

Измерение сечения  $e^+e^- \rightarrow \pi^+\pi^-$   
с детектором КМД-3 на коллайдере  
ВЭПП-2000 и его последствия для  
проблемы аномального  
магнитного момента мюона

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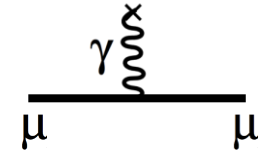
Научная сессия  
Объединенного  
ученого совета  
по физическим  
наукам СО РАН

29 ноября 2023

# The basics

**Gyromagnetic ratio  $g$**  connects magnetic moment  $\mu$  and spin  $s$

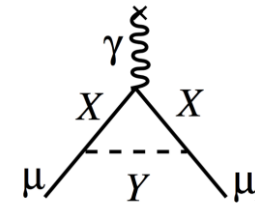
$$\vec{\mu}_s = g \frac{e}{2m} \vec{S}$$



For point-like particle  $g = 2$

**Anomalous magnetic moment  $a$**  arises in higher-orders

$$a = (g - 2)/2$$



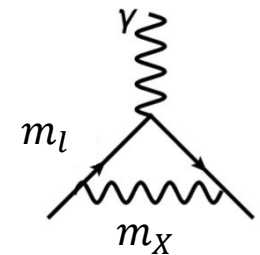
$$a_e \approx a_\mu \approx \frac{\alpha}{2\pi} \approx 10^{-3} \quad (\text{QED dominated})$$

**Idea of experiment:** by comparing measured value of  $a$  with the theory prediction we probe extra contributions beyond theory expectations

$$a_\mu(\text{strong})/a_\mu(\text{QED}) \approx 6 \times 10^{-5} \quad a_\mu(\text{weak})/a_\mu(\text{QED}) \approx 10^{-6}$$

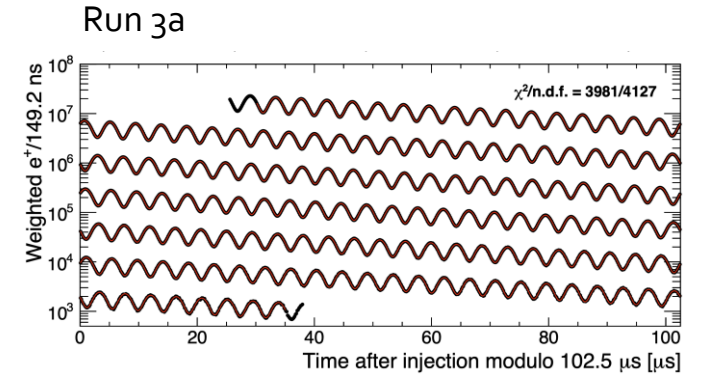
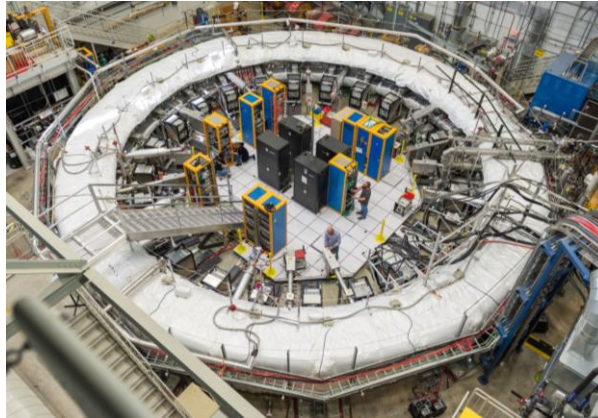
**Why muon?** For massive fields there is natural scaling, which enhances contribution to  $a_\mu$  by  $(m_\mu/m_e)^2 \sim 43000$  compared to  $a_e$

$$\Delta a \sim \left( \frac{m_l}{m_X} \right)^2$$



# Generations of $a_\mu$ measurements

FNAL Run 2-3  
(USA)



$$g_\mu(\text{эксп}) = 0.001\,165\,920\,55(24) \quad \text{FNAL}_{2023}$$

$$g_\mu(\text{теория}) = 0.001\,165\,918\,10(43) \quad \text{WP}_{2020}$$

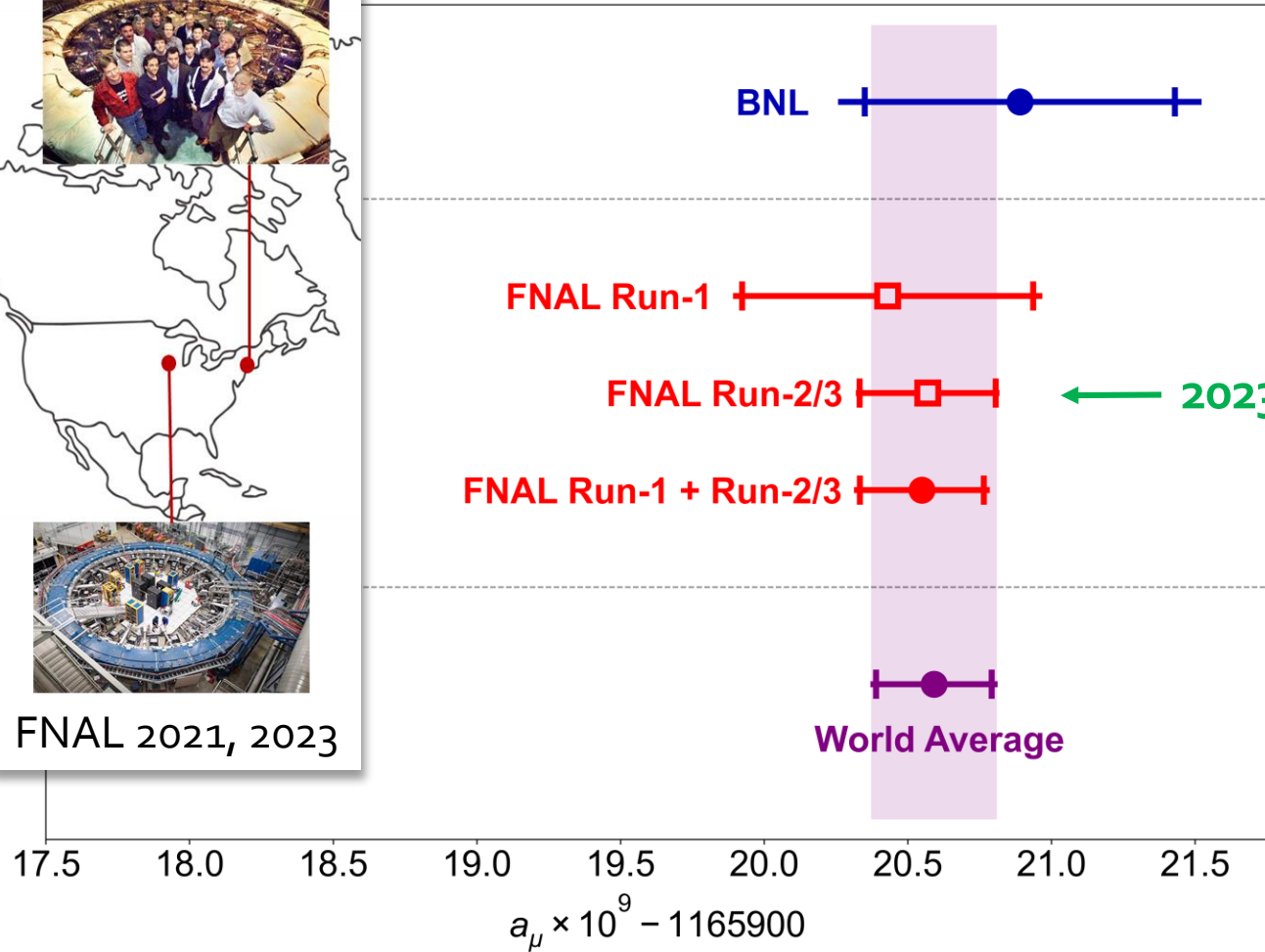
QED

Strong

Weak

Contributions of known interactions

# Muon G-2 2023 result



$$a_\mu(\text{Exp}) = 0.00116592059(22) \quad [190 \text{ ppb}]$$

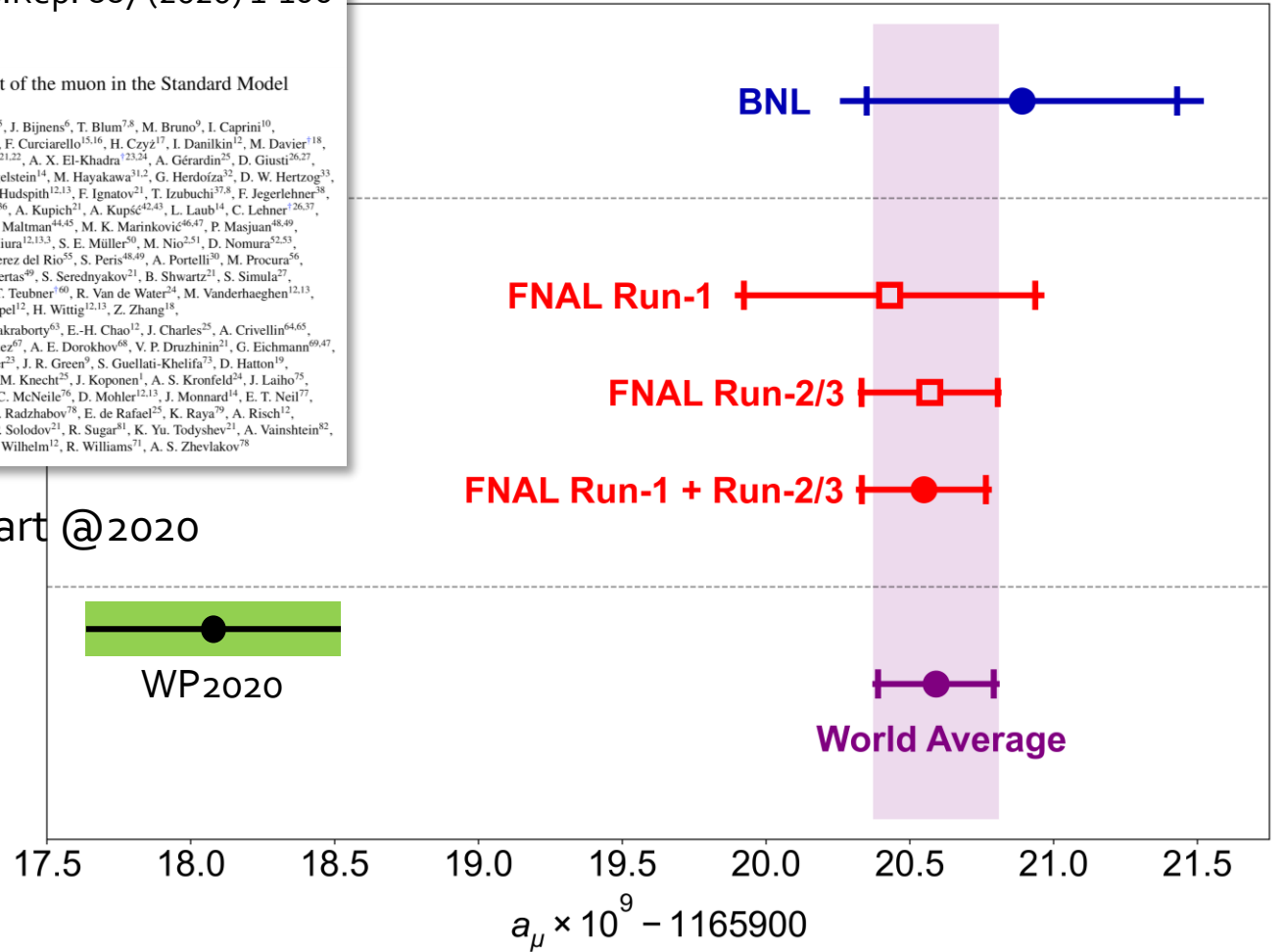
# Experiment vs SM prediction

## Muon G-2 Theory Initiative Consortium of >100 theorists and experimental physicists "White paper", Phys.Rep. 887 (2020) 1-166

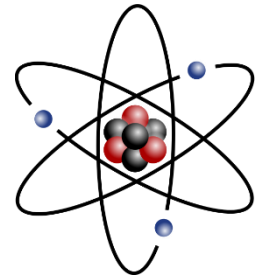
The anomalous magnetic moment of the muon in the Standard Model

T. Aoyama<sup>1,2,3</sup>, N. Asmussen<sup>4</sup>, M. Benayoun<sup>5</sup>, J. Bijnens<sup>6</sup>, T. Blum<sup>7,8</sup>, M. Bruno<sup>9</sup>, I. Caprini<sup>10</sup>,  
C. M. Carloni Calame<sup>11</sup>, M. Cè<sup>9,12,13</sup>, G. Colangelo<sup>14</sup>, F. Curciarello<sup>15,16</sup>, H. Czyz<sup>17</sup>, I. Danilkin<sup>12</sup>, M. Davier<sup>18</sup>,  
C. T. H. Davies<sup>19</sup>, M. Della Morte<sup>20</sup>, S. I. Eidelman<sup>21,22</sup>, A. X. El-Khadra<sup>23,24</sup>, A. Gérardin<sup>25</sup>, D. Giusti<sup>26,27</sup>,  
M. Golterman<sup>28</sup>, Steven Gottlieb<sup>29</sup>, V. Gülpers<sup>30</sup>, F. Hagelstein<sup>14</sup>, M. Hayakawa<sup>31,2</sup>, G. Herdoíza<sup>32</sup>, D. W. Hertzog<sup>33</sup>,  
A. Hoecker<sup>34</sup>, M. Hoferichter<sup>14,35</sup>, B.-L. Hoid<sup>36</sup>, R. J. Hudspith<sup>12,13</sup>, F. Ignatov<sup>21</sup>, T. Izubuchi<sup>37,8</sup>, F. Jegerlehner<sup>38</sup>,  
L. Jin<sup>7,8</sup>, A. Keshavarzi<sup>39</sup>, T. Kinoshita<sup>40,41</sup>, B. Kubis<sup>36</sup>, A. Kupich<sup>21</sup>, A. Kupś<sup>42,43</sup>, L. Laub<sup>14</sup>, C. Lehner<sup>26,37</sup>,  
L. Lellouch<sup>25</sup>, I. Logashenko<sup>21</sup>, B. Malaescu<sup>5</sup>, K. Maltman<sup>44,45</sup>, M. K. Marinkovic<sup>46,47</sup>, P. Masjuan<sup>48,49</sup>,  
A. S. Meyer<sup>27</sup>, H. B. Meyer<sup>12,13</sup>, T. Mibe<sup>11</sup>, K. Miura<sup>12,13,3</sup>, S. E. Müller<sup>50</sup>, M. Nio<sup>2,51</sup>, D. Nomura<sup>52,53</sup>,  
A. Nyffeler<sup>12</sup>, V. Pascalutsa<sup>12</sup>, M. Passera<sup>54</sup>, E. Perez del Rio<sup>55</sup>, S. Peris<sup>48,49</sup>, A. Portelli<sup>30</sup>, M. Procura<sup>56</sup>,  
C. F. Redmer<sup>12</sup>, B. L. Roberts<sup>57</sup>, P. Sánchez-Puertas<sup>49</sup>, S. Serednyakov<sup>21</sup>, B. Shwartz<sup>21</sup>, S. Simula<sup>27</sup>,  
D. Stöckinger<sup>58</sup>, H. Stöckinger-Kim<sup>58</sup>, P. Stoffer<sup>59</sup>, T. Teubner<sup>60</sup>, R. Van de Water<sup>21</sup>, M. Vanderhaeghe<sup>12,13</sup>,  
G. Venanzoni<sup>61</sup>, G. von Hippel<sup>12</sup>, H. Wittig<sup>12,13</sup>, Z. Zhang<sup>18</sup>,  
M. N. Achasov<sup>21</sup>, A. Bashir<sup>62</sup>, N. Cardoso<sup>47</sup>, B. Chakraborty<sup>63</sup>, E.-H. Chao<sup>12</sup>, J. Charles<sup>25</sup>, A. Crivellin<sup>64,65</sup>,  
O. Deineka<sup>12</sup>, A. Denig<sup>12,13</sup>, C. DeTar<sup>66</sup>, C. A. Dominguez<sup>67</sup>, A. E. Dorokhov<sup>68</sup>, V. P. Druzhinin<sup>21</sup>, G. Eichmann<sup>69,47</sup>,  
M. Fael<sup>70</sup>, C. S. Fischer<sup>71</sup>, E. Gámiz<sup>72</sup>, Z. Gelzer<sup>25</sup>, J. R. Green<sup>9</sup>, S. Guellati-Khelifa<sup>73</sup>, D. Hattori<sup>19</sup>,  
N. Hermansson-Truedsson<sup>14</sup>, S. Holz<sup>36</sup>, B. Hörz<sup>74</sup>, M. Knecht<sup>25</sup>, J. Koponen<sup>1</sup>, A. S. Kronfeld<sup>24</sup>, J. Laiho<sup>75</sup>,  
S. Leupold<sup>42</sup>, P. B. Mackenzie<sup>24</sup>, W. J. Marciano<sup>37</sup>, C. McNeile<sup>76</sup>, D. Mohler<sup>12,13</sup>, J. Monnard<sup>14</sup>, E. T. Neil<sup>77</sup>,  
A. V. Nesterenko<sup>68</sup>, K. Ottnad<sup>12</sup>, V. Pauk<sup>12</sup>, A. E. Radzhabov<sup>78</sup>, E. de Rafael<sup>25</sup>, K. Raya<sup>79</sup>, A. Risch<sup>12</sup>,  
A. Rodríguez-Sánchez<sup>8</sup>, P. Roig<sup>80</sup>, T. San José<sup>12,13</sup>, E. P. Solodov<sup>21</sup>, R. Sugar<sup>81</sup>, K. Yu. Todyshev<sup>21</sup>, A. Vainshtein<sup>82</sup>,  
A. Vaquero Avilés-Casco<sup>66</sup>, E. Wei<sup>71</sup>, J. Wilhelm<sup>12</sup>, R. Williams<sup>71</sup>, A. S. Zhevlakov<sup>78</sup>

State-of-art @2020

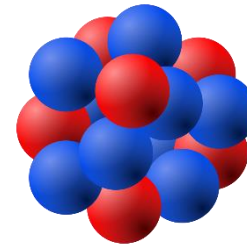


# SM prediction for $a_\mu$



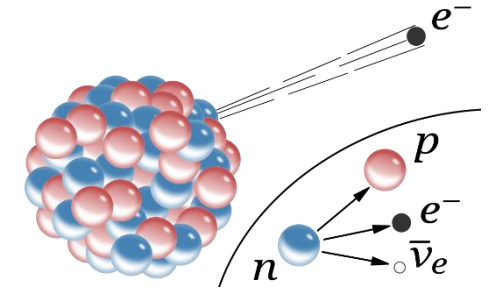
Electromagnetic interactions

0.001 165 847 19 (0.1)



Strong interactions

0.000 000 069 37 (43)



Weak interactions

0.000 000 001 54 (1)

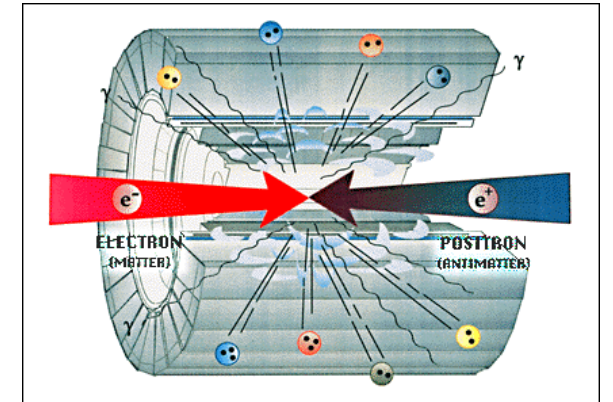
$$a_\mu = 0.001\ 165\ 918\ 10\ (43)$$

The uncertainty is dominated by contribution of strong interactions

# Contribution of exclusive hadronic cross sections to $a_\mu$

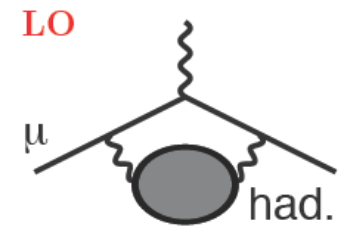
Hadronic contribution can be calculated via dispersion relation, using measured cross section of hadron production in  $e^+e^-$  annihilation:

$$a_\mu^{had}(LO) = \frac{1}{4\pi^3} \int \sigma^0(e^+e^- \rightarrow X) K_\mu(s) ds$$



In exclusive approach, we calculate  $a_\mu$  integral for each final state and sum them:

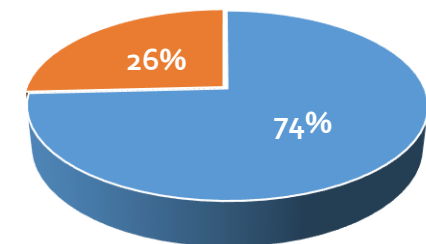
$$a_\mu^{had}(LO) = \sum_{X=\pi^0\gamma, \pi^+\pi^-, \dots} \frac{1}{4\pi^3} \int \sigma^0(e^+e^- \rightarrow X) K_\mu(s) ds$$



$e^+e^- \rightarrow \pi^+\pi^-$  gives by far the largest contribution to the integral – about 74% (and the largest contribution to uncertainty)

$\sigma(e^+e^- \rightarrow \pi^+\pi^-)$  required to be measured with <1% precision ( $\rightarrow 0.1\%$ )

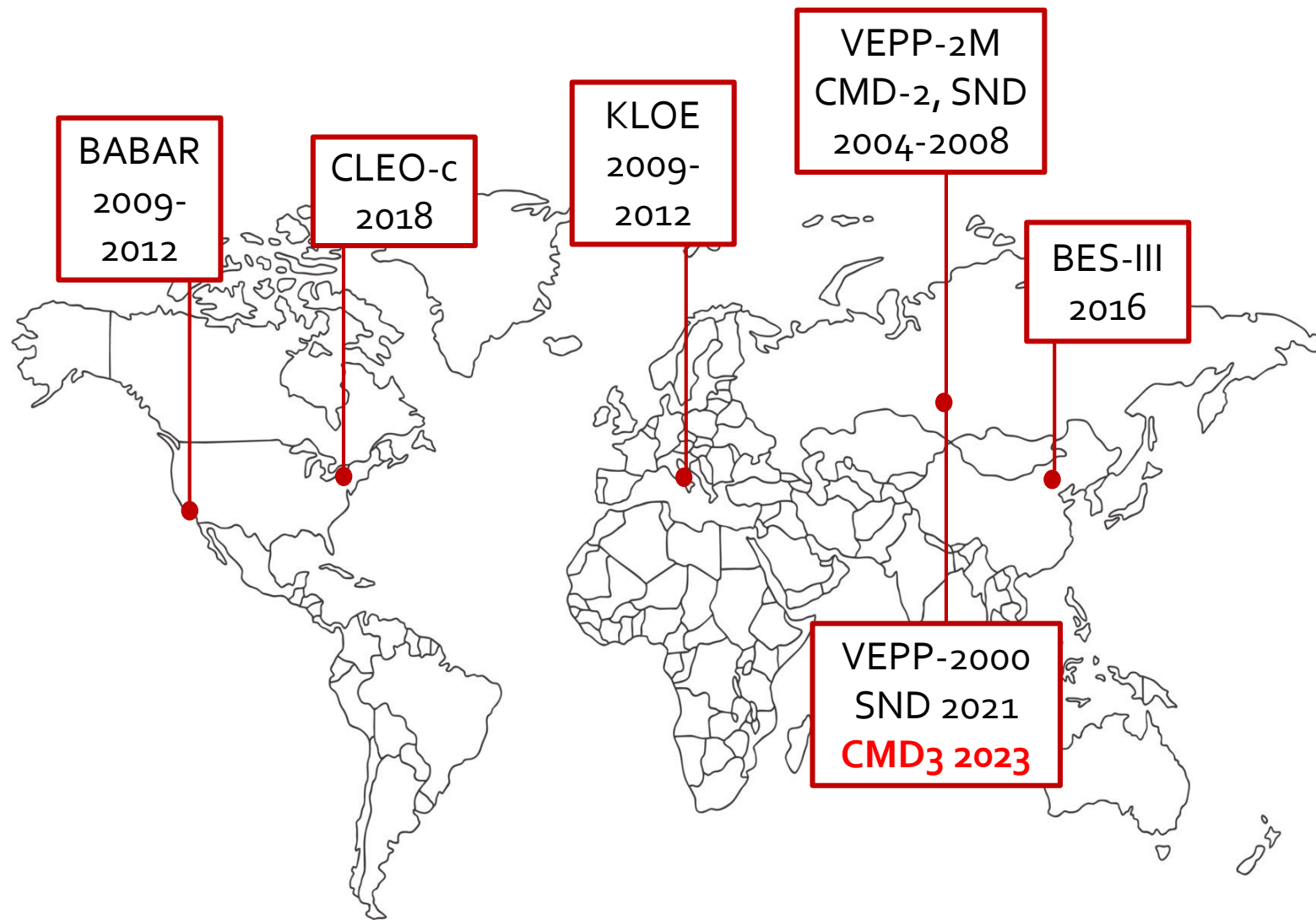
All the rest



$e^+e^- \rightarrow \pi^+\pi^-$

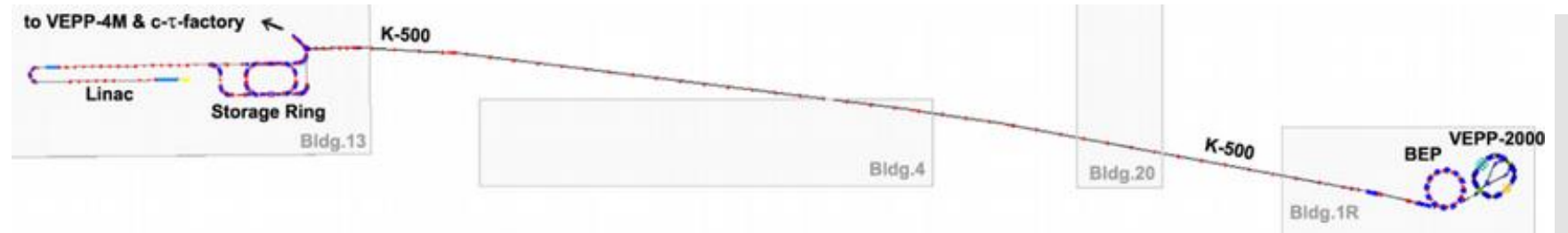
# Measurements of $e^+e^- \rightarrow \pi^+\pi^-$

There are several measurements of  $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$  with sub-percent systematic accuracy



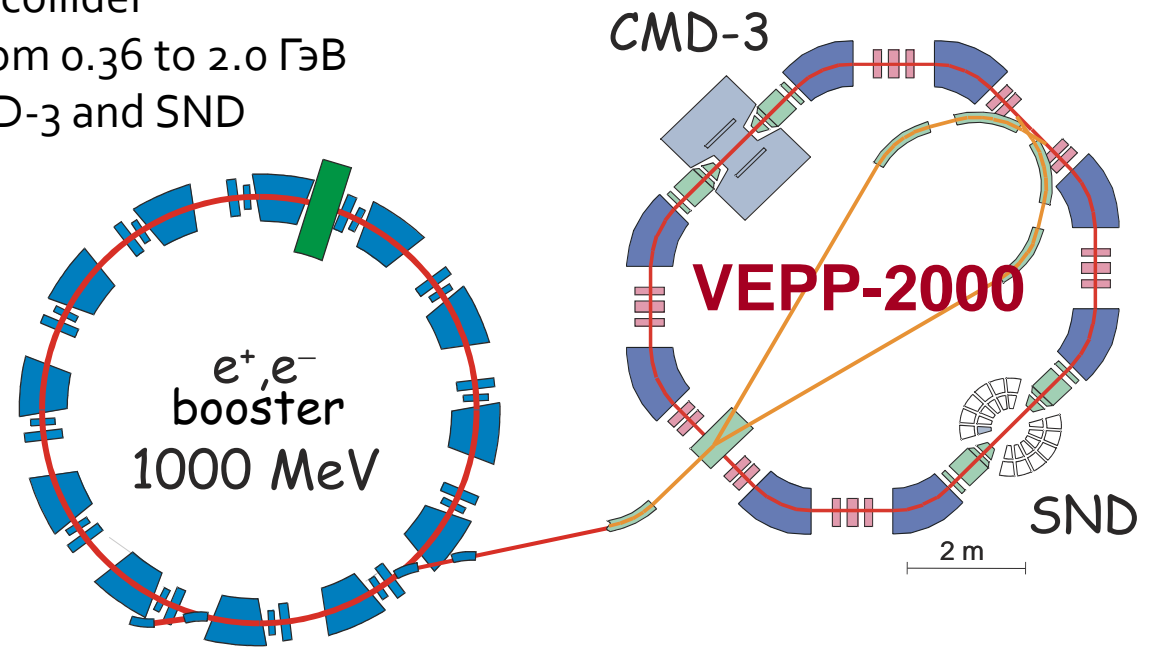


# VEPP-2000 collider



Electron-positron collider  
 Covers c.m. energy range from 0.36 to 2.0 ГэВ  
 Two experiments – CMD-3 and SND

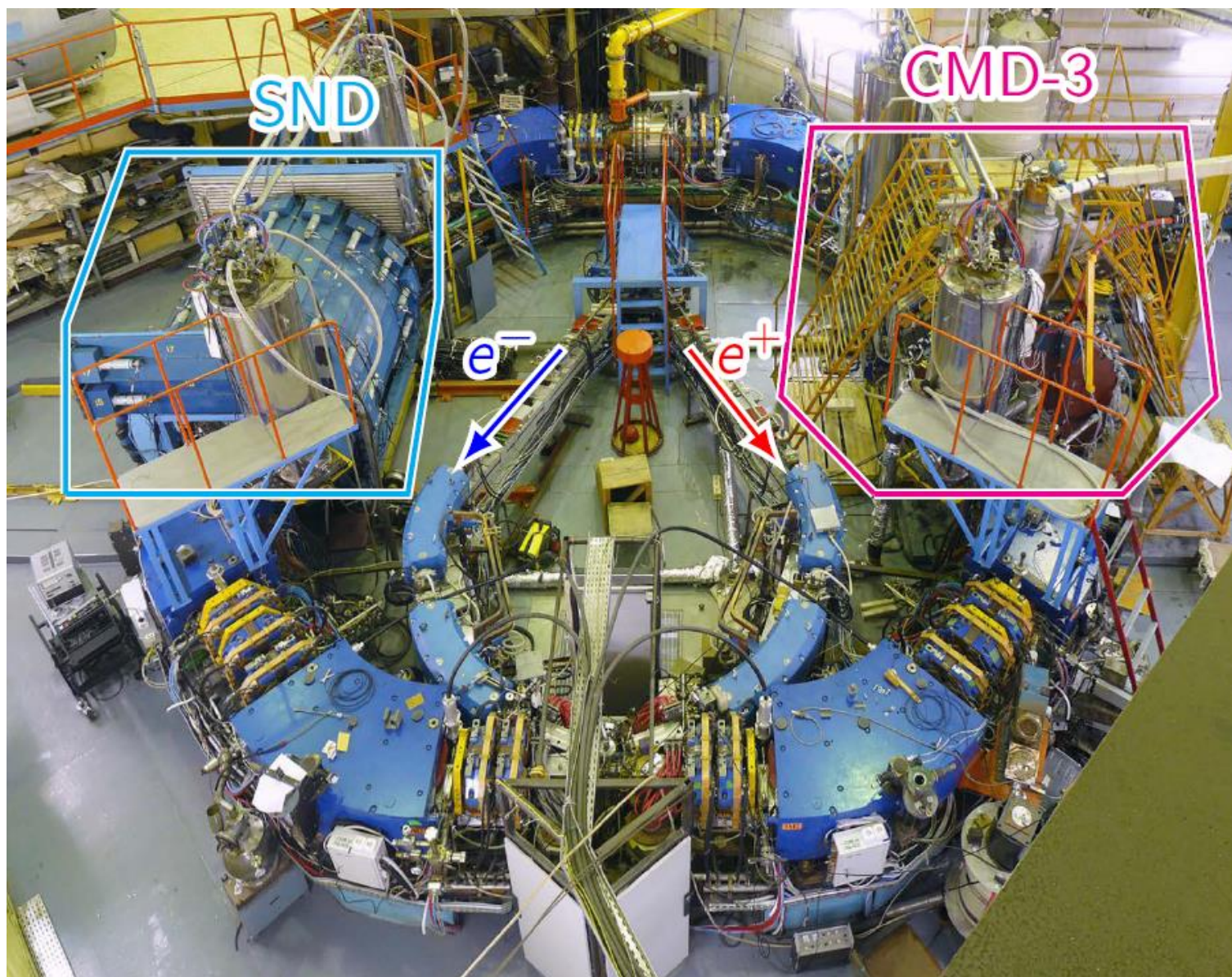
Design parameters @ 1 GeV	
Circumference	24.388 m
Beam energy	150 ÷ 1000 MeV
N of bunches	1×1
N of particles	1×10 <sup>11</sup>
Betatron tunes	4.14 / 2.14
Beta*	8.5 cm
BB parameter	0.1
Luminosity	1×10 <sup>32</sup> cm <sup>-2</sup> s <sup>-1</sup>



“Round beam” optics

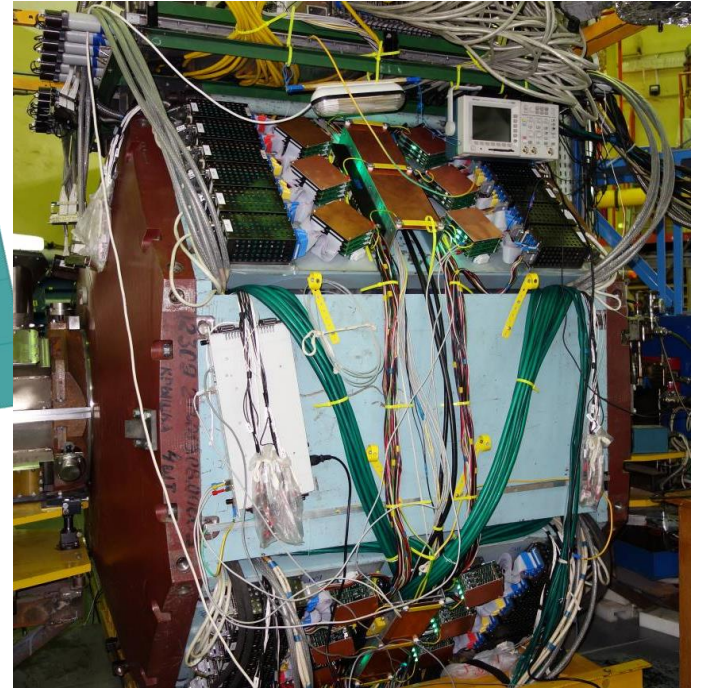
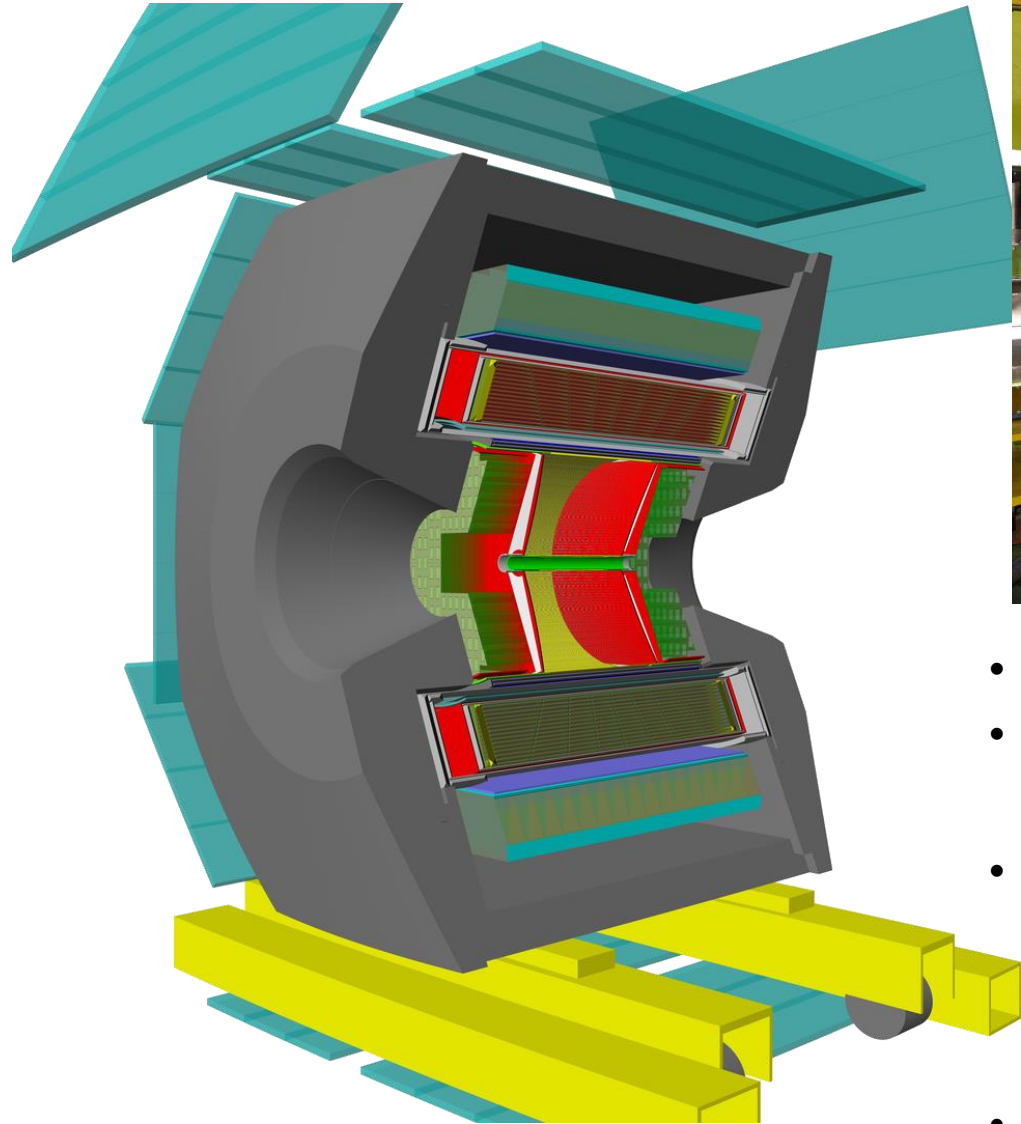
Energy monitoring by Compton backscattering ( $\sigma_{\sqrt{s}} \approx 0.1$  MeV)

# VEPP-2000



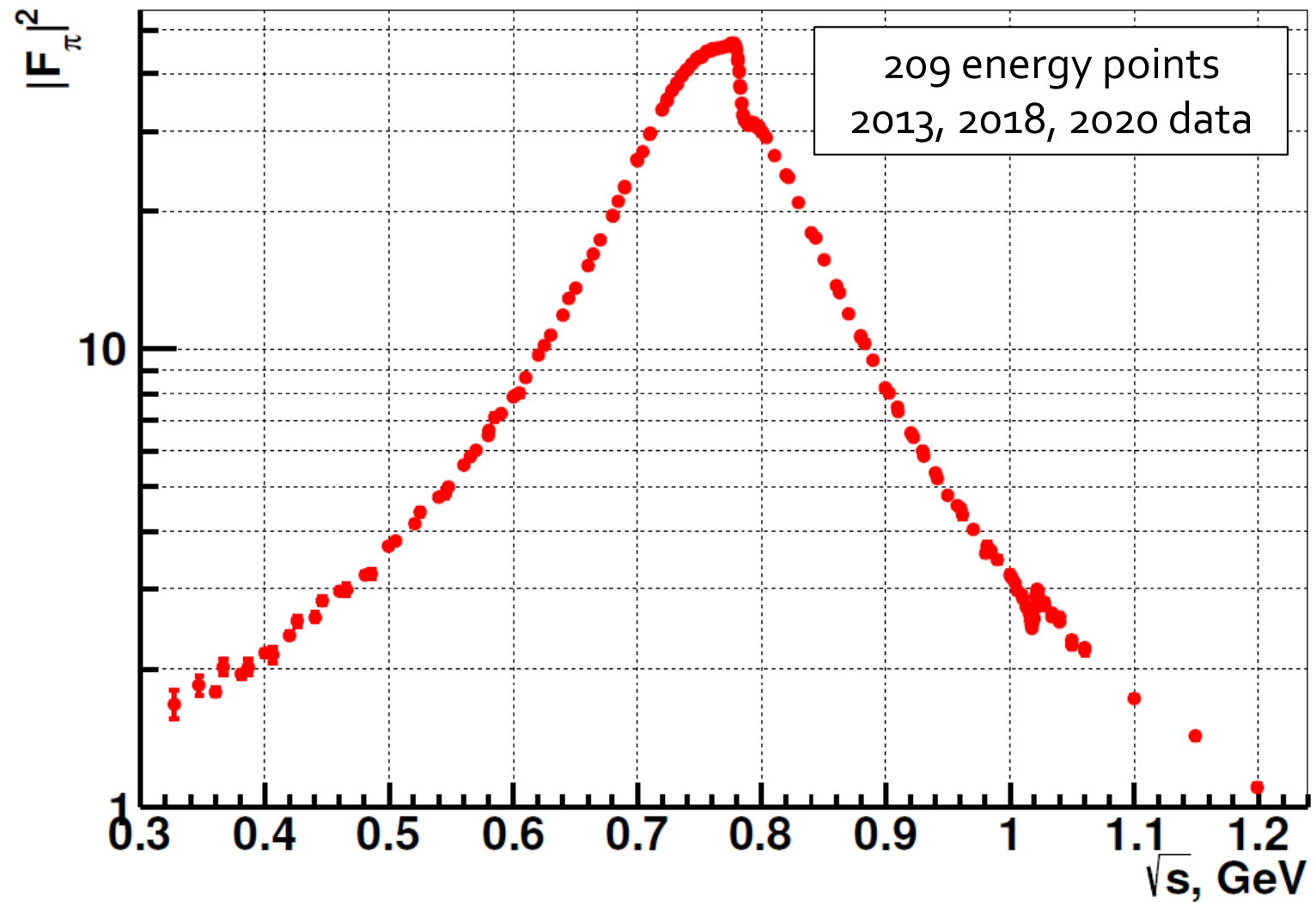
# CMD-3 Detector

\*Cryogenic  
Magnetic Detector

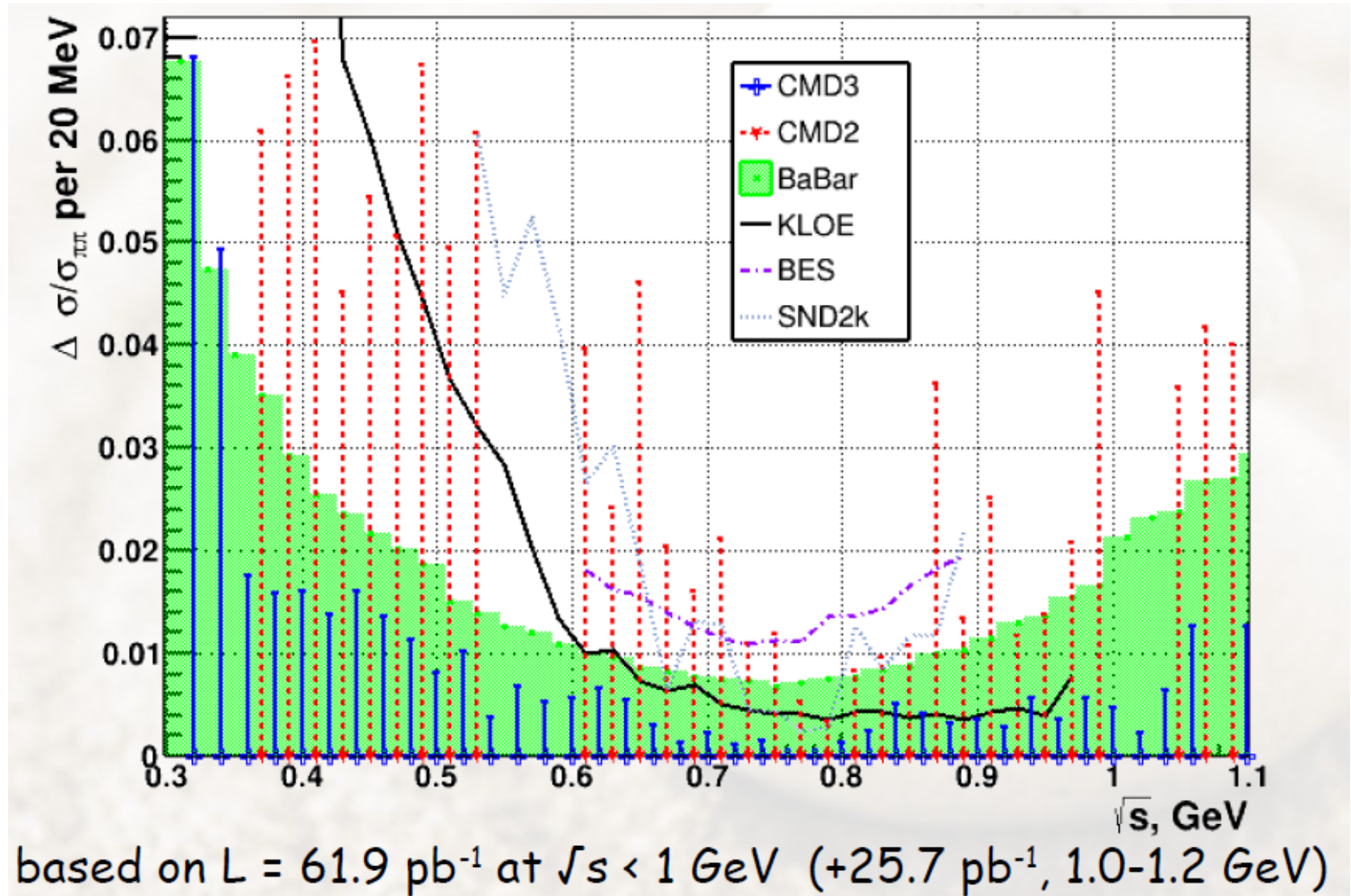


- Magnetic field 1.0-1.3 T
- Drift chamber
  - $\sigma_{R\phi} \sim 100 \mu, \sigma_z \sim 2 - 3 \text{ mm}$
- EM calorimeter (LXE, CsI, BGO),  $13.5 X_0$ 
  - $\sigma_E/E \sim 3\% - 10\%$
  - $\sigma_\theta \sim 5 \text{ mrad}$
- TOF
- Muon counters

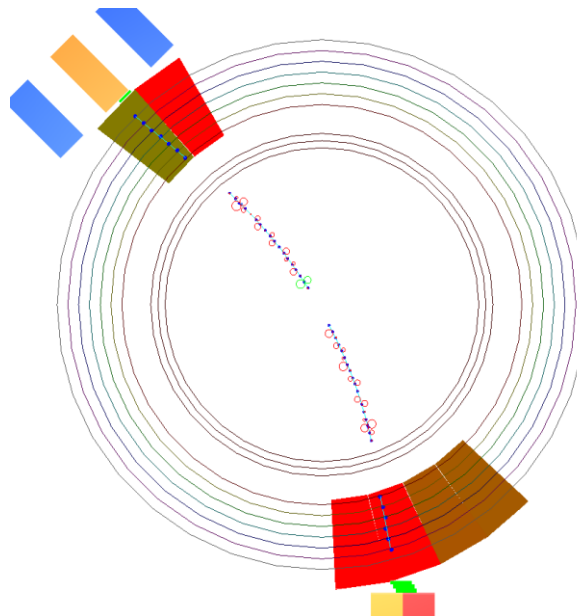
# Measurement of $e^+e^- \rightarrow \pi^+\pi^-$ at CMD-3



# Statistical precision of CMD-3 data



# Three methods of separation of $e^+e^-$ , $\mu^+\mu^-$ , $\pi^+\pi^-$



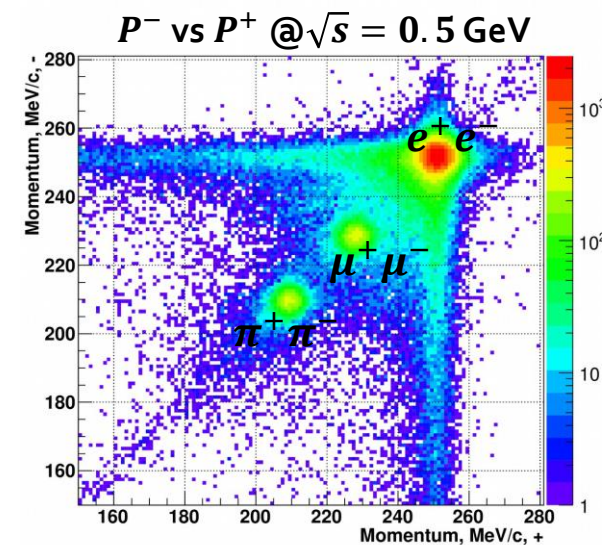
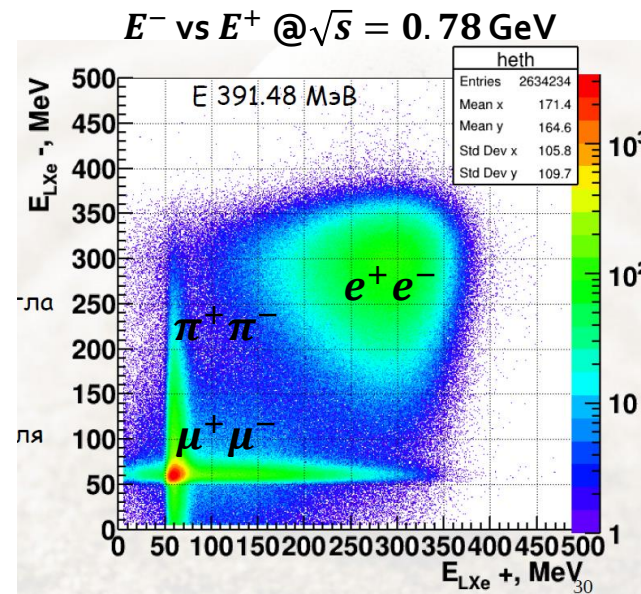
Example of  $e^+e^- \rightarrow \pi^+\pi^-$  event

Similar events:  $e^+e^- \rightarrow \mu^+\mu^-$ ,  $e^+e^- \rightarrow e^+e^-$

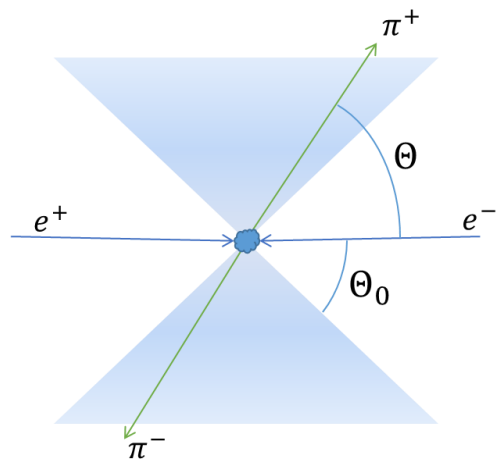
Unique feature of CMD-3:  
three independent methods to measure  $N_{\pi\pi}/N_{ee}$ !

- Energy deposition distribution
- Momentum distribution
- Angular distribution

Agree to 0.2%!



# Measurement of polar angle

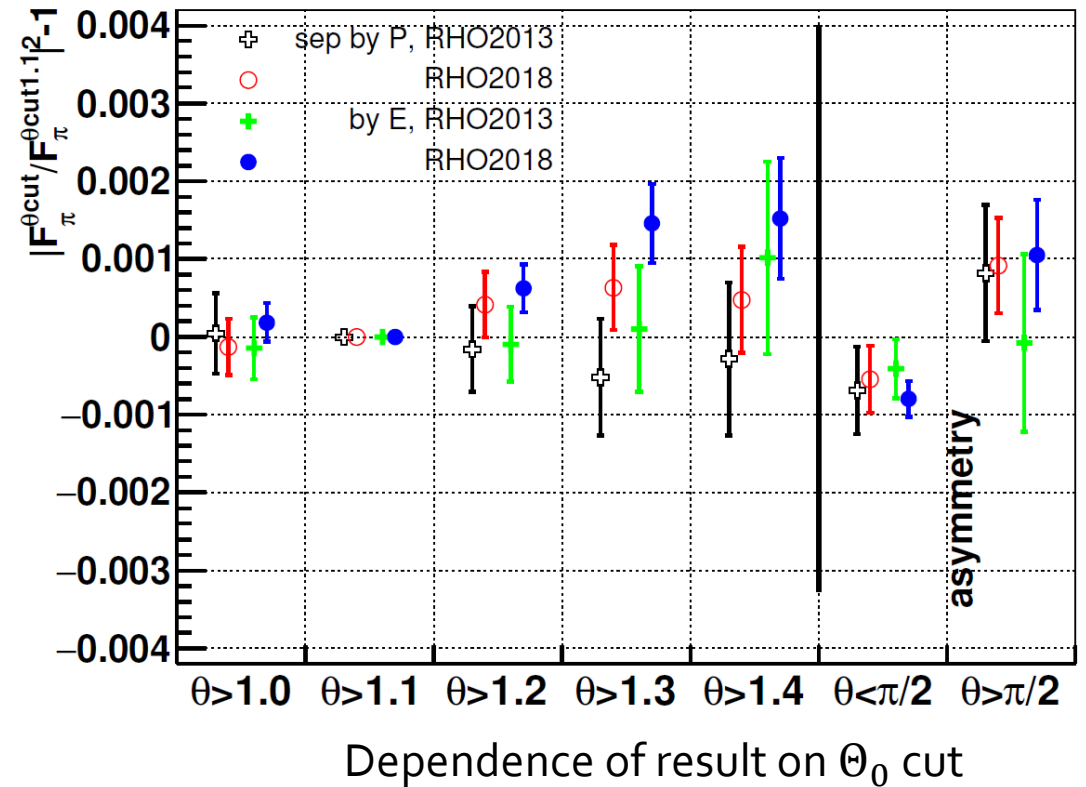


$\Theta$  angle is measured by drift chamber via charge division

Two detector systems with strips readout, LXe calorimeter and Z-chamber, are used for precise calibration and monitoring of DC

We need to precisely know the fiducial volume ( $\Theta_0$  cut).

$$|F_\pi|^2 = \left( \frac{N_{\pi\pi}}{N_{ee}} - \Delta_{bg} \right) \cdot \frac{\sigma_{ee}^0 \cdot (1 + \delta_{ee}) \cdot \epsilon_{ee}}{\sigma_{\pi\pi}^0 \cdot (1 + \delta_{\pi\pi}) \cdot \epsilon_{\pi\pi}}$$

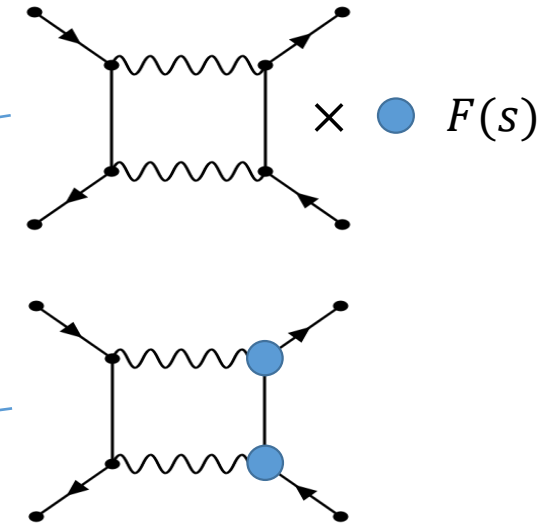
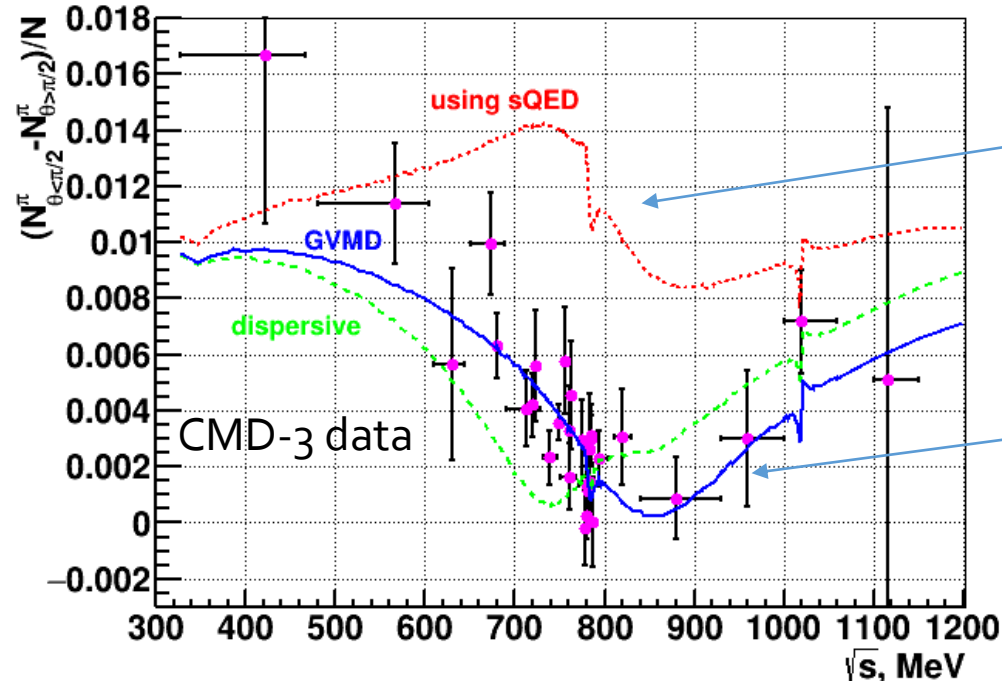
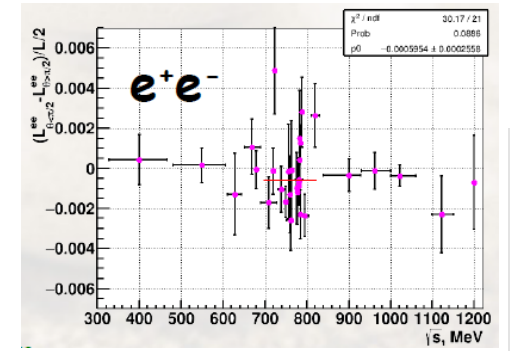


Factor 10 smaller compared to CMD-2, SND2k!

# Charge asymmetry in $e^+e^- \rightarrow \pi^+\pi^-$

Charge asymmetry in  $e^+e^- \rightarrow \pi^+\pi^-$  is due to interference between ISR/FSR and between one- and two-photon exchange

$$A = (N_{\Theta < \pi/2}^{\pi} - N_{\Theta > \pi/2}^{\pi})/N$$



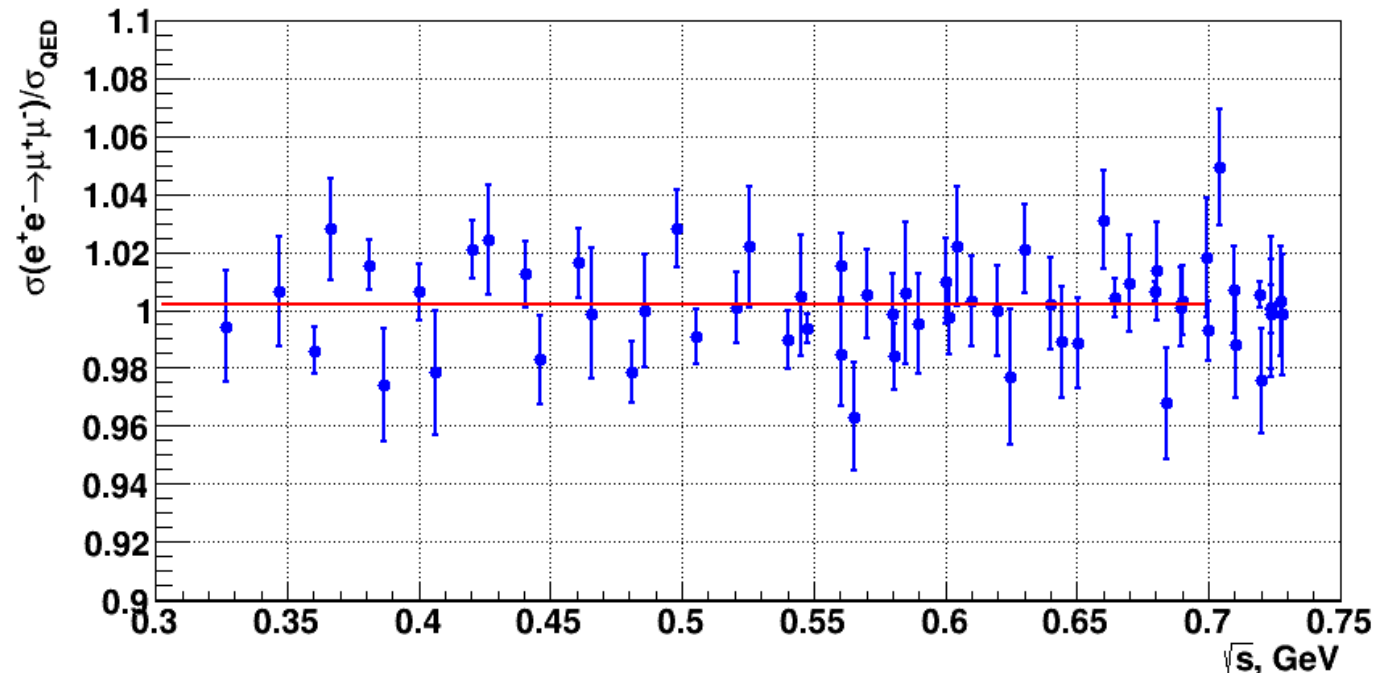
The theoretical model by Lee, Ignatov, PLB 833 (2022) 137283 (GVDM) describes well the CMD-3 data  
 Recent calculation in dispersive formalism Colangelo et al., JHEP 08 (2022) 295 confirms the effect.



# Measurement of $e^+e^- \rightarrow \mu^+\mu^-$

$e^+e^- \rightarrow \mu^+\mu^-$  events are identified as a by-product of analysis, which allows to measure  $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$  and compare it to QED prediction

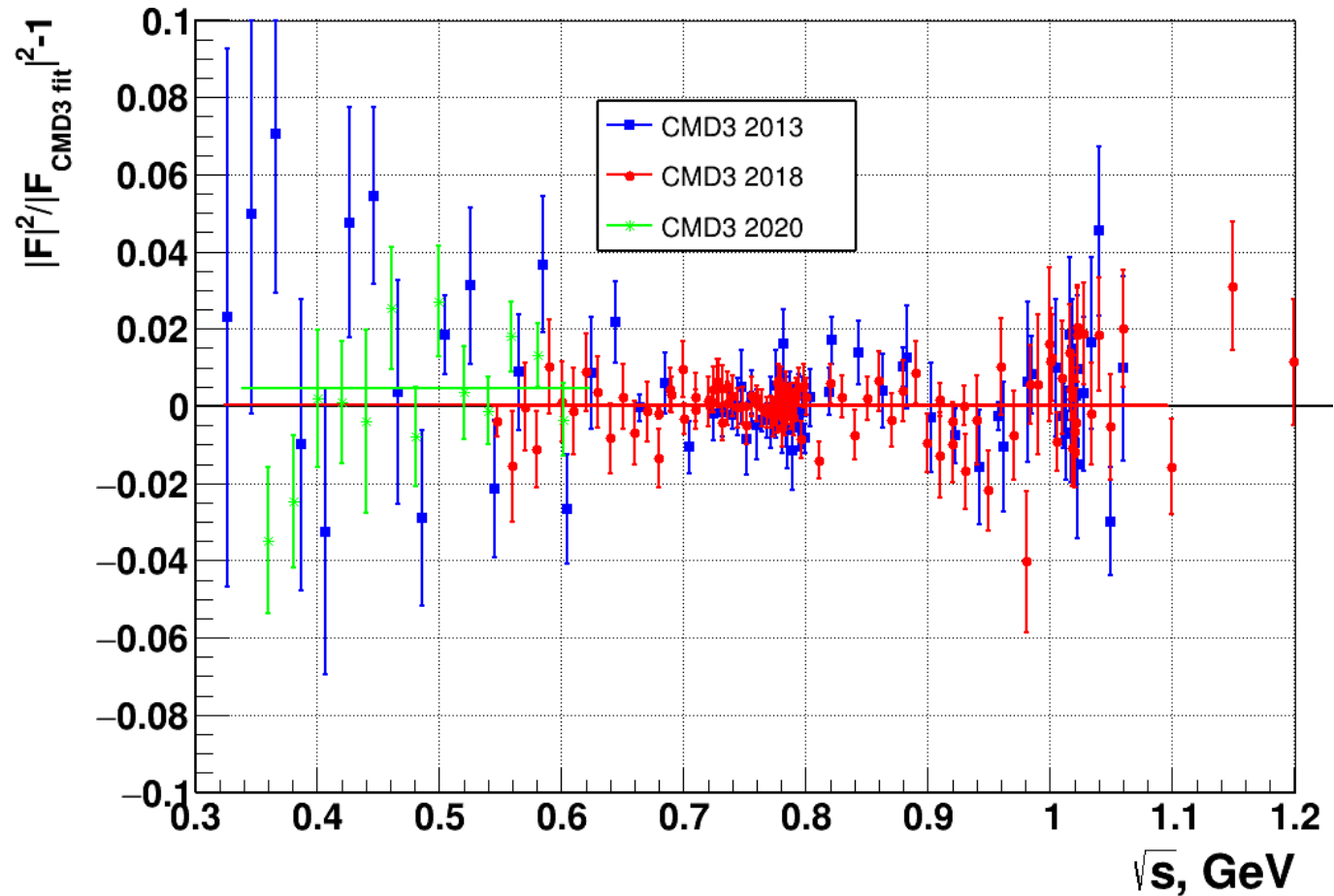
$$\sigma(e^+e^- \rightarrow \mu^+\mu^-)_{CMD3} / \sigma(e^+e^- \rightarrow \mu^+\mu^-)_{QED}$$



**+0.17 ± 0.16 %**

Powerful cross-check of  $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$  measurement! All ingredients are tested: event separation, detection efficiencies, radiative corrections.

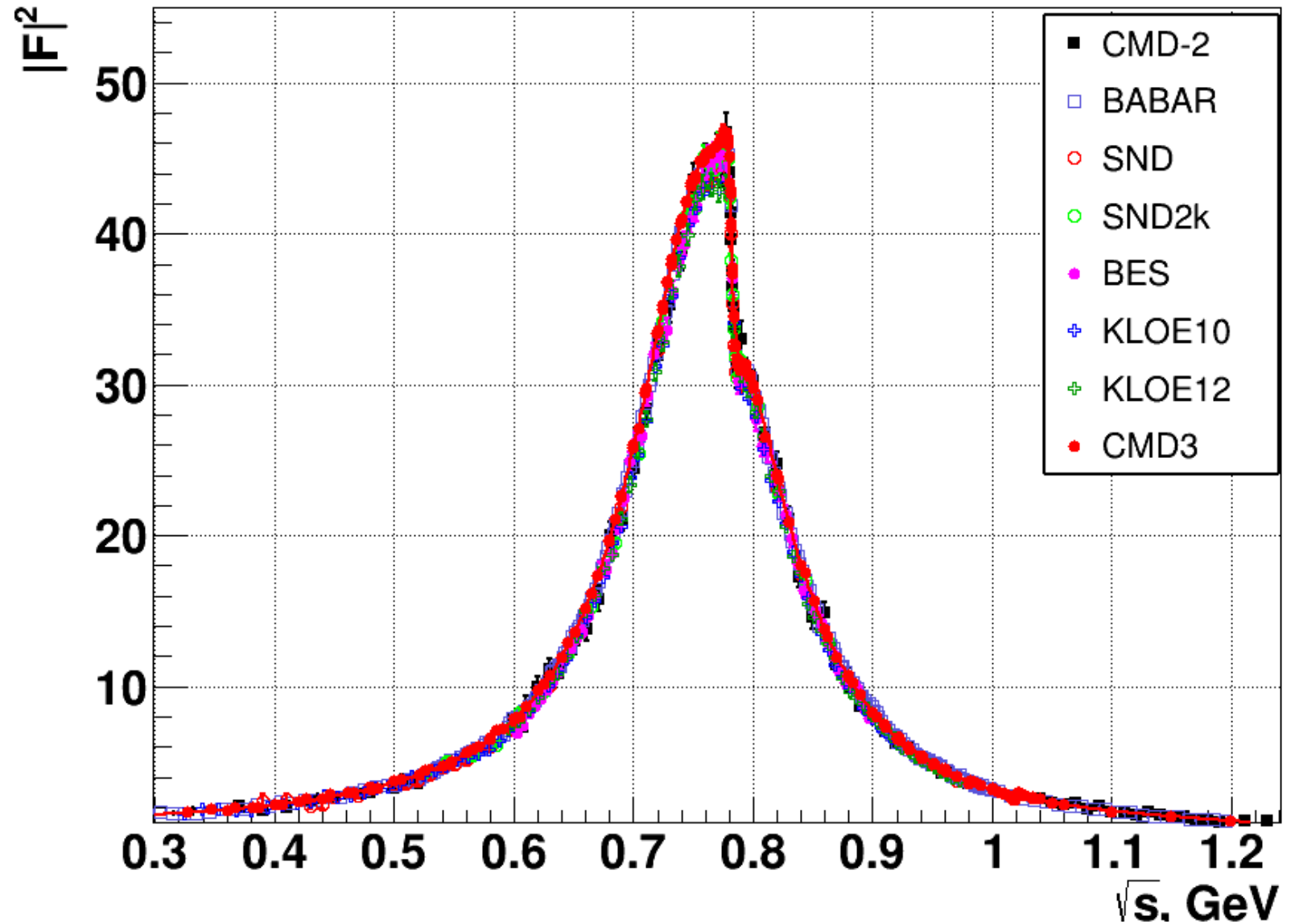
# Comparison of data taking seasons



Results based on 2013, 2018 and 2020 data only agree to  $\sim 0.1\%$ !  
The detector performance and run conditions were significantly different for these runs.

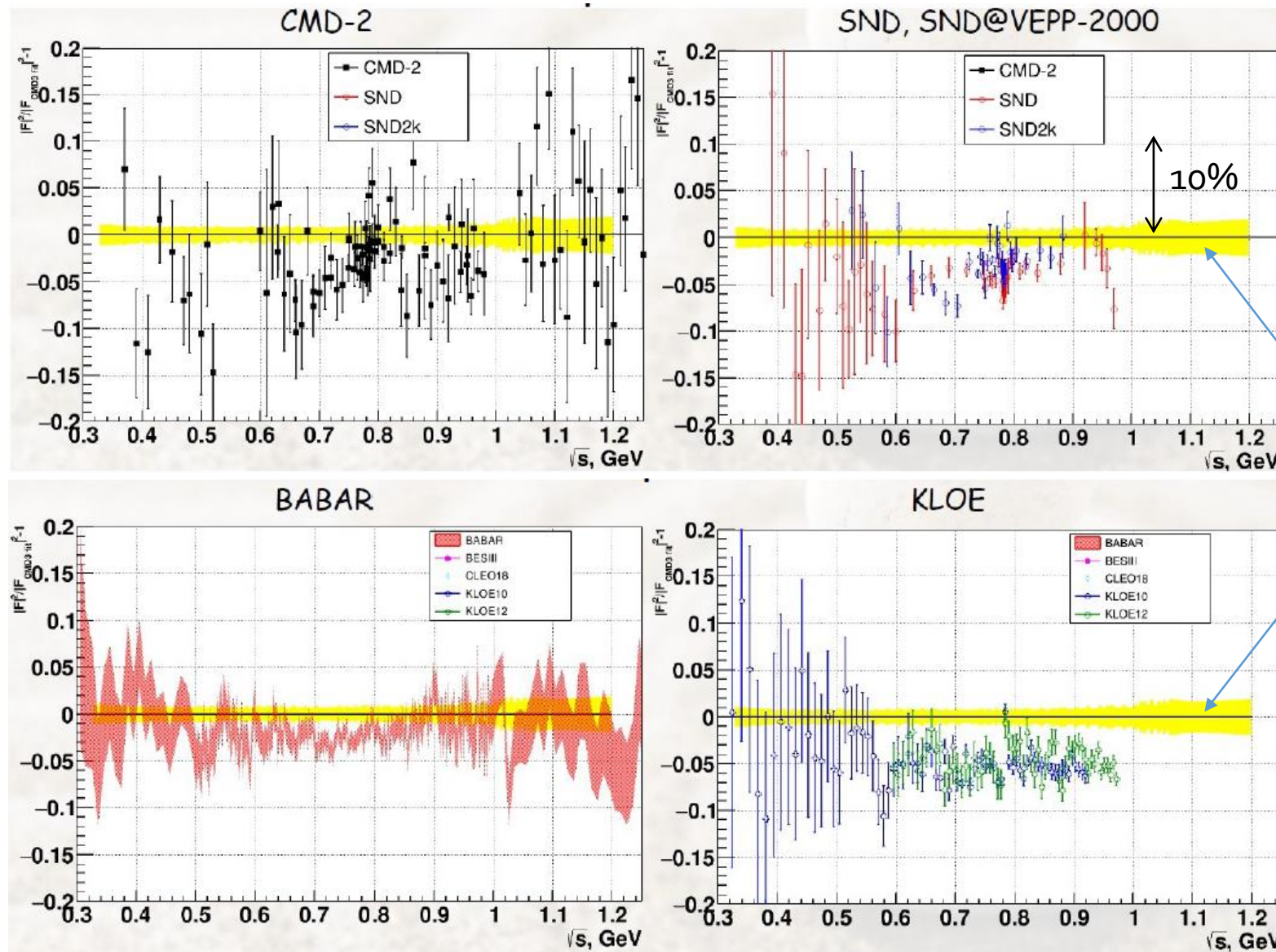
# Comparison to other measurements

At first glance, they look close to each other...



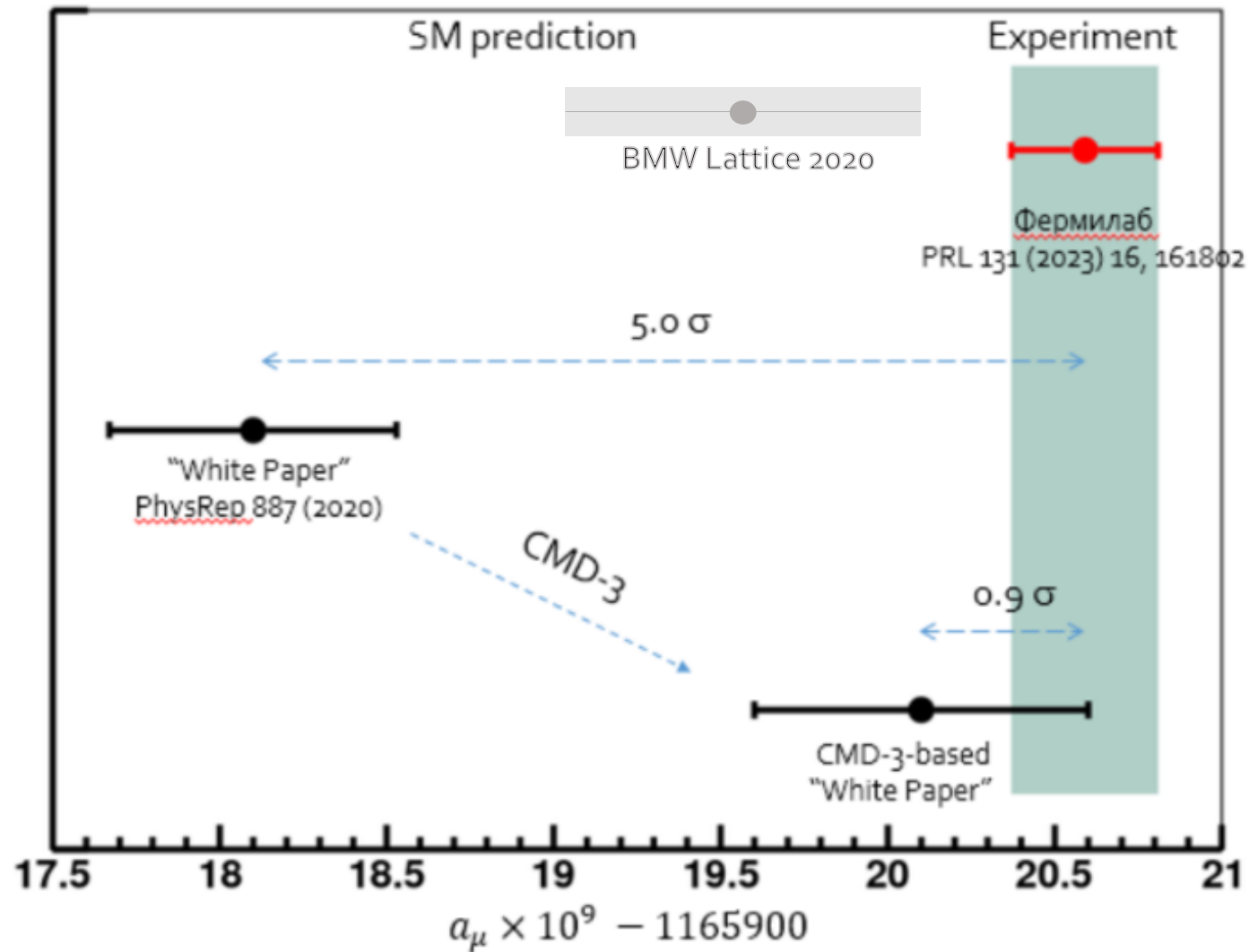
CMD-3 is systematically above previous measurements by ~2-5%

# Comparison to other measurements



CMD-3

# Experiment vs SM prediction



# Результат

- На детекторе КМД-3 измерено сечение  $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$  в области энергий от 0.32 до 1.2 ГэВ в системе центра масс
  - Лучшая статистическая точность в мире
  - Наиболее детальный анализ систематических ошибок, уникальные методы перекрестных проверок
  - «Побочные» измерения: зарядовая асимметрия в  $e^+e^- \rightarrow \pi^+\pi^-$ , сечение  $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ , параметры векторных мезонов,...
  - >10 лет работы
- Результат КМД-3 привел к пересмотру устоявшегося мнения о наличии противоречия между измеренной величиной аномального магнитного момента мюона и предсказанием Стандартной модели
- Результат КМД-3 вызвал большой резонанс в сообществе физики элементарных частиц
  - Проведены рабочие совещания, посвященные результату и детальной проверке анализа данных
  - Ведутся новые независимые измерения/обработки данных, которые должны подтвердить/опровергнуть результат КМД-3
- На ВЭПП-2000 мы планируем провести новый цикл измерений с целью повысить точность в 2-3 раза

Публикации (направлены в PRL/PRD):

- 1. F.V.Ignatov et al. (CMD-3 Collaboration) Measurement of the pion formfactor with CMD-3 detector and its implication to the hadronic contribution to muon (g-2) // arXiv:2309.12910 [hep-ex]
- 2. F.V.Ignatov et al. (CMD-3 Collaboration) Measurement of the  $e^+e^- \rightarrow \pi^+\pi^-$  cross section from threshold to 1.2 GeV with the CMD-3 detector // arXiv:2302.08834 [hep-ex]