Study of charmoniumlike states by amplitude analyses at Belle

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Charmonium production:

- From B decays 711 fb⁻¹, 772 \times 10⁶ $B\bar{B}$ pairs.
- Double charmonium production (all energies): 980 fb⁻¹.





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Charged Z_c^+ states

$\mathbb{Z}_{c}(4430)^+$ quantum numbers



Full 4-dimensional amplitude analysis of $\bar{B}^0 \rightarrow \psi(2S)K^-\pi^+$. Result: $J^P = 1^+$ is preferred, 0^- , 1^- , 2^- and 2^+ hypotheses are excluded at the levels of 3.4σ , 3.7σ , 4.7σ and 5.1σ , respectively. Parameters: $M = 4485^{+22+28}_{-22-11} \text{ MeV}/c^2$, $\Gamma = 200^{+41+26}_{-46-35} \text{ MeV}$. The J^P measurement was confirmed by LHCb in PRL **112**, 222002

The J^r measurement was confirmed by LHCb in PRL **112**, 222002 (2014) with much higher significance.

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A new state $Z_c(4200)^+$ was observed with $M = 4196^{+31}_{-17}^{+29}_{-13} \text{ MeV}/c^2$, $\Gamma = 370^{+70}_{-70}^{+70}_{-132} \text{ MeV}$, 6.2σ . $J^P = 1^+$ is preferred, 0^- , 1^- , 2^- , 2^+ are excluded at the levels of 6.1σ , 7.4σ , 4.4σ and 7.0σ , respectively. Evidence for $Z_c(4430)^+ \rightarrow J/\psi\pi^+$ was found.

The $Z_c(4200)^+$ is not confirmed, but it may be the same state as the $Z_c(4240)^+$ found by LHCb in $\bar{B}^0 \to \psi(2S)K^-\pi^+$ (if it has $J^P = 1^+$ that is excluded at 1σ only).

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$X^*(3860)$ in $e^+e^- ightarrow J/\psi Dar{D}$

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The X(3915) was seen by first seen by Belle in $B \rightarrow J/\psi \omega K$ [PRL **94**, 182002 (2005)]:



Then the X(3915) was seen by BaBar in the same decay mode [PRD **82**, 011101 (2010)].



It was also seen by Belle [PRL **104**, 092001 (2010)] and BaBar [PRD **86**, 072002 (2012)] in $\gamma\gamma \rightarrow J/\psi\omega$. The BaBar analysis measured $J^P = 0^+$:



Resulting identification was: $X(3915) = \chi_{c0}(2P)$

But this identification is doubtful because of: low width; $\chi_{c0}(2P) \rightarrow J/\psi\omega$ is OZI-suppressed; the difference between the X(3915) and $\chi_{c2}(2P)$ masses is too small, ... [see F. K. Guo and U. G. Meissner, PRD **86**, 091501 (2012), S. L. Olsen PRD **91**, 057501 (2015)].

An alternative $\chi_{c0}(2P)$ candidate is searched for in the process $e^+e^- \rightarrow J/\psi D\bar{D}$.

Reconstruction

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$$J/\psi \rightarrow e^+e^-, \mu^+\mu^-.$$

- One *D* is reconstructed, the other is identified by the recoil mass $(M_{\rm rec}(J/\psi, D))$. Both D^0 and D^+ are used.
- $D^0 \to K^- \pi^+, K^0_S \pi^+ \pi^-, K^- \pi^+ \pi^0, K^- \pi^+ \pi^+ \pi^-$ (4 channels).
- $D^+ \to K_S^0 \pi^+, K^- \pi^+ \pi^+, K_S^0 \pi^+ \pi^0, K^- \pi^+ \pi^+ \pi^0, K_S^0 \pi^+ \pi^+ \pi^-$ (5 channels).
- Separation of signal and background using the MLP neural network.
- Global optimization of the selection requirements for (4 variables per *D* channel: signal regions in *M_{J/ψ}*, *M_D*, *M_{rec}(J/ψ, D*) and MLP output cutoff value).

Resulting sample: 103 events with 24.9 \pm 1.1 \pm 1.6 background events.



$$S(\Phi) = \sum_{\substack{\lambda_{\text{beam}} = -1, 1 \\ \lambda_{\ell\ell} = -1, 1}} \left| \sum_{X^*} A_{\lambda_{\text{beam}} \lambda_{\ell\ell}}(\Phi) A_{X^*}(M_{D\bar{D}}) \right|^2,$$
(1)

Here, $A_{\lambda_{\text{beam}} \lambda_{\ell\ell}}(\Phi)$ is the signal amplitude calculated using the helicity formalizm (the phase space Φ is 6-dimensional). For resonance, A_{X^*} = relativistic Breit-Wigner. For nonresonant amplitude,

 $A_{X^*} = \sqrt{F_{D\bar{D}}(M_{D\bar{D}})}$, where $F_{D\bar{D}}(M_{D\bar{D}})$ is the nonresonant amplitude form factor ($F_{D\bar{D}} = 1$ by default). Alternatives: mass dependence of NRQCD prediction for $e^+e^- \rightarrow \psi \chi_c$ [PRD **77**, 014002 (2008)], $F_{D\bar{D}} = M_{D\bar{D}}^{-4}$ [Victor Chernyak, based on PLB **612**, 215 (2005)].







Red dashed line - only background and nonresonant amplitudes, blue solid line - X^* , $J^{PC} = 0^{++}$.



Fit results in the default model. For the 2^{++} hypothesis, there are three solutions (fit is started 1000 times from random initial values in order to check for that).

JPC	Mass, MeV/ c^2	Width, MeV	Significance (Wilks)
0++	3862^{+26}_{-32}	201^{+154}_{-67}	9.1σ
2++	3879_{-17}^{+20}	171^{+129}_{-62}	8.0σ
2++	$3879^{+\bar{1}\bar{7}}_{-17}$	148^{+108}_{-50}	8.0σ
2++	3883^{+26}_{-24}	227^{+201}_{-125}	8.0σ



Global significance is determined from $\Delta(-2 \ln L)$ distributions.

Model	Significance (global)
Default (constant nonresonant)	8.5σ
NRQCD nonresonant	7.6σ
$M_{D\bar{D}}^{-4}$ nonresonant	6.5σ
Background mass calculation	8.4σ
Optimization $(a = 4)$	8.1σ
Optimization $(a = 6)$	8.1σ

The minimal significance is $6.5\sigma \Rightarrow$ the $X^*(3860)$ is observed.

Error source	Mass	Width
Nonresonant amplitude model	$^{+40.2}_{-0.0}$	+0.0 -82.0 +0.0
Signal model	-10.2	-4.0 +32.6
Fit blas	+0.0	$^{-0}_{+71.1}$
Background mass calculation	$^{-3.1}_{+0.0}$	-0.0 +40.0 -0.0
D mass	± 0.2	
Total	$\substack{+40.2\\-13.3}$	$\substack{+87.9\\-82.1}$

Comparison of $J^{PC} = 0^{++}$ and 2^{++}

PRD **95**, 112003 (2017)

Toy MC pseudoexperiments are generated in accordance with the fit results with $J^{PC} = 0^{++}$ and 2^{++} and fitted by both hypotheses. Result:



The $J^{PC} = 2^{++}$ hypothesis is excluded at the level of 3.8σ for the default model (shown in the histogram) and 2.5σ with systematic uncertainty. The confidence level of the $J^{PC} = 0^{++}$ hypothesis is 77% (default model).

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Cross section measurement

PRD 95, 112003 (2017)

The Born cross section of $e^+e^- \rightarrow J/\psi X^*(3860)(\rightarrow D\bar{D})$ is measured at each energy point:

Data set	Energy, GeV	$\sigma^{(\mathrm{Born})}_{e^+e^- ightarrow J/\psi X^*(3860)(ightarrow Dar{D})}$, fb
$\Upsilon(1S)$	9.46	$77^{+66}_{-66}{}^{+9}_{-7}$
Ƴ(2 <i>S</i>)	10.02	$6.9^{+12.6}_{-12.6}{}^{+0.9}_{-0.7}$
Ƴ(3 <i>S</i>)	10.36	$77^{+85}_{-85}^{+11}_{-8}$
Continuum	10.52	$5.5^{+5.7}_{-5.7}$
Ƴ(4 <i>S</i>)	10.58	$21.7^{+3.9}_{-4.3}^{+2.9}_{-2.1}$
Υ(5 <i>S</i>)	10.87	$17.9^{+7.2}_{-7.3}{}^{+2.4}_{-1.8}$

NRQCD [PRD **77**, 014002 (2008)] (usually smaller than measured cross sections): $\sigma_{e^+e^- \rightarrow J/\psi_{\chi_{c0}(2P)}}(10.6 \text{ GeV}) = 9.1 \text{ fb}$



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- A new charmoniumlike state $X^*(3860)$ is observed (6.5 σ with systematic error).
- Parameters: $M = 3862^{+26}_{-32}^{+40} \text{ MeV}/c^2$, $\Gamma = 201^{+154}_{-67}^{+88} \text{ MeV}.$
- The $J^{PC} = 0^{++}$ hypothesis is favored over the 2^{++} hypothesis at the level of 2.5σ .



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- Quantum numbers: $J^{PC} = 0^{++}$.
- Production: in S-wave. Same for $\chi_{c0}(1P)$, measured in PRD **70**, 071102 (2004).
- The $\chi_{c0}(2P)$ mass in Ebert-Faustov-Galkin model [PRD **67**, 014027 (2003)]: 3854 MeV/ c^2 , in Godfrey-Isgur model [PRD **32**, 189 (1985)]: 3916 MeV/ c^2 .
- Mass difference (potential models: $\sim 0.6 0.9$): $r_c = (m_{\chi_{c2}(2P)} - m_{\chi_{c0}(2P)})/(m_{\chi_{c2}(1P)} - m_{\chi_{c0}(1P)}) = 0.46^{+0.25}_{-0.34}$
- Decay: $\chi_{c0}(2P)$ should primarily decay to $D\overline{D}$ (observation mode for the X*(3860)) and not $J/\psi\omega$.
- The X*(3860) agrees with the peak in $\gamma\gamma$ data ($M = 3837.6 \pm 11.5 \text{ MeV}/c^2$, $\Gamma = 221 \pm 19 \text{ MeV}$).

The $X^*(3860)$ is a better $\chi_{c0}(2P)$ candidate than the X(3915).

Thank you for attention!