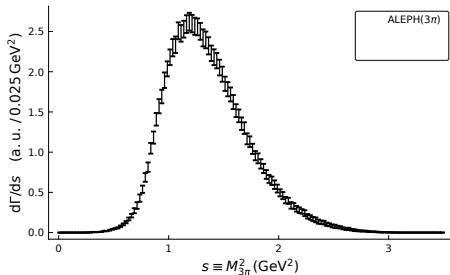
- Isospin 1 due to the charge

- Negative G-parity $\Rightarrow$ positive C-parity

$$\Rightarrow J^{PC} = 1^{++}$$

Fit to ALEPH data

[data from ALEPH, Phys.Rept.421 (2005)]

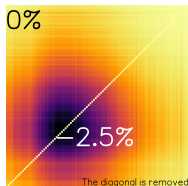


χ^2 function

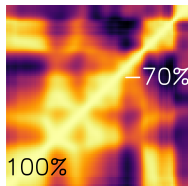
$$\chi^2(c, m, g) =$$

$$(\vec{D} - \vec{M}(c, m, g))^T C_{\text{stat}}^{-1} (\vec{D} - \vec{M}(c, m, g)),$$

- Stat. errors $\sim \times 5$ Systematic errors
- Stat. cov. matrix is used in the fit
- Syst. cov. matrix – in the bootstrap



Stat. cov. matrix



Syst. cov. matrix

Fit to ALEPH data

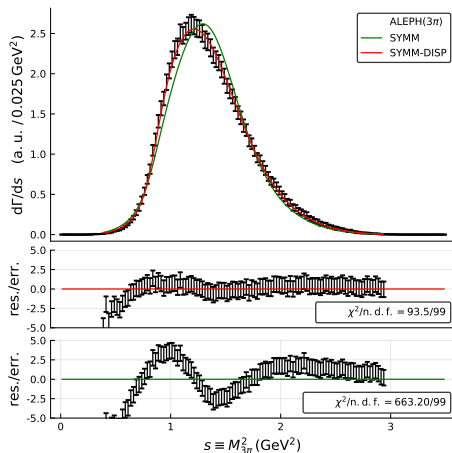
[data from ALEPH, Phys.Rept.421 (2005)]

χ^2 function

$$\chi^2(c, m, g) =$$

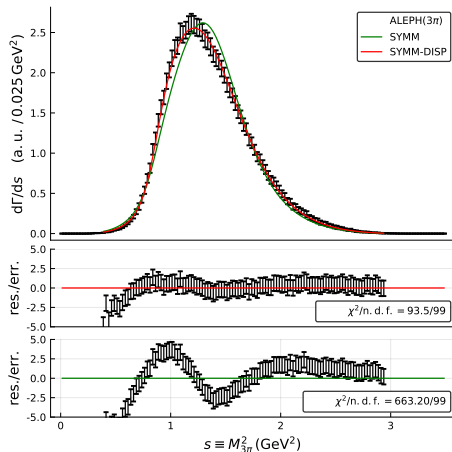
$$(\vec{D} - \vec{M}(c, m, g))^T C_{\text{stat}}^{-1} (\vec{D} - \vec{M}(c, m, g)),$$

- Stat. errors $\sim \times 5$ Systematic errors
- Stat. cov. matrix is used in the fit
- Syst. cov. matrix – in the bootstrap



Fit to ALEPH data

[data from ALEPH, Phys.Rept.421 (2005)]



χ^2 function

$$\chi^2(c, m, g) =$$

$$(\vec{D} - \vec{M}(c, m, g))^T C_{\text{stat}}^{-1} (\vec{D} - \vec{M}(c, m, g)),$$

- Stat. errors $\sim \times 5$ Systematic errors
- Stat. cov. matrix is used in the fit
- Syst. cov. matrix – in the bootstrap

Breit-Wigner mass and width

- The BW mass is given by $\text{Re}[t^{-1}(m^2)] = 0$
- The BW width $\Leftarrow \text{Im}[g^2/t(m^2)] = m\Gamma_{\text{BW}}$

Preliminary results:

$$m_{\text{BW}}^{a_1(1260)} = (1.246 \pm 0.003) \text{ GeV},$$

$$\Gamma_{\text{BW}}^{a_1(1260)} = (0.394 \pm 0.005) \text{ GeV}.$$

Tour to the complex plane

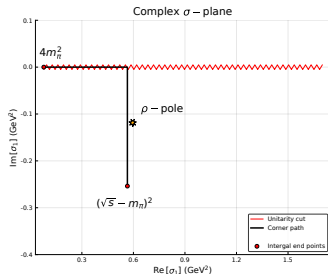
Analytical continuation

$$|t_I^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

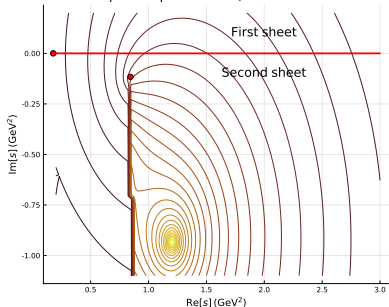
- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.



Tour to the complex plane

Complex s - plane in the QTB-DISP - model

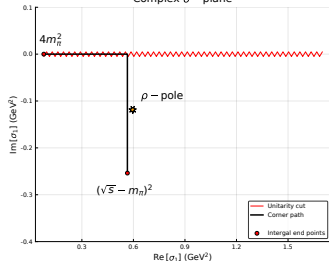
Analytical continuation

$$|t_l^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

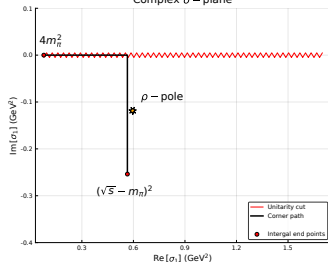
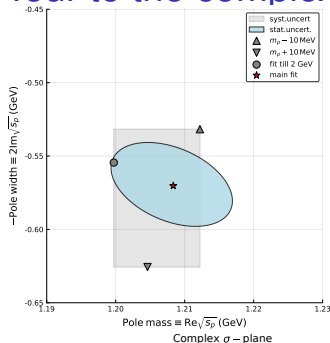
- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Complex σ - plane

Tour to the complex plane



Analytical continuation

$$|t_I^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Prel. results for the pole position

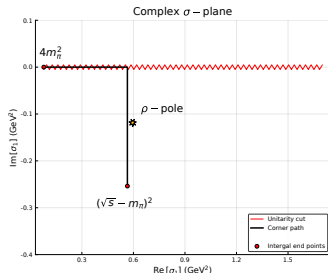
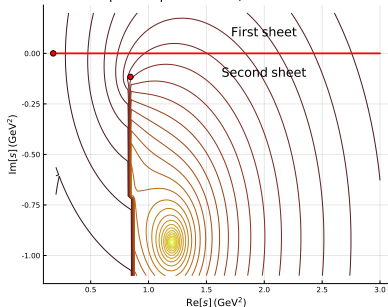
$$m_{\rho^{(a_1(1260))}} = (1.208 \pm 0.004) \text{ GeV},$$

$$\Gamma_{\rho^{(a_1(1260))}} = (0.569 \pm 0.012) \text{ GeV}.$$

Systematic studies:

- variation of ρ -shape
- variation of fit range

Tour to the complex plane

Complex s - plane in the QTB-DISP - model

Analytical continuation

$$|t_l^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Prel. results for the pole position

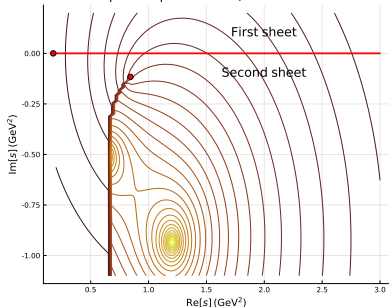
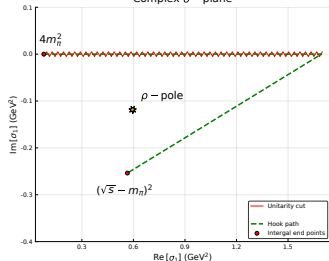
$$m_{\rho^{(a_1(1260))}} = (1.208 \pm 0.004) \text{ GeV},$$

$$\Gamma_{\rho^{(a_1(1260))}} = (0.569 \pm 0.012) \text{ GeV}.$$

Systematic studies:

- variation of ρ -shape
- variation of fit range

Tour to the complex plane

Complex s -plane in the QTB-DISP-modelComplex σ -plane

Analytical continuation

$$|t_l^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Prel. results for the pole position

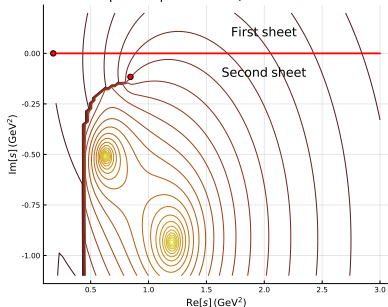
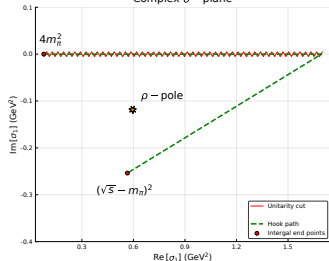
$$m_p^{(a_1(1260))} = (1.208 \pm 0.004) \text{ GeV},$$

$$\Gamma_p^{(a_1(1260))} = (0.569 \pm 0.012) \text{ GeV}.$$

Systematic studies:

- variation of ρ -shape
- variation of fit range

Tour to the complex plane

Complex s -plane in the QTB-DISP-modelComplex σ -plane

Analytical continuation

$$|t_l^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Prel. results for the pole position

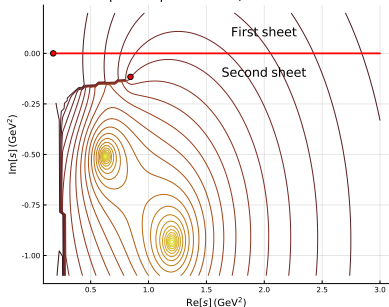
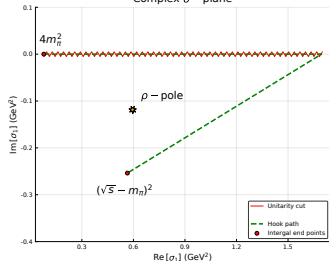
$$m_p^{(a_1(1260))} = (1.208 \pm 0.004) \text{ GeV},$$

$$\Gamma_p^{(a_1(1260))} = (0.569 \pm 0.012) \text{ GeV}.$$

Systematic studies:

- variation of ρ -shape
- variation of fit range

Tour to the complex plane

Complex s -plane in the QTB-DISP-modelComplex σ -plane

Analytical continuation

$$|t_l^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Error propagation using bootstrap analysis.

Prel. results for the pole position

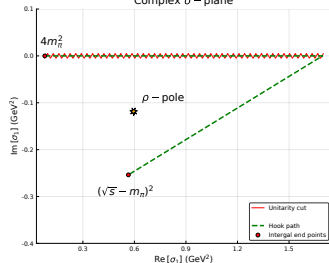
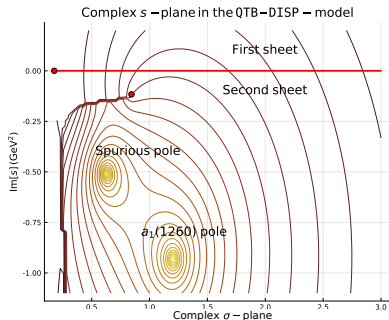
$$m_p^{(a_1(1260))} = (1.208 \pm 0.004) \text{ GeV},$$

$$\Gamma_p^{(a_1(1260))} = (0.569 \pm 0.012) \text{ GeV}.$$

Systematic studies:

- variation of ρ -shape
- variation of fit range

The spurious pole

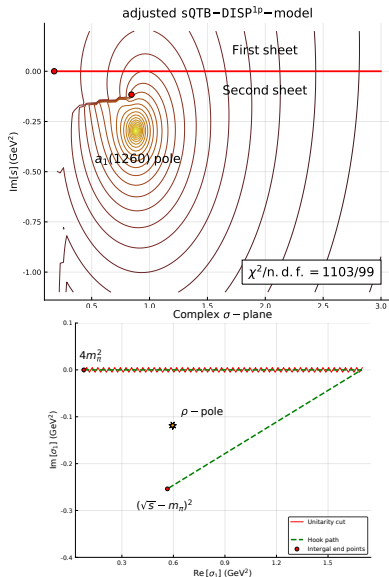


The reason of the spurious pole

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{\textcolor{red}{s}}$$

- Not every pole is the resonance

The spurious pole



The reason of the spurious pole

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Not every pole is the resonance

Can we find the model without the pole

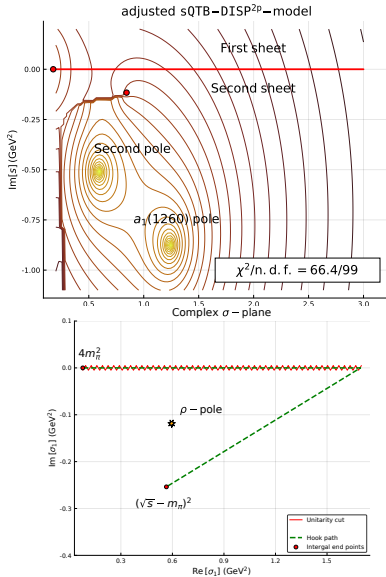
- Denoting $\hat{t}(s) = t(s)/s$
 $2\text{Im } \hat{t}(s) = \hat{t}^*(s) (s\rho(s)) \hat{t}(s),$

- New model

$$c/(P_k(s) - ig^2 s \hat{\rho}(s)/2)$$

$$P_1(s) = m^2 - s \quad \text{does not fit the data}$$

The spurious pole



The reason of the spurious pole

$$\rho(s) \sim \int_{4m_\pi^2}^{(\sqrt{s}-m_\pi)^2} \frac{d\sigma}{2\pi} U_\rho(\sigma) \frac{\lambda^{1/2}(s, \sigma, m_\pi^2)}{s}$$

- Not every pole is the resonance

Can we find the model without the pole

- Denoting $\hat{t}(s) = t(s)/s$

$$2\text{Im} \hat{t}(s) = \hat{t}^*(s) (s\rho(s)) \hat{t}(s),$$

- New model

$$c/(P_k(s) - ig^2 s \hat{\rho}(s)/2)$$

$$P_1(s) = m^2 - s \quad \text{does not fit the data}$$

$$P_2(s) = s(m^2 - s) + h \quad \text{gives a good fit and the same pole}$$

Summary

- Three-body unitarity is important for hadron spectroscopy
 - ▶ both experimental and lattice
 - ▶ The light 1^{++} sector is the first place to study it
- We found a class of models which can satisfy unitarity exactly:
 - ▶ Factorization and final state interaction.
- Free parameters are constrained using ALEPH data (where is Belle, BES data !?).
- Analytical continuation:
 - ▶ requires an accurate integral path deformation.
 - ▶ The pole of the $a_1(1260)$ resonance was found.

The spurious pole

- The dynamics in the channel requires strong influence of the left singularities (the threshold region is important).
- Systematic studies are ongoing.

Thank you for the attention

Thanks to

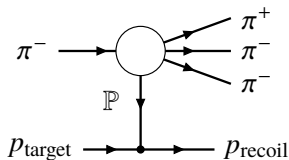
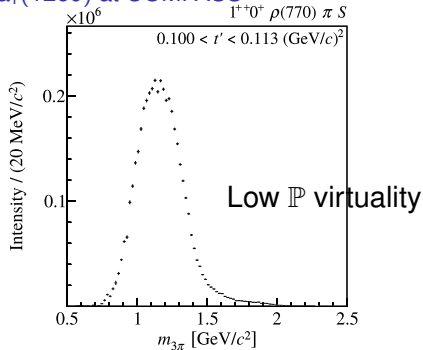
B. Ketzer, COMPASS colleagues,
A. Jackura, A. Pilloni, A. Szczepaniak (JPAC)

EXTRA MATERIAL

1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

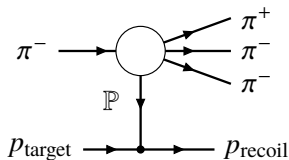
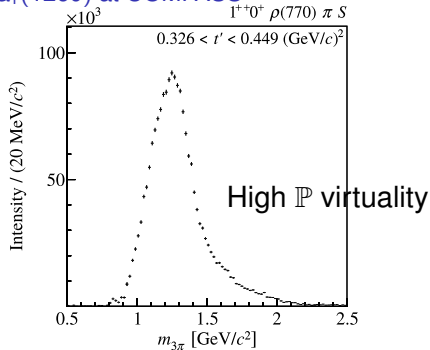
[COMPASS, PRD 95 (2017)]



1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

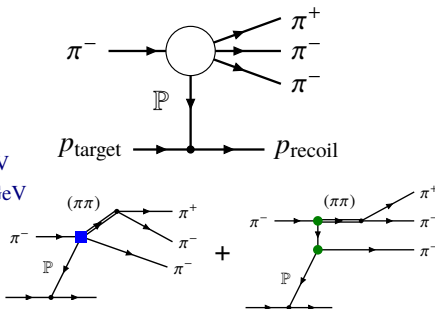
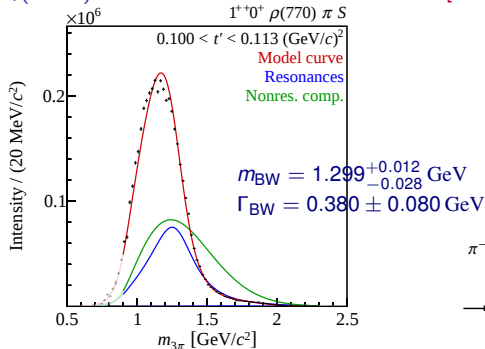
[COMPASS, PRD 95 (2017)]



1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

[COMPASS, PRD 95 (2017)]



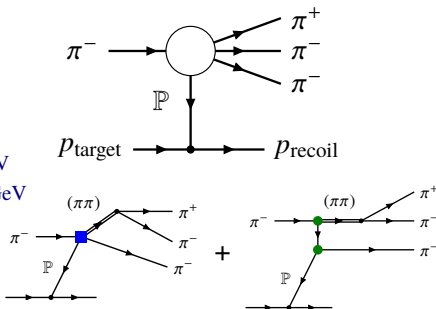
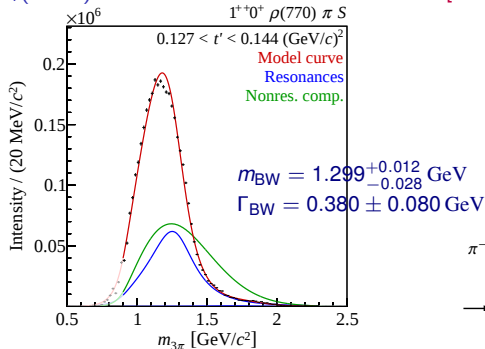
COMPASS Model

- Partial wave analysis to isolate 1^{++} sector.
- Coherent addition of the short and the long range diagrams $\Rightarrow$ (no unitarity)

1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

[COMPASS, PRD 95 (2017)]



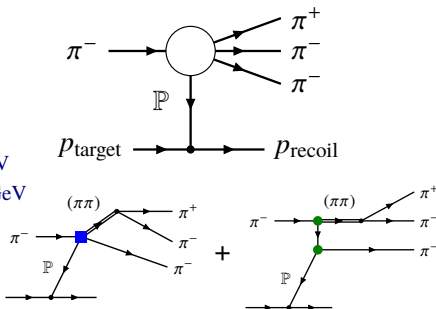
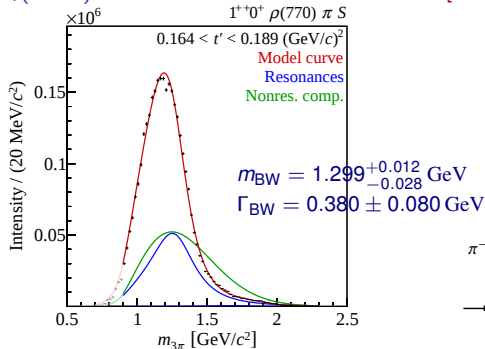
COMPASS Model

- Partial wave analysis to isolate 1^{++} sector.
- Coherent addition of the short and the long range diagrams $\Rightarrow$ (no unitarity)

1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

[COMPASS, PRD 95 (2017)]



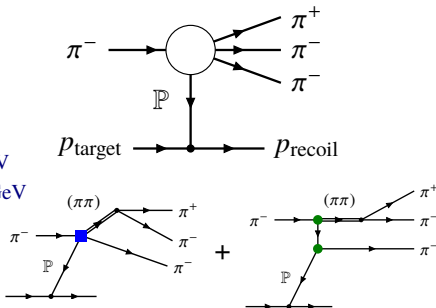
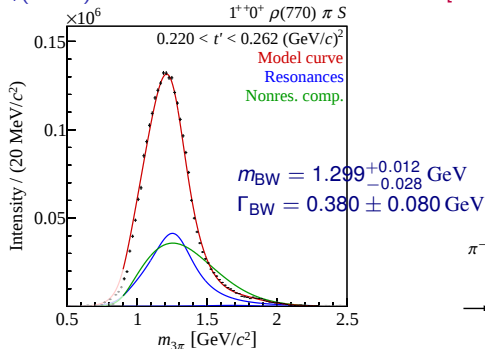
COMPASS Model

- Partial wave analysis to isolate 1^{++} sector.
- Coherent addition of the short and the long range diagrams $\Rightarrow$ (no unitarity)

1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

[COMPASS, PRD 95 (2017)]



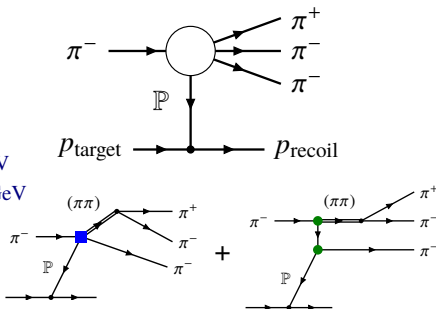
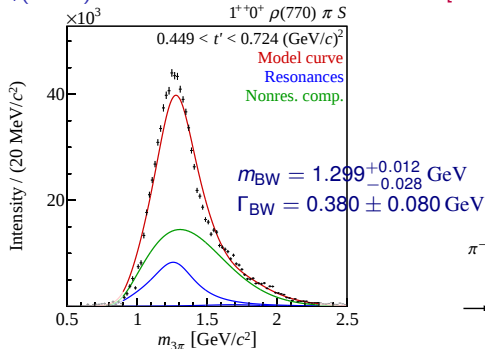
COMPASS Model

- Partial wave analysis to isolate 1^{++} sector.
- Coherent addition of the short and the long range diagrams $\Rightarrow$ (no unitarity)

1^{++} interaction in scattering experiments

$a_1(1260)$ at COMPASS

[COMPASS, PRD 95 (2017)]



COMPASS Model

- Partial wave analysis to isolate 1^{++} sector.
- Coherent addition of the short and the long range diagrams $\Rightarrow$ (no unitarity)

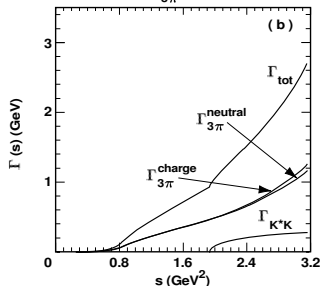
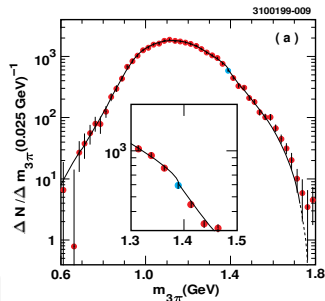
3π at CLEO, [Phys.Rev. D61 (2000) 012002]

Partial wave analysis:

- combined fit with 7 waves.
- 5 isobars ρ , $\rho(1450)$, $f_2(1270)$, σ , $f_0(1370)$.
- significant contribution of $K^*\bar{K}$

Main fit results

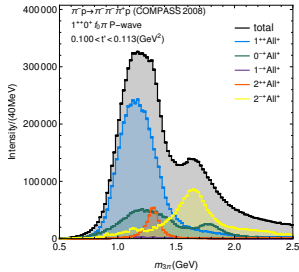
- $m_{a_1} \approx 1.33 \text{ GeV}$,
 $\Gamma_{a_1}(m_{a_1}) \approx 0.814 \text{ GeV}$,
- Branching fractions:
 - $\text{Br}(a_1 \rightarrow \rho\pi) \approx 60\%$,
 - $\text{Br}(a_1 \rightarrow \sigma\pi) \approx 20\%$,
 - $\text{Br}(a_1 \rightarrow f_0(1370)\pi) \approx 7\%$,
 - $\text{Br}(a_1 \rightarrow K^*\bar{K}) \approx 2.2\%$.



$J^{PC} = 1^{++}$ sector

$a_1(1420)$ phenomenon

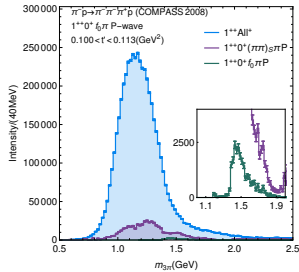
[PRL 115 (2015) 082001]



$J^{PC} = 1^{++}$ sector

$a_1(1420)$ phenomenon

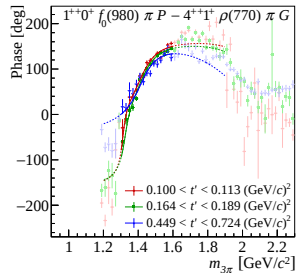
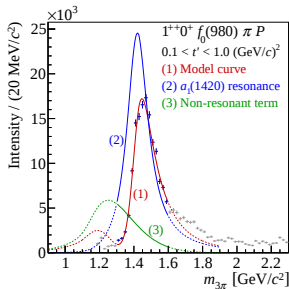
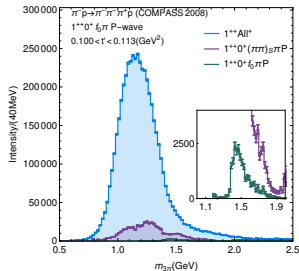
[PRL 115 (2015) 082001]



$J^{PC} = 1^{++}$ sector

$a_1(1420)$ phenomenon

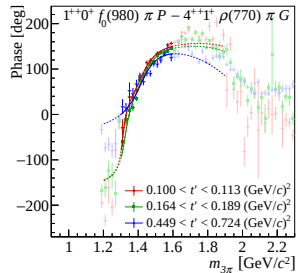
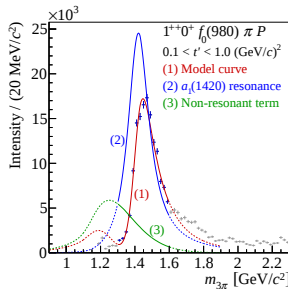
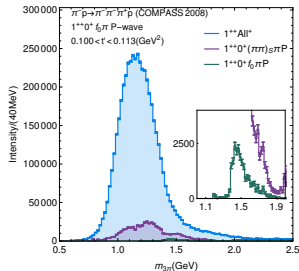
[PRL 115 (2015) 082001]



$J^{PC} = 1^{++}$ sector

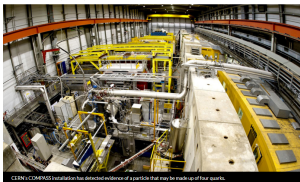
$a_1(1420)$ phenomenon

[PRL 115 (2015) 082001]



COMPASS: SEARCHING FOR NEW PARTICLES

New particle may be made of four quarks



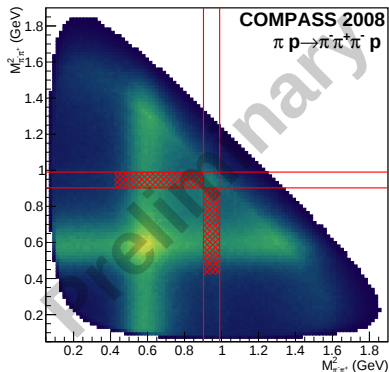
Not something ordinary

- Does not fit to the radial excitation trajectory
- Too close to ground state $a_1(1260)$
- Width narrower than ground state
- Mass is very close to the $K^*\bar{K}$ threshold

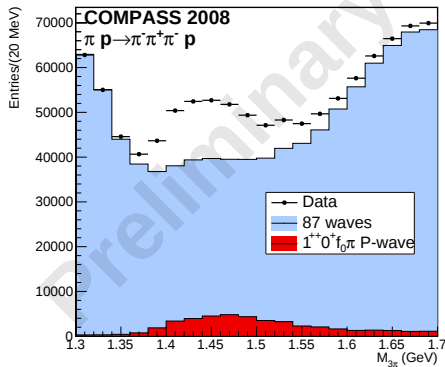
Selection of $a_1(1420)$

Enhancement in COMPASS data $\pi p \rightarrow 3\pi p$

3π Dalitz plot, $\{1.4 \text{ GeV} < M_{3\pi} < 1.5 \text{ GeV}\}$



" f_0 -band-cut" in the Dalitz plot



- A simple cut on the Dalitz plot (f_0 band) enhances bump in $M_{3\pi}$
- Partial wave analysis proves that the origin is the $a_1(1420)$.