

Baryon form factors at BESIII

Kai Zhu (IHEP, Beijing)

On behalf of BESIII collaboration

International Workshop on e^+e^- collisions from Phi to Psi

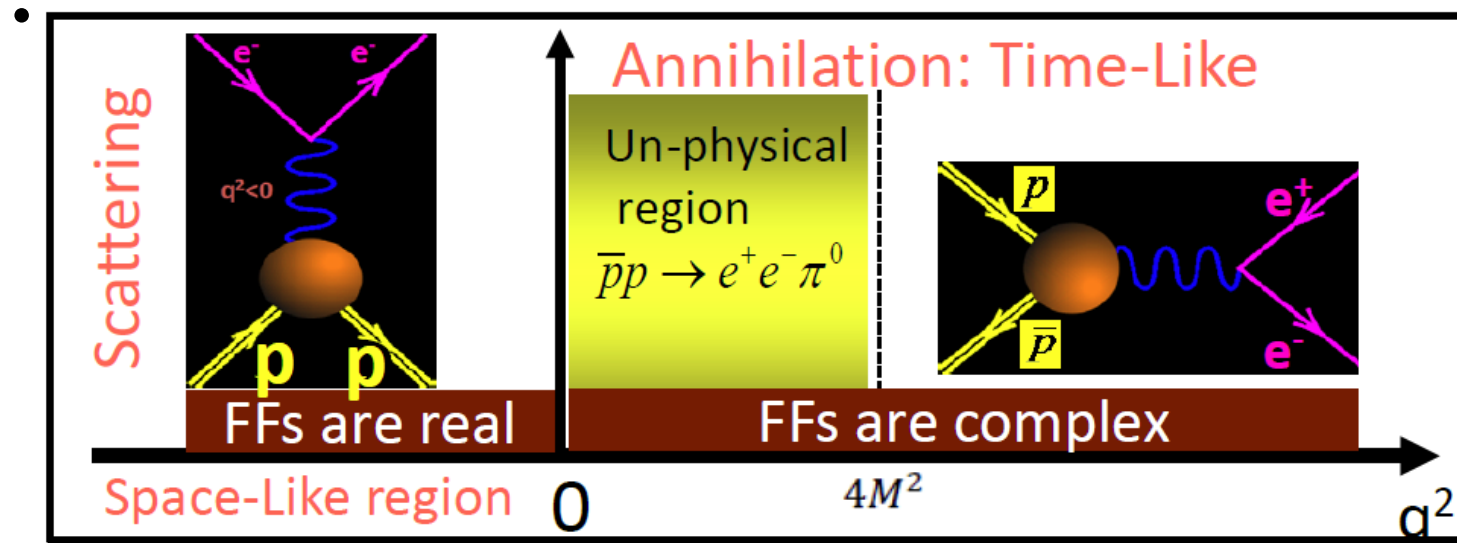
25 Feb. 2019 – 1 Mar. 2019, Budker INP, Novosibirsk

Outline

- Introduction (form factors and BESIII data)
- Time-like form factors of **proton**
 - ISR un-tag method with data at higher energies
 - Scan data
- Time-like form factors of Λ
- Time-like form factors of Λ_c
- Summary and outlook

Form factors of proton

- Understanding the inner structure of proton

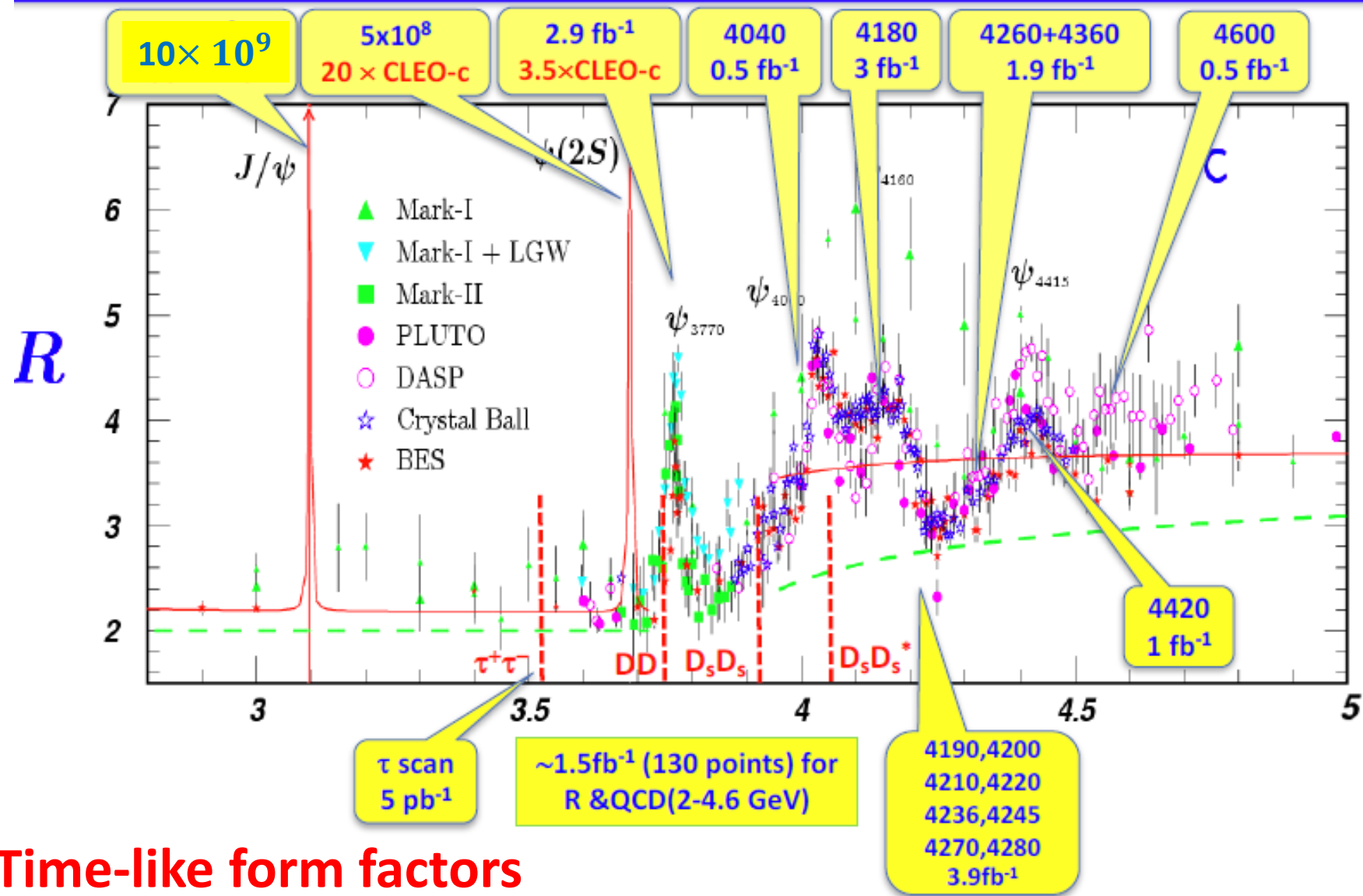


- Form factors represent the charge distribution in momentum space and are related to the cross sections directly
 - Probe the size of the nucleus
 - Test QCD scaling

Form factors for proton

- Space-like
 - Many measurements via scattering
 - At JLab, the ratio was measured precisely with an uncertainty of $\sim 1\%$, based on which the proton electric and magnetic radii could be extracted
- Time-like
 - Measurements via $p\bar{p}$ collision or $e^+e^- \rightarrow p\bar{p}$
 - $\sigma(e^+e^- \rightarrow p\bar{p})$ $\frac{d\sigma_{p\bar{p}}(s)}{d\Omega} = \frac{\alpha^2\beta C}{4s} [|G_M|^2(1 + \cos^2\theta_p) + \frac{4m_p^2}{s}|G_E|^2\sin^2\theta_p]$
 - Lacking accurate data on the angular distributions. Most experiments assume $|G_E| = |G_M| = |G_{eff}|$
 - Only two experiments measure $|G_E/G_M|$, with inconsistent results (Babar & PS170)

BESIII data: the largest e^+e^- collision samples in τ -charm region



Proton form factors at BESIII

Both scan and ISR can be used at BESIII

- High-energy data sets (3.773 ~ 4.6 GeV)
 - Untagged ISR technique
 - Tagged ISR technique
- Low-energy scan data sets (2.0 ~ 3.08 GeV)
 - Phys. Rev. D. 91, 112004 (2015), the first set
 - Full set; expected to supersede the previous results

Proton form factors at BESIII

Both scan and ISR can be used at BESIII

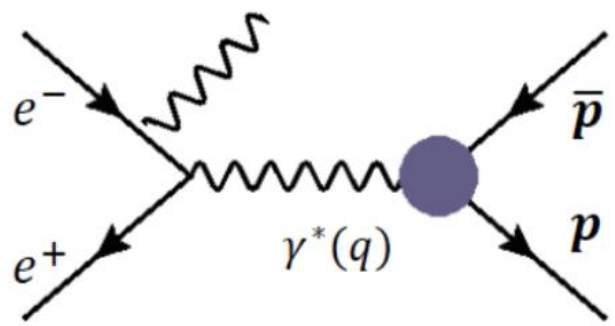
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*This
talk*

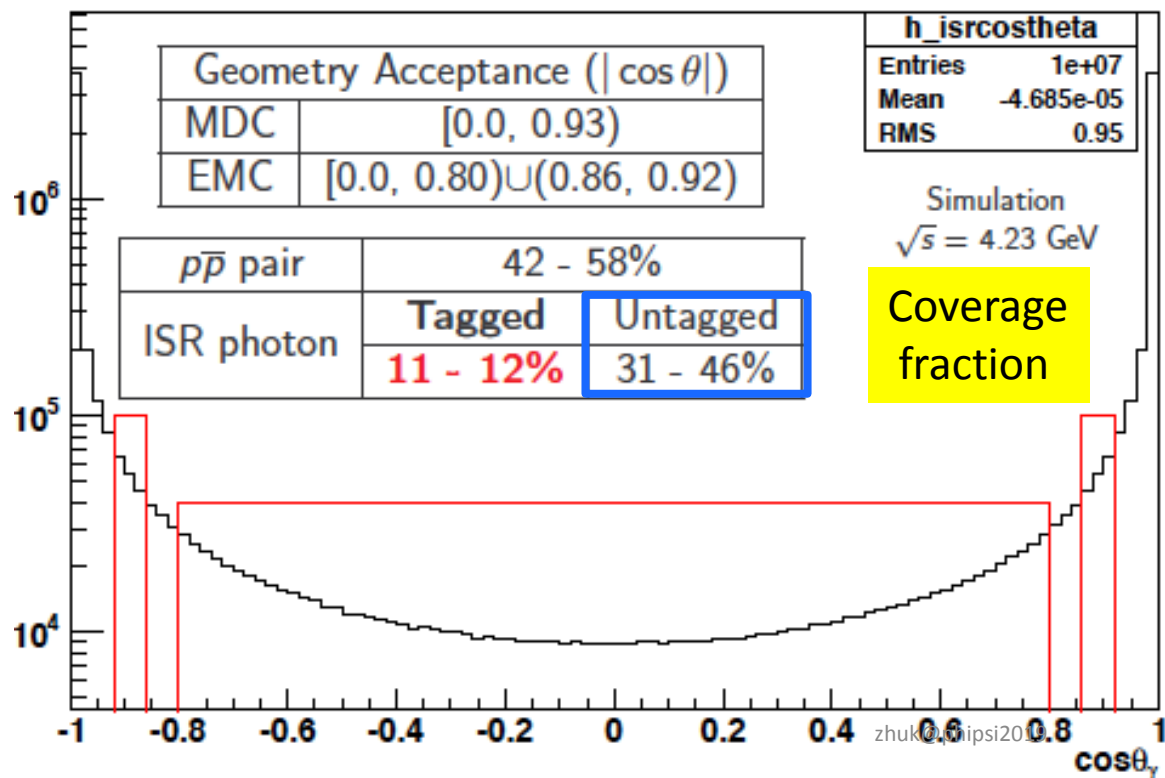
Untagged ISR technique

Proton form factors results at BESIII with untag ISR method

arXiv:1902.00665 [hep-ex]



E_{cm} [GeV]	L [pb^{-1}]
3.773	2931.8
4.008	481.96
4.226	1053.9
4.258	825.67
4.358	539.84
4.416	1041.3
4.600	585.4



Total: 7.4 fb⁻¹

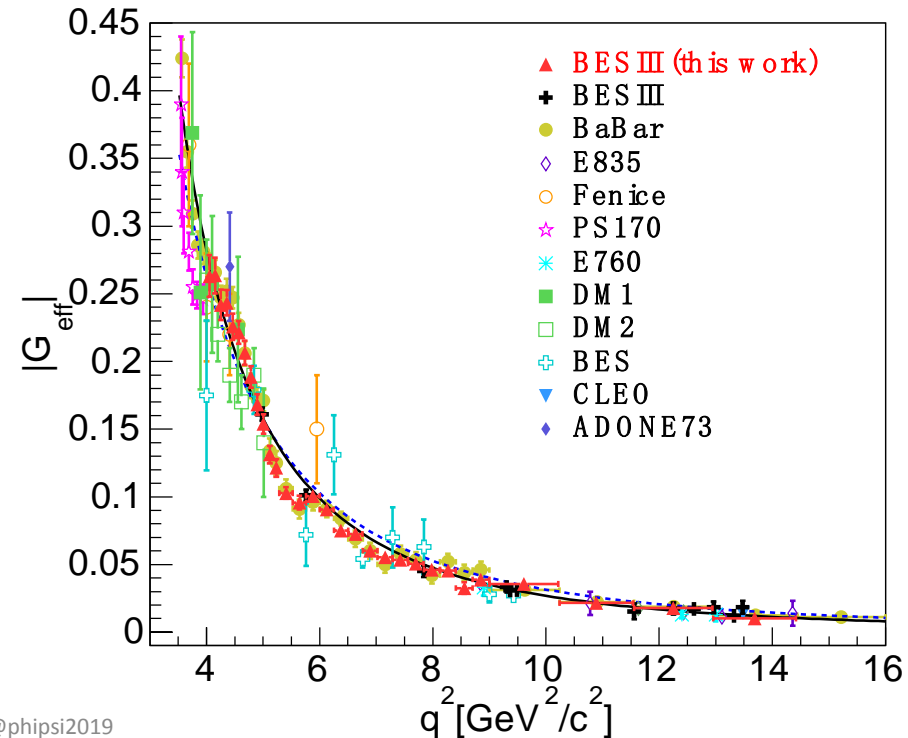
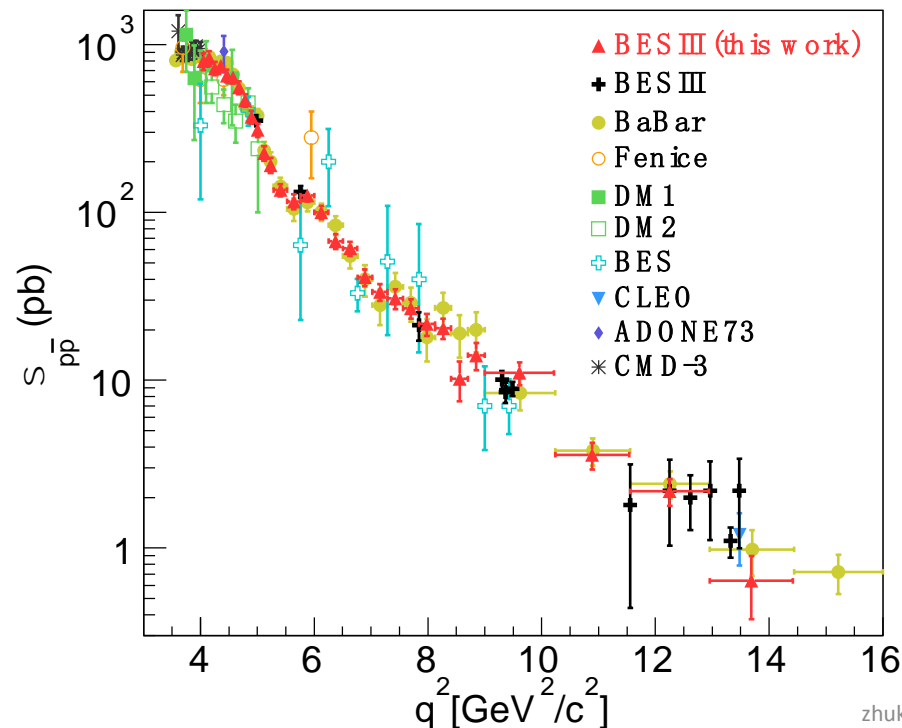
- **ISR analysis:** continuous q^2 range
- **Untagged ISR analysis:** high statistic

Total cross section and effective FFs [arXiv:1902.00665 \[hep-ex\]](https://arxiv.org/abs/1902.00665) in 30 intervals of the $M_{p\bar{p}}$ between 2.0 and 3.8 GeV

The ISR and the Born cross sections are related by the radiator function $W(s,x)$:

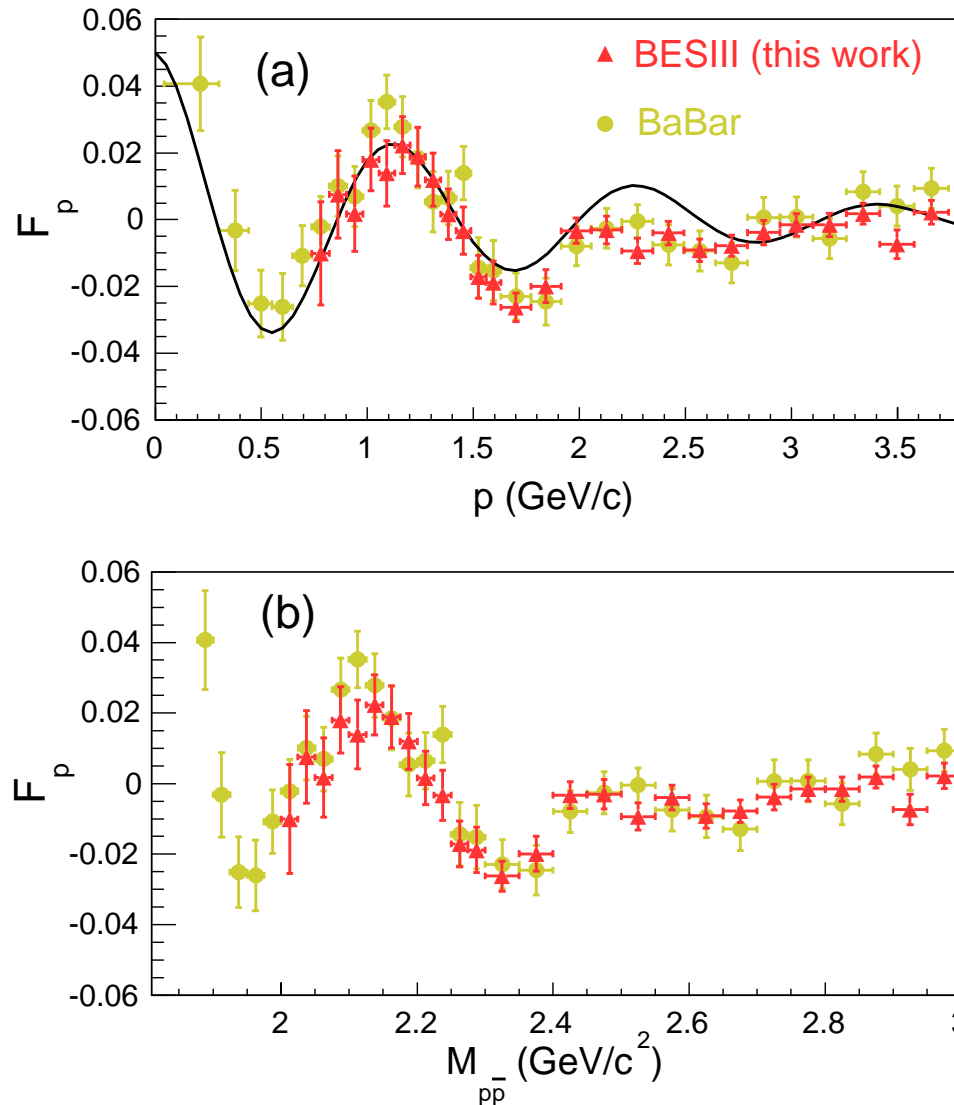
$$\frac{d\sigma_{p\bar{p}\gamma_{ISR}}(q^2)}{dq^2} = \frac{1}{s} W(s,x) \sigma_{p\bar{p}}(q^2),$$

$$x = 1 - \frac{q^2}{s}, q^2 = M_{p\bar{p}}^2$$



Oscillations

arXiv:1902.00665 [hep-ex]



- The oscillations can be extracted as
$$F_\rho = |G_{eff}| - F^0$$
(F^0 : regular behavior over the long range)

Confirmed observation by BaBar
Phys. Rev. Lett. 114, 232301 (2015)
Phys. Rev. C 93, 035201 (2016)

Re-scattering effect?
Resonance contribution?

R= |GE|/|GM| in 3 intervals of $M_{p\bar{p}}$ between 2.0 and 3.0 GeV

arXiv:1902.00665 [hep-ex]

$$\frac{d\sigma}{d\cos\theta_p} = A(H_M + \frac{R^2}{\tau} H_E)$$

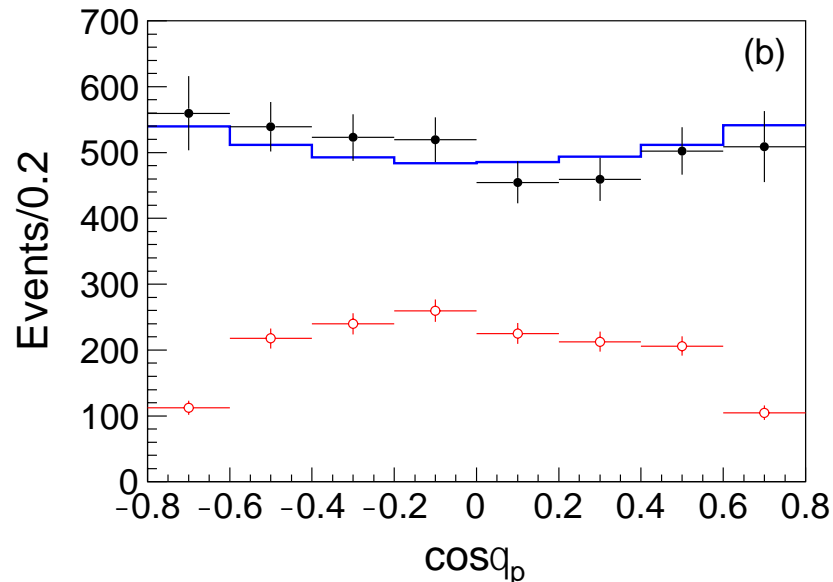
$H_M(\cos\theta_p, M_{p\bar{p}})$ from MC ($G_E=0$)
 $H_E(\cos\theta_p, M_{p\bar{p}})$ from MC ($G_M = 0$)

Data of 7 samples are combined

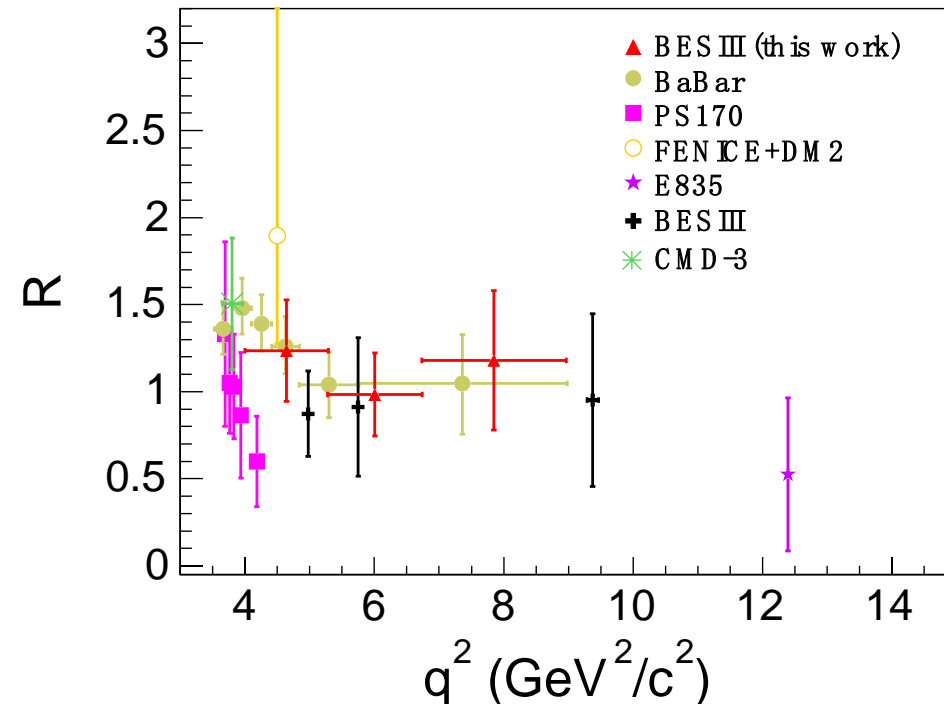
After bkg. subtraction (red)

After efficiency correction (black)

Fit (blue)



$M_{p\bar{p}}$ [GeV/ c^2]	Fitting range ($\cos\theta_p$)	R
2.0 - 2.3	[-0.6,0.6]	1.24 ± 0.29
2.3 - 2.6	[-0.8,0.8]	0.98 ± 0.24
2.6 - 3.0	[-0.8,0.8]	1.18 ± 0.40



(Full) low-energy scan data sets
688.5 pb⁻¹, 22 energies, 2.0 to 3.08 GeV

Cross section of $e^+e^- \rightarrow p\bar{p}$ and effective FF

■ Strategy:

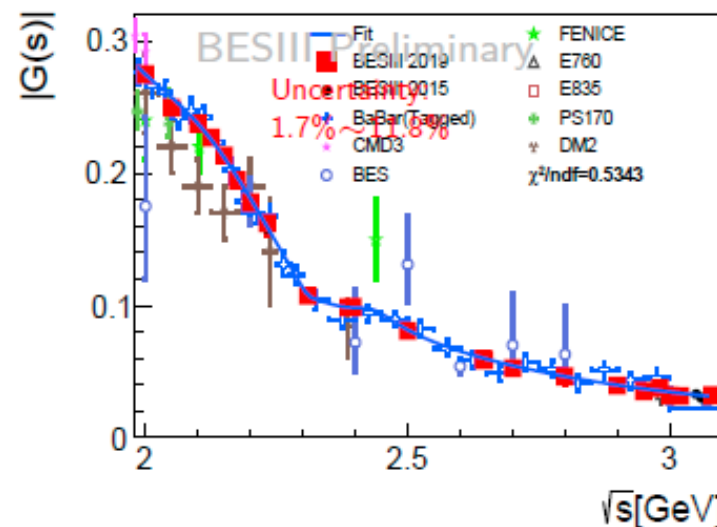
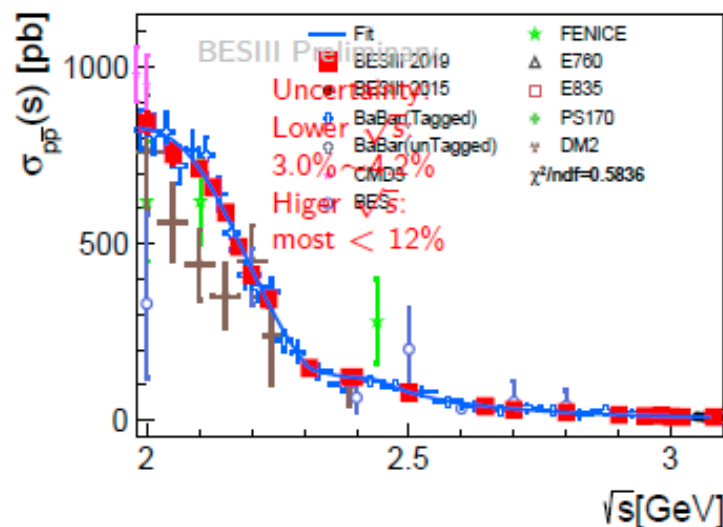
- In TL region:

$$\frac{d\sigma_{p\bar{p}}(s)}{d\Omega} = \frac{\alpha^2 \beta C}{4s} [|G_M|^2 (1 + \cos^2 \theta_p) + \frac{4m_p^2}{s} |G_E|^2 \sin^2 \theta_p]$$

- Assume $|G| = |G_E| = |G_M|$, the effective FF is

$$|G| = \sqrt{\frac{\sigma_{p\bar{p}}}{\frac{4\pi\alpha^2\beta C}{3s} \left(1 + \frac{2m_p^2}{s}\right)}}$$

- In time-like region, BESIII result is an unprecedented accuracy.



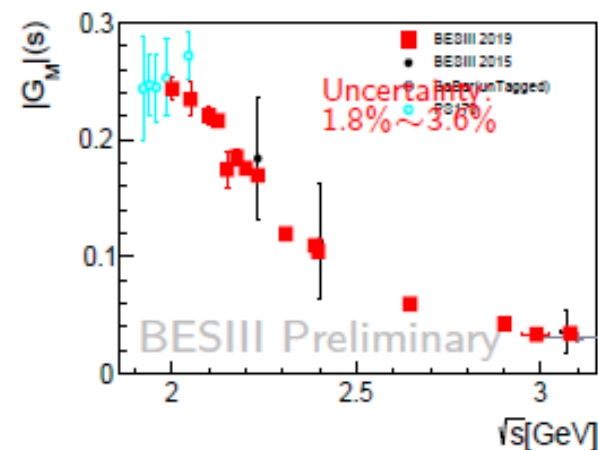
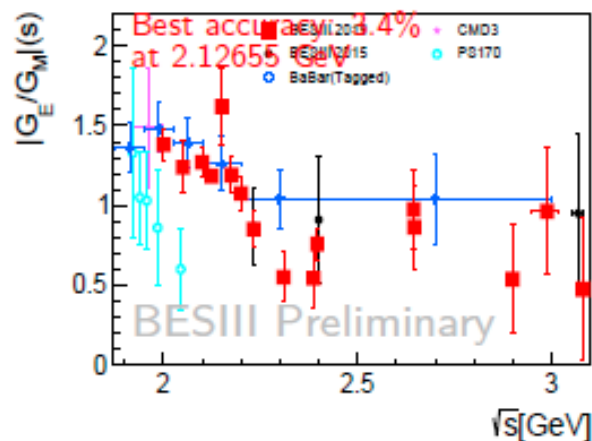
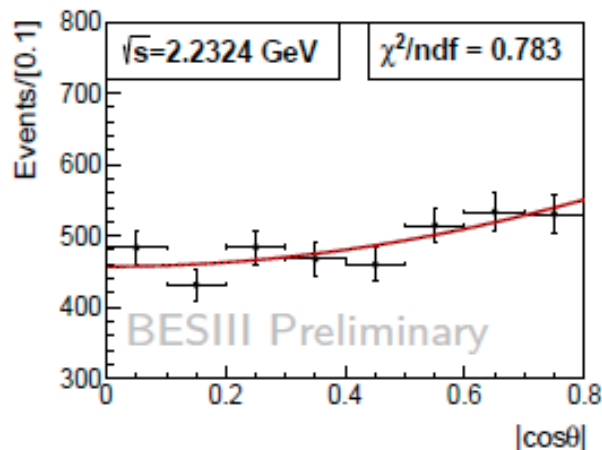
Measurement of $|G_E/G_M|$ and magnetic FF

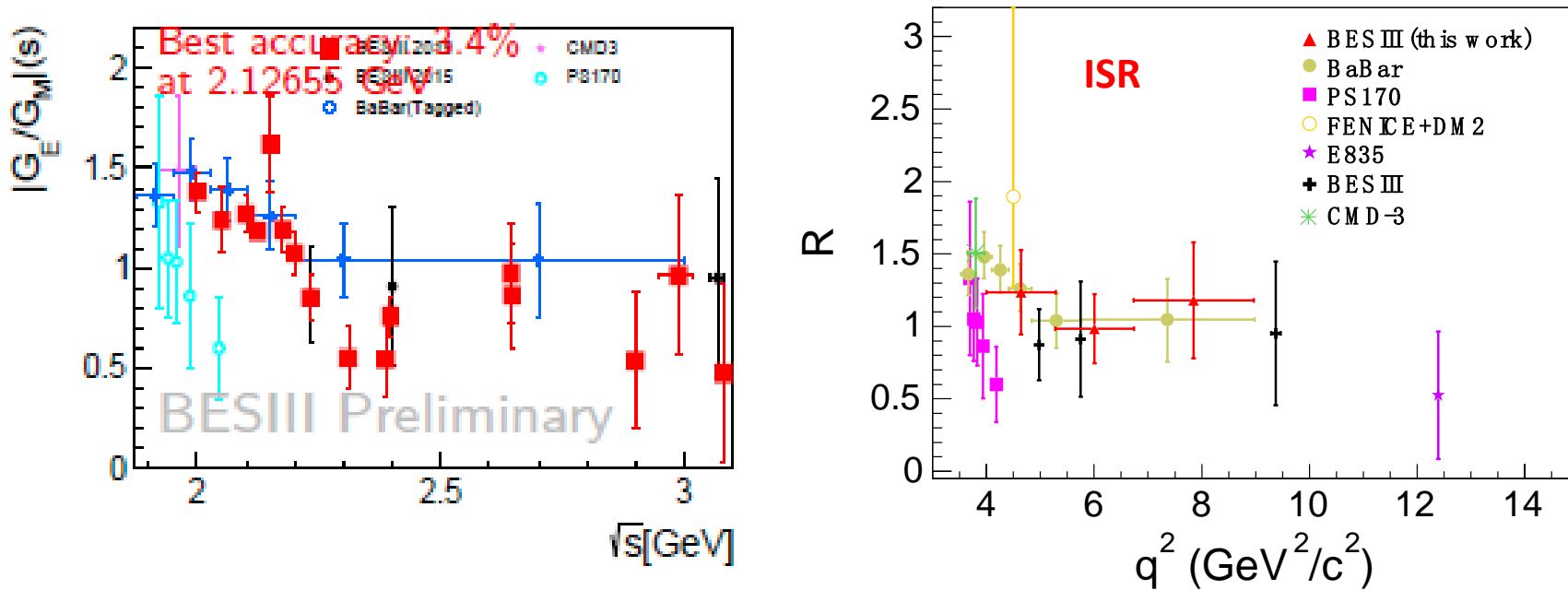
■ Strategy:

- Fit on the polar angular distribution of proton:

$$\frac{dN}{\epsilon(1+\delta) \times d \cos \theta_p} = \frac{\mathcal{L} \hbar c \pi \alpha^2 \beta C}{2s} |G_M|^2 [(1 + \cos^2 \theta_p) + \frac{4m_p^2}{s} |\frac{G_E}{G_M}|^2 (1 - \cos^2 \theta_p)]$$

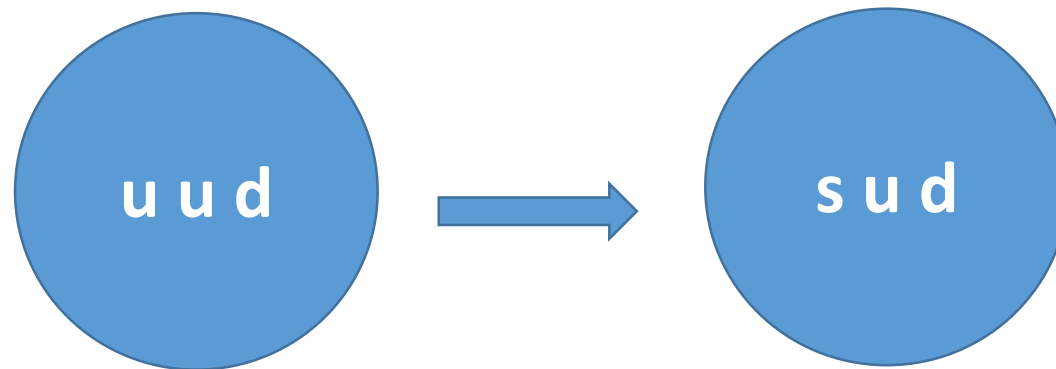
- Uncertainties include statistical and systematical.





- $|G_E/G_M|$ providing an uncertainty comparable to the space-like region for **the first time**.
- BESIII measurement of $|G_E/G_M|$ strongly favors BaBar's compared to that of PS170.

Λ form factors



Hyperon structure

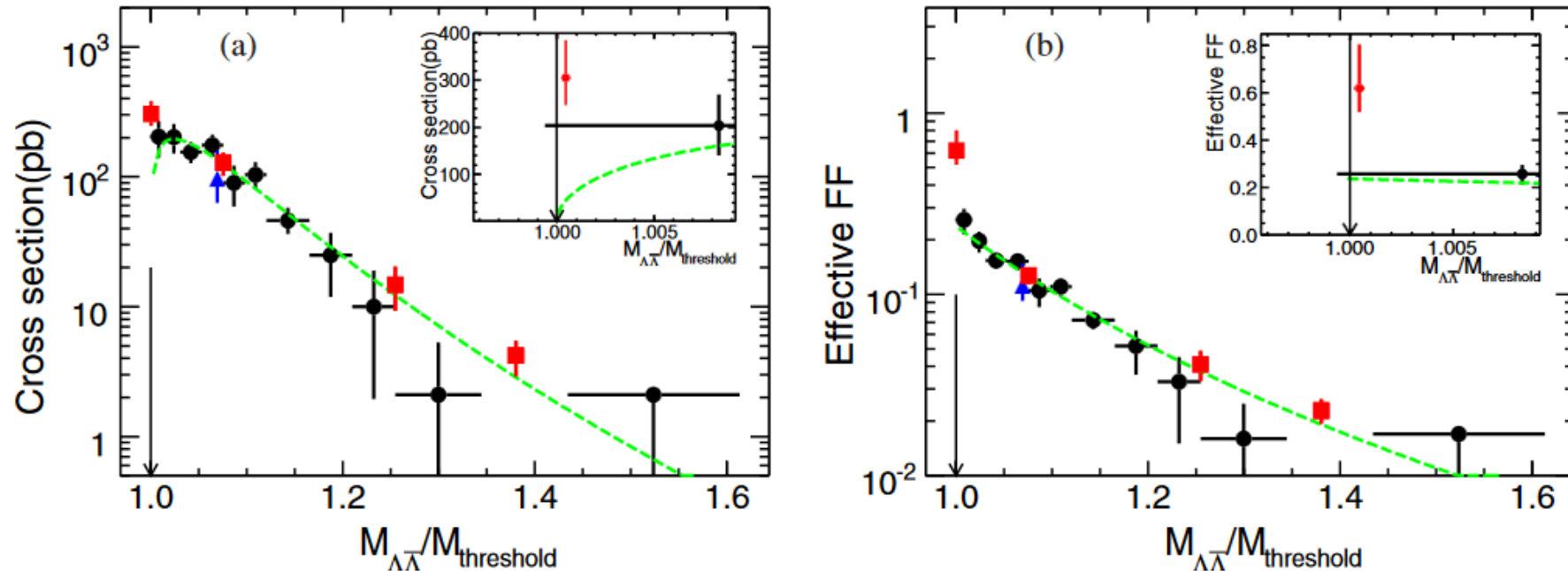
- ❑ Electromagnetic Form Factors (EMFFs)
 - fundamental hadron structure observables
 - describe the deviation from the point-like case
 - related to the charge- and magnetization density
- ❑ EMFFs of nucleon can be studied in:
 - elastic scattering, $e^- N \rightarrow e^- N$, space-like
 - annihilation, $e^+ e^- \rightarrow N \bar{N}$, $N \bar{N} \rightarrow e^+ e^-$, time-like
- ❑ Hyperons are difficult to study in the space-like region
 - they are unstable — hyperon targets are unfeasible
 - the quality of hyperon beams is in general not sufficient
- ❑ $e^+ e^-$ annihilation offers the best opportunity to study hyperon structure

Previous BESIII measurements

PHYSICAL REVIEW D **97**, 032013 (2018)

Observation of a cross-section enhancement near mass threshold in $e^+e^- \rightarrow \Lambda\bar{\Lambda}$

$$\sqrt{s} = 2.2324, 2.400, 2.800 \text{ and } 3.080 \text{ GeV}$$



Results of the cross section and effective EMFFs

- The cross section $\sigma = \frac{N_{\text{signal}}}{L\epsilon(1+\delta)Br(\Lambda \rightarrow p\pi^-)Br(\bar{\Lambda} \rightarrow \bar{p}\pi^+)}$
- ISR and vacuum polarization factor $1 + \delta$ is from ConExc
 - ϵ is the detection efficiency, L is the luminosity
 - $\sigma = 119.0 \pm 5.3(\text{stat.}) \pm 7.3(\text{sys.}) \text{ pb}^1$ $\sqrt{s} = 2.396 \text{ GeV}$

- Effective form factors are related to σ , $|G(q^2)| = \sqrt{\frac{\sigma(q^2)}{(1+\frac{1}{2\tau})(\frac{4\pi\alpha^2\beta}{3q^2})}}$
- $|G| = 0.123 \pm 0.003(\text{stat.}) \pm 0.004(\text{sys.})$

$\alpha \approx \frac{1}{137}$ is the fine structure constant,

$\beta = \sqrt{1 - \frac{1}{\tau}}$ is the velocity, $\tau = \frac{q^2}{4m_\Lambda^2}$.

Previous measurements

	$\sigma(\text{pb})$	$ G $	Reference
BESIII $\sqrt{s} = 2.40 \text{ GeV}$	$128 \pm 19 \pm 18$	$0.127 \pm 0.009 \pm 0.009$	Phys. Rev. D 97 , 032013 (2018)
BaBar $\sqrt{s} = 2.35\text{-}2.40 \text{ GeV}$	176 ± 34	0.152 ± 0.016	Phys. Rev. D 76 , 092006 (2007)

¹The systematic uncertainty is dominated by a conservative estimate of the contribution of $|\vec{p}|(\Lambda)$

Joint decay distribution

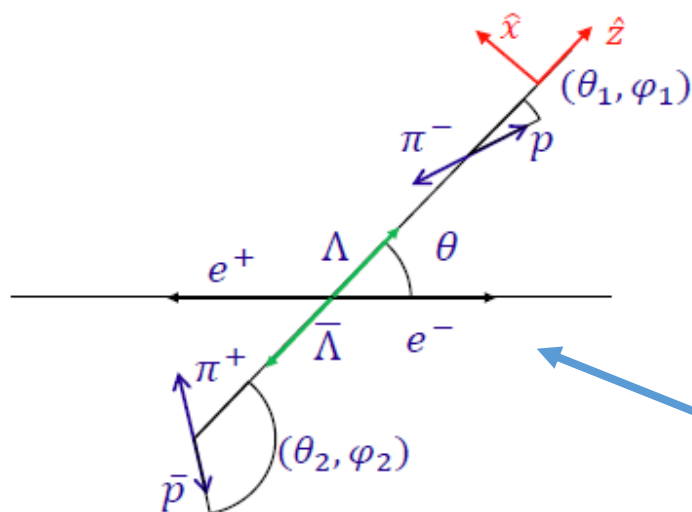
□ Space-like region:

- $e^- B \rightarrow e^- B$ scattering
- $q^2 = (p_{ie} - p_{fe})^2 < 0$
- G_E and G_M real numbers

□ Time-like region:

- $e^+ e^- \leftrightarrow B \bar{B}$, $q^2 \geq 4M_B^2 > 0$
- $G_E(q^2) = |G_E(q^2)|e^{i\Phi_E}$, $G_M(q^2) = |G_M(q^2)|e^{i\Phi_M}$
- Relative phase: $\Delta\Phi = \Phi_E - \Phi_M$

$$\begin{aligned} \mathcal{W}(\xi) = & \mathcal{T}_0 + \eta \mathcal{T}_5 \\ & - \alpha_\Lambda^2 \left(\mathcal{T}_1 + \sqrt{1 - \eta^2} \cos(\Delta\Phi) \mathcal{T}_2 + \eta \mathcal{T}_6 \right) \\ & + \alpha_\Lambda \sqrt{1 - \eta^2} \sin(\Delta\Phi) (\mathcal{T}_3 - \mathcal{T}_4), \end{aligned}$$

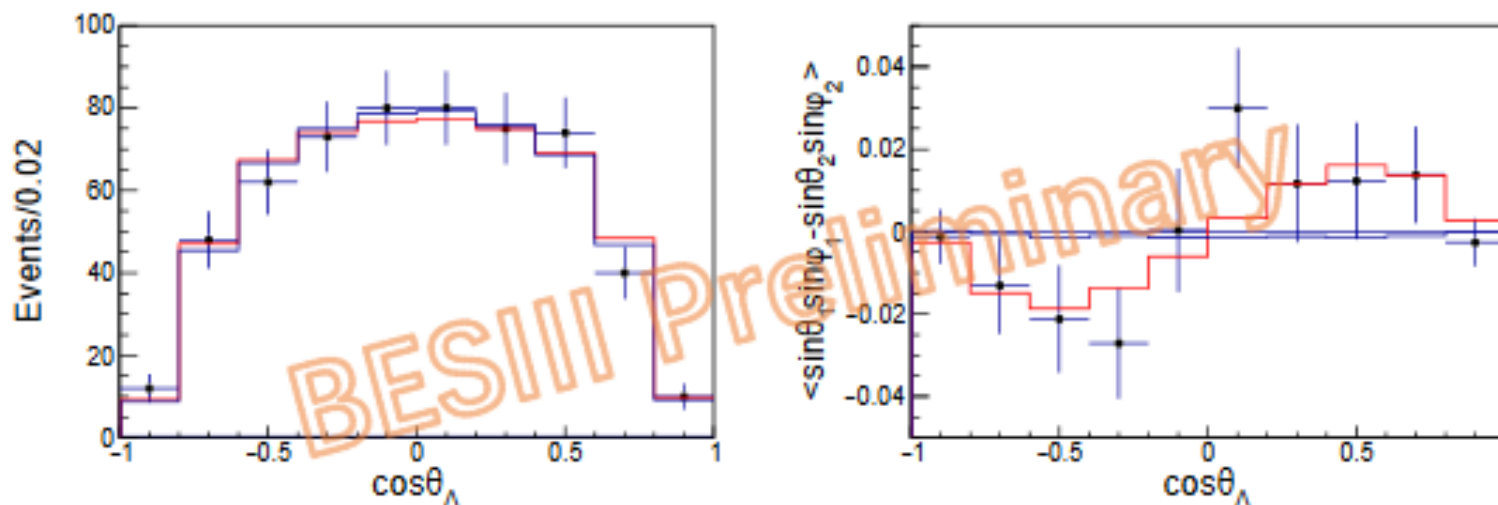


$$\begin{aligned} \mathcal{T}_0(\xi) &= 1, \\ \mathcal{T}_1(\xi) &= \sin^2 \theta \sin \theta_1 \sin \theta_2 \cos \phi_1 \cos \phi_2 + \cos^2 \theta \cos \theta_1 \cos \theta_2, \\ \mathcal{T}_2(\xi) &= \sin \theta \cos \theta (\sin \theta_1 \cos \theta_2 \cos \phi_1 + \cos \theta_1 \sin \theta_2 \cos \phi_2), \\ \mathcal{T}_3(\xi) &= \sin \theta \cos \theta \sin \theta_1 \sin \phi_1, \\ \mathcal{T}_4(\xi) &= \sin \theta \cos \theta \sin \theta_2 \sin \phi_2, \\ \mathcal{T}_5(\xi) &= \cos^2 \theta, \\ \mathcal{T}_6(\xi) &= \cos \theta_1 \cos \theta_2 - \sin^2 \theta \sin \theta_1 \sin \theta_2 \sin \phi_1 \sin \phi_2. \end{aligned}$$

$|G_E/G_M|$ and relative phase

Fit by maximal likelihood

α_Λ is determined to be 0.75 ± 0.01
by preliminary BESIII result

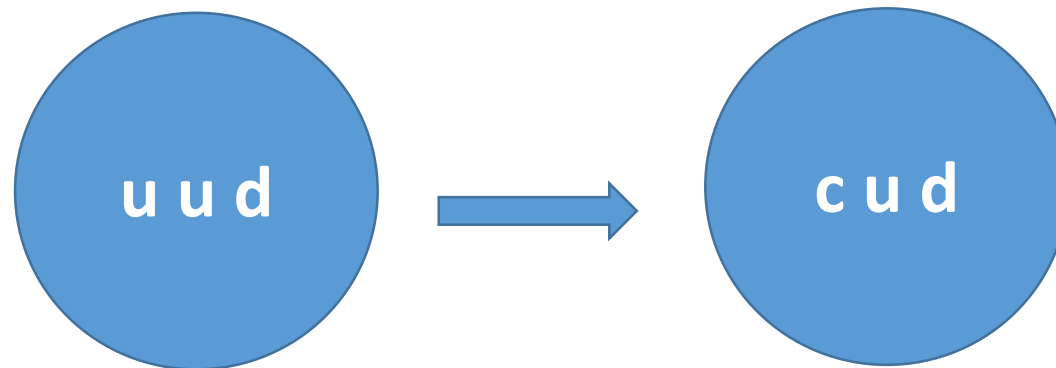


$$R = 0.94 \pm 0.16(stat.) \pm 0.03(sys.) \pm 0.02(\alpha_\Lambda)$$

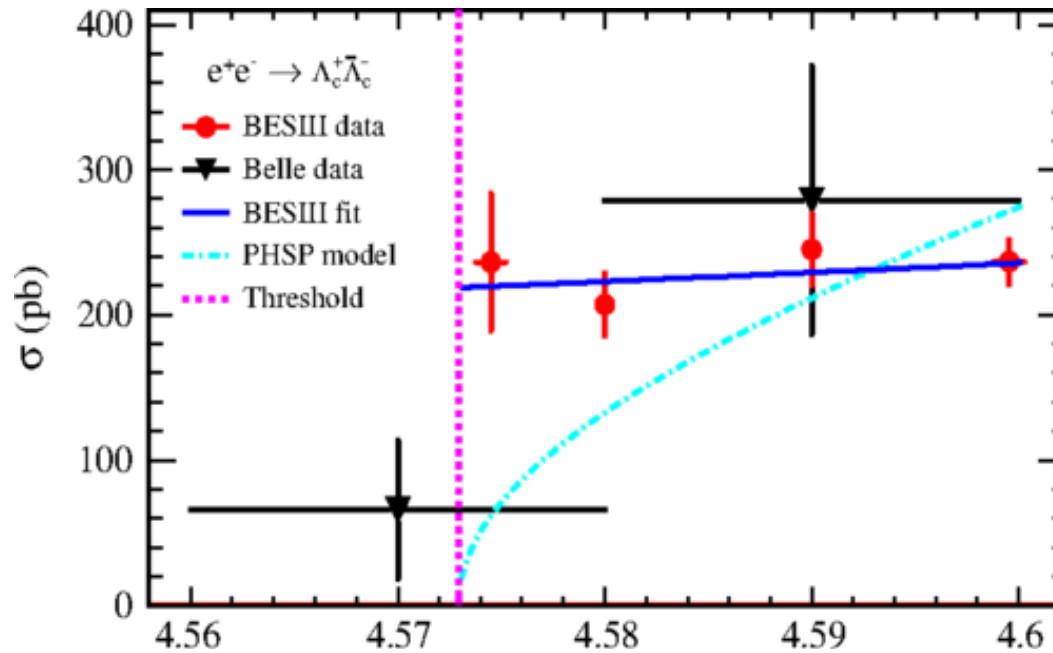
$$\Delta\Phi = 42^\circ \pm 16^\circ(stat.) \pm 8^\circ(sys.) \pm 6^\circ(\alpha_\Lambda)$$

The statistical significance of $\Delta\Phi$ is 4.3σ . **First measurement of $\Delta\phi$**

Λ_c form factors near threshold



$e^+e^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-$ near threshold Phys. Rev. Lett. 120, 132001 (2018)



$\sqrt{s} = 4574.5, 4580.0, 4590.0$
and **4599.5 MeV**

$$|G_E/G_M|:$$

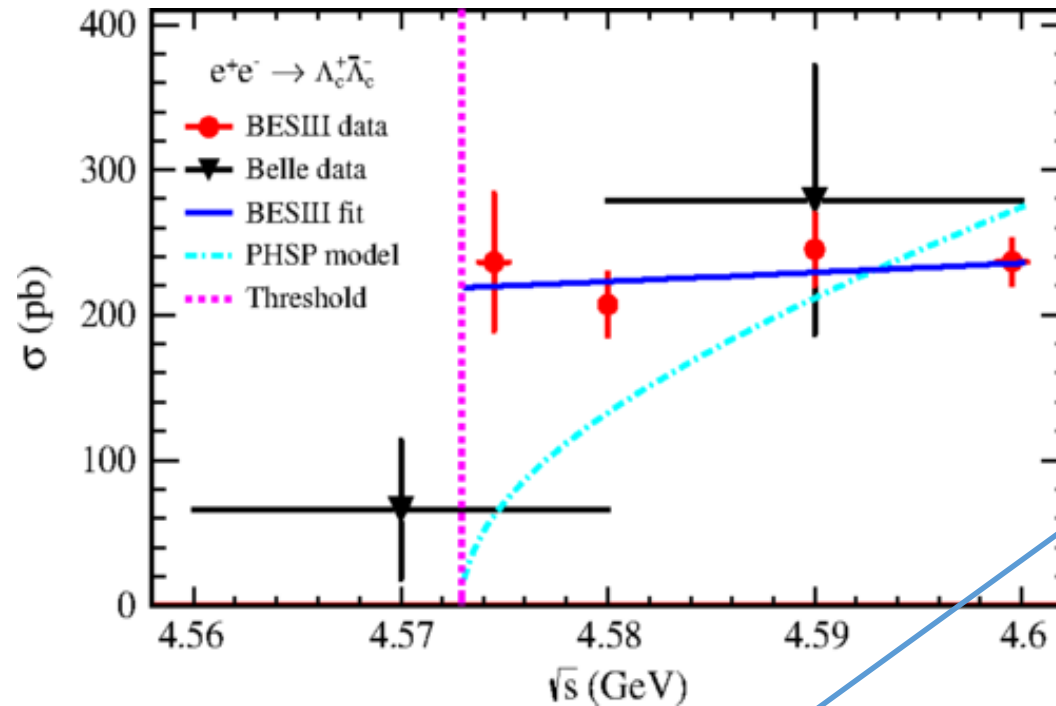
$$1.14 \pm 0.14 \pm 0.07$$

$$1.23 \pm 0.05 \pm 0.03$$

- At thr there is indeed a step in $\sigma(e^+e^- \rightarrow \Lambda_c \Lambda_{c\text{bar}})$,
- Followed by a kind of a plateau
- At thr $\sigma(e^+e^- \rightarrow \Lambda_c \Lambda_{c\text{bar}})$ is close to the pointlike value, once the Coulomb enhancement factor is taken into account:

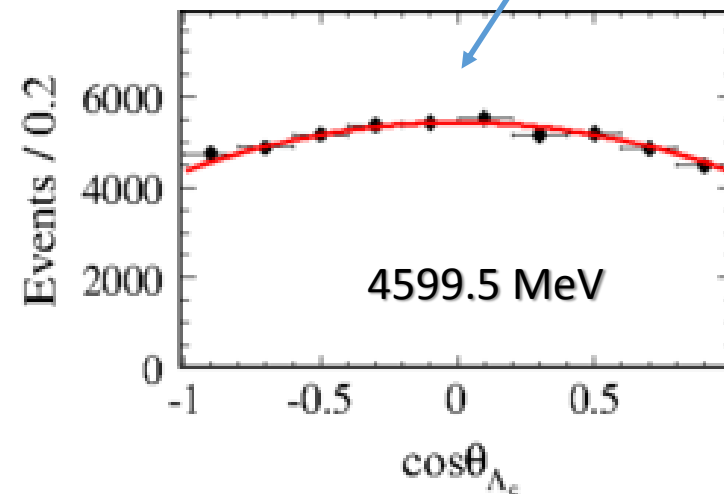
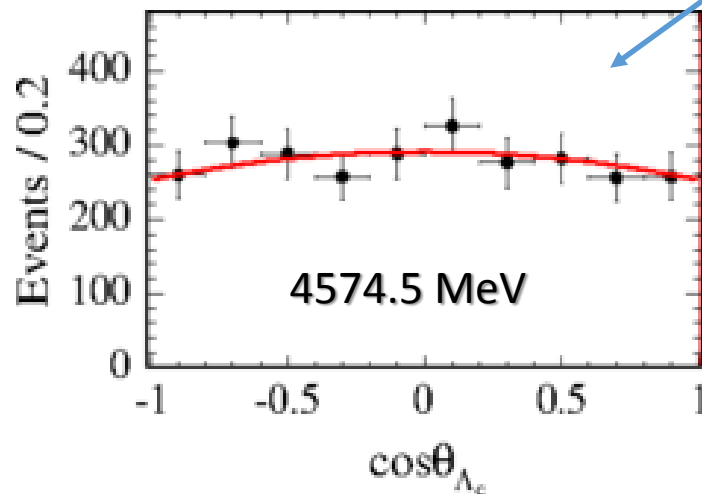
$$\sigma(e^+e^- \rightarrow \Lambda_c \Lambda_{c\text{bar}})_{\text{pointl}} \approx \pi^2 \alpha^3 / (2M_B^2) \approx 145 \text{ pb}$$

$e^+e^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-$ near threshold [Phys. Rev. Lett. 120, 132001 \(2018\)](#)



$\sqrt{s} = 4574.5, 4580.0, 4590.0$
 and **4599.5 MeV**

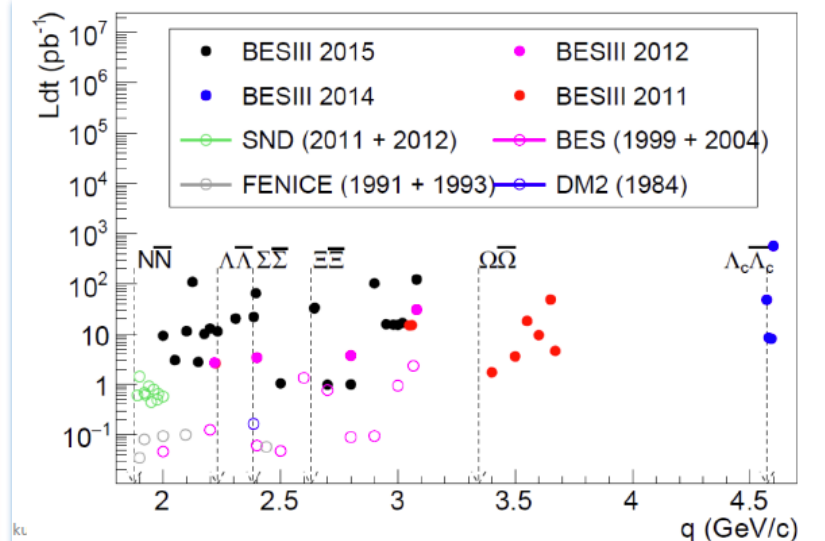
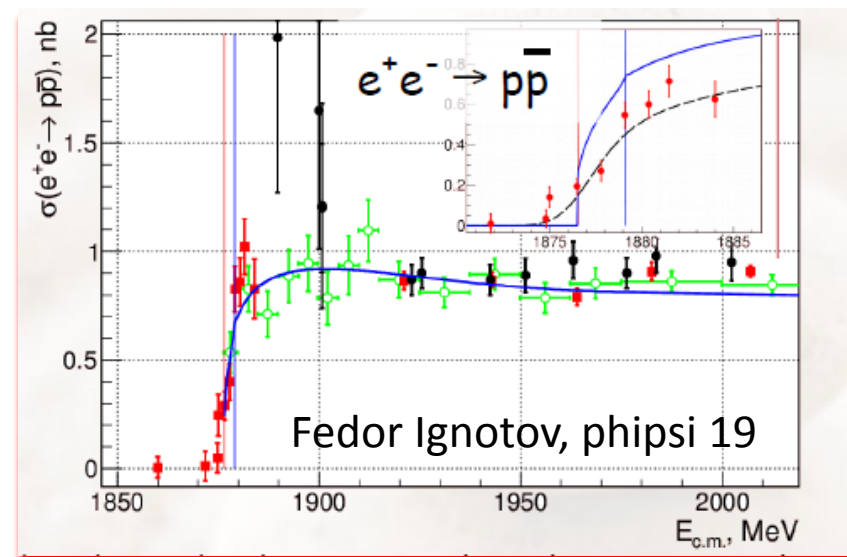
$|G_E/G_M|:$
 $1.14 \pm 0.14 \pm 0.07$
 $1.23 \pm 0.05 \pm 0.03$



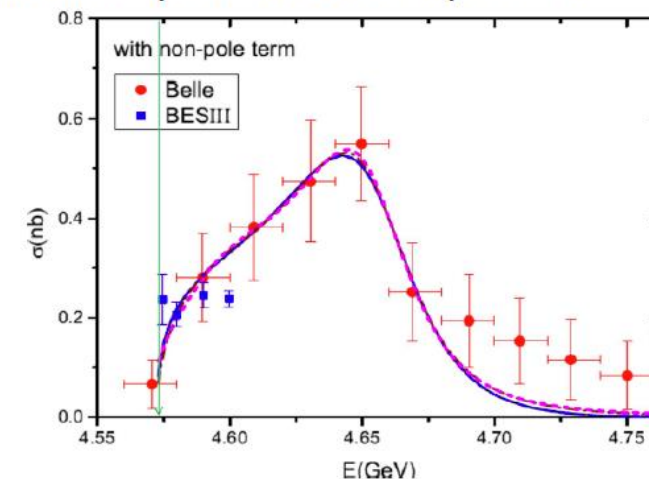
Summary and outlook (I)

- Proton form factors have been measured at BESIII with ISR and scan technique
 - Result in significantly improved precision
 - Some novel/unexpected features observed/confirmed
- Λ form factors with relative phase between G_E and G_M first time
- Λ_c effective form factors and $|G_E/G_M|$ near threshold

Summary and outlook (II)



BESIII data up to 4.9 GeV expected.



- Will come
 - Proton form factors with varied technique ($E_{\text{cm}} < 2 \text{ GeV}$)
 - Form factors of other baryons, such as neutron (will be released soon), Ω , Σ , Ξ , as well as polarizations
 - Λ_c form factors at higher energy



Thanks for your attention!

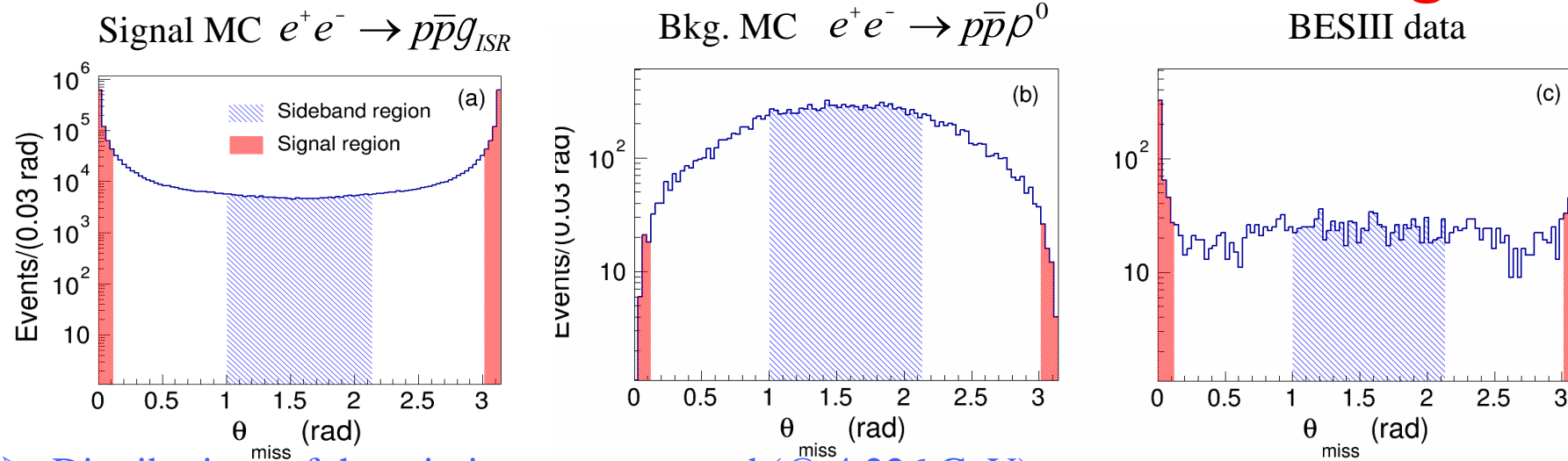
Backup

Background estimation based on Sideband method

- Angular distributions of the missing momentum (@ 4.226 GeV):

Un-tag ISR

BESIII data

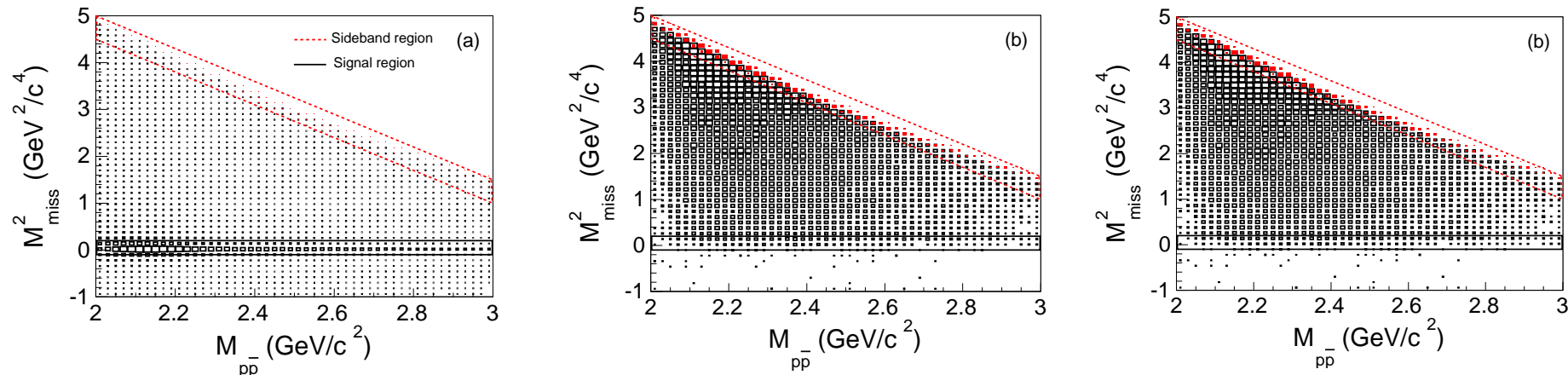


- Distributions of the missing mass squared (@ 4.226 GeV):

Signal MC $e^+e^- \rightarrow p\bar{p}g_{ISR}$

Bkg. MC from two-photon channel

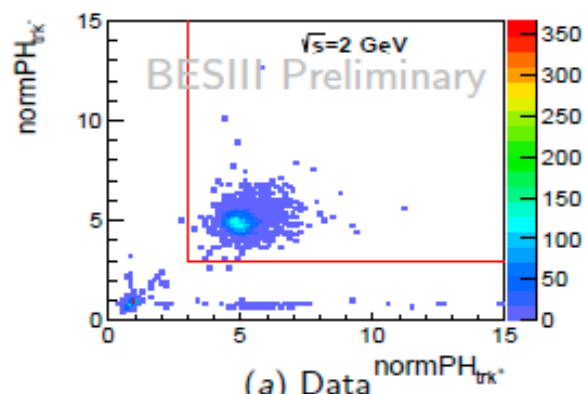
BESIII data



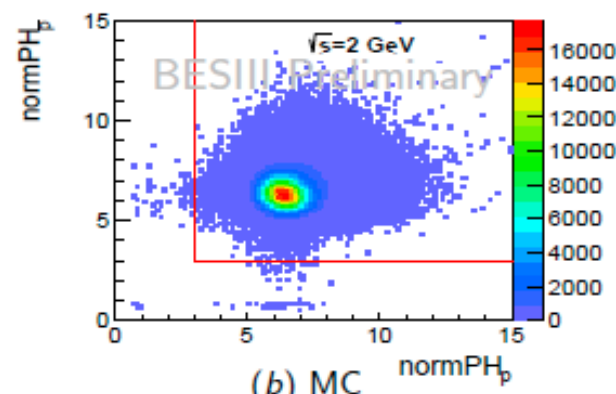
Scan data (selection I)

■ Particle identification

- At (2.0~2.15) GeV, use normalized pulse height,
- At (2.175~3.08) GeV, use dE/dx and TOF.

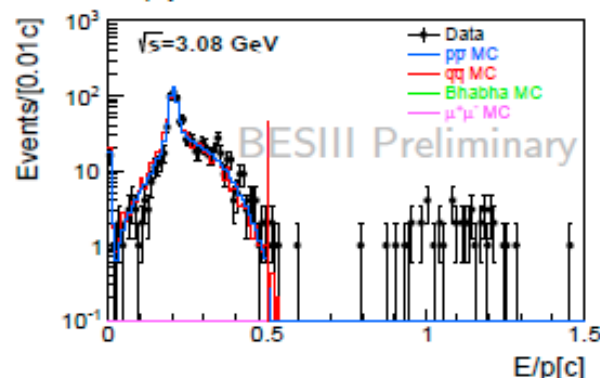


(a) Data

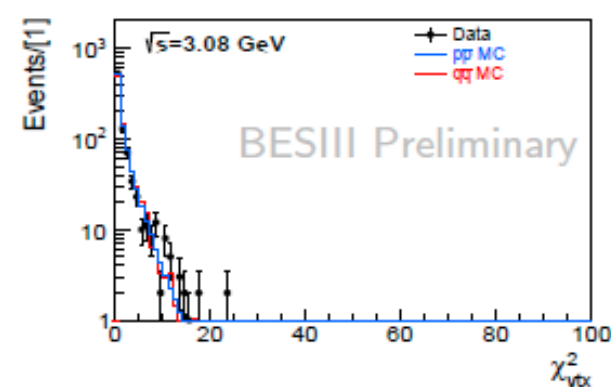


(b) MC

■ Use E/p criteria.



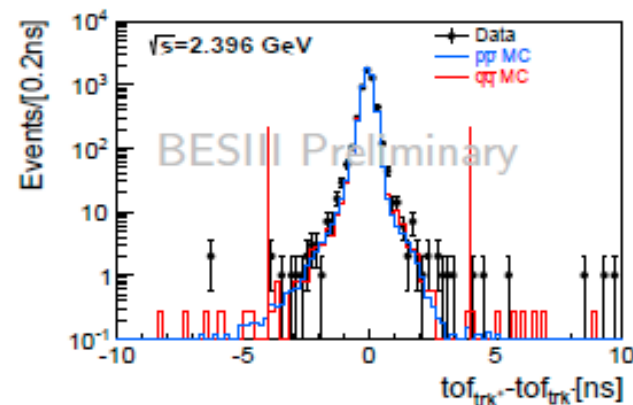
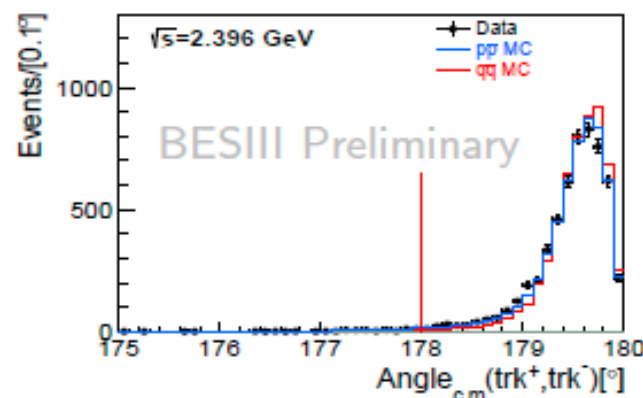
■ Vertex fit.



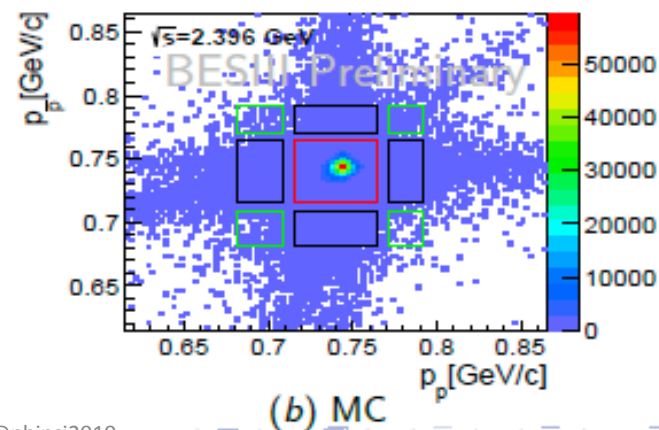
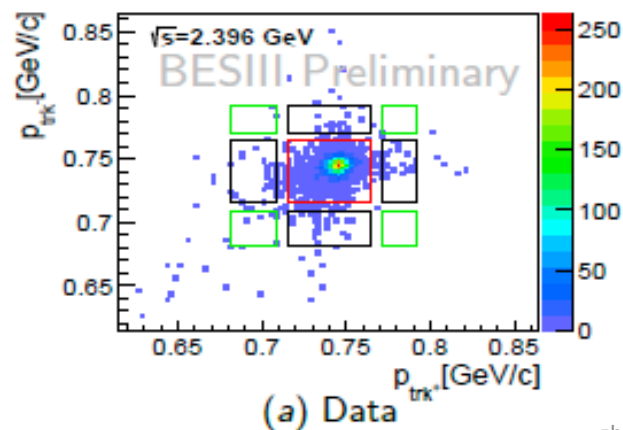
*The red line is $p\bar{p}$ signal events selected from the $q\bar{q}$ samples.

Scan data (selection II)

- Require angle between p and \bar{p} in center-of-mass criteria.
- Require TOF information.

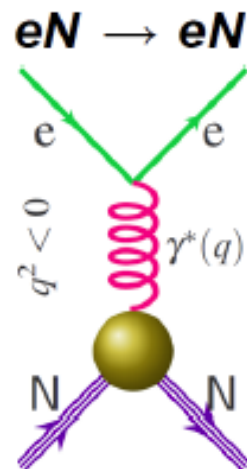


- Momentum window for p and \bar{p} :

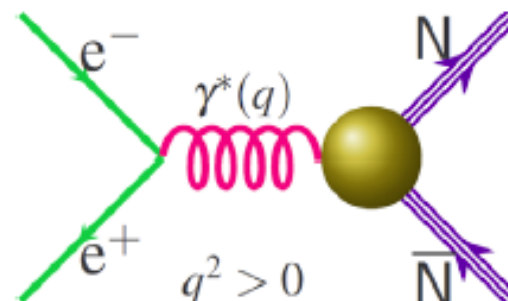


Electromagnetic Form Factors

Space-like:
FF real



$$e^+e^- \leftrightarrow N\bar{N}, \Lambda\bar{\Lambda}, \dots$$



Time-like:
FF complex

Vector current, **two form factors** (F_1 and F_2)

$$\Gamma_\mu = e\bar{u}(p')[F_1(q^2)\gamma_\mu + \frac{\kappa}{2M_N}F_2(q^2)i\sigma_{\mu\nu}q^\nu]u(p)e^{iqx}$$

Dirac

$$F_1^p(q^2 = 0) = 1$$

$$F_1^n(q^2 = 0) = 0$$

Pauli

$$F_2^p(q^2) = 1$$

$$F_2^n(q^2) = 1$$

Sachs

$$G_E = F_1 + \frac{\kappa q^2}{4M^2}F_2$$

$$G_M = F_1 + \kappa F_2$$

$$G_E(4M_p^2) = G_M(4M_p^2)$$

G.S. Huang: Baryon FF @BESIII



Baryon-pair production near threshold

- The Born cross section for $e^+e^- \rightarrow \gamma^* \rightarrow B\bar{B}$, can be expressed in terms of electromagnetic form factor G_E and G_M :

$$\sigma_{B\bar{B}}(m) = \frac{4\pi\alpha^2 C \beta}{3m^2} [|G_M(m)|^2 + \frac{1}{2\tau} |G_E(m)|^2]$$

$\alpha = \frac{1}{137}$ is fine structure constant, $\beta = \sqrt{1 - 4m_B^2/m^2}$ is the velocity,
 $\tau = m^2/4m_B^2$

- The Coulomb factor $C = \begin{cases} \frac{\pi\alpha}{\beta} \frac{1}{1 - \exp(-\frac{\pi\alpha}{\beta})} & \text{for a charged } B\bar{B} \text{ pair} \\ 1 & \text{for a neutral } B\bar{B} \text{ pair} \end{cases}$

- For the neutral pair production, the cross section **should be 0 at threshold**, and is expected to increase with the velocity near the threshold.