# **Charm Physics**



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# Overview

Understanding QCD in charmed hadrons;

CP violation in charm;

SM in rare charm decays;

Dark Matter search in rare charm decays;

From B anomalies to NP in charm;

Signatures of NP in chargned current and FCNC charm decays;

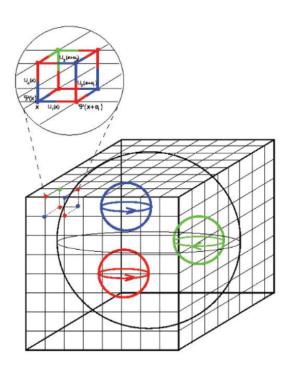
Summary and Outlook

# Theory goals

# Deepening our knowledge of SM → QCD

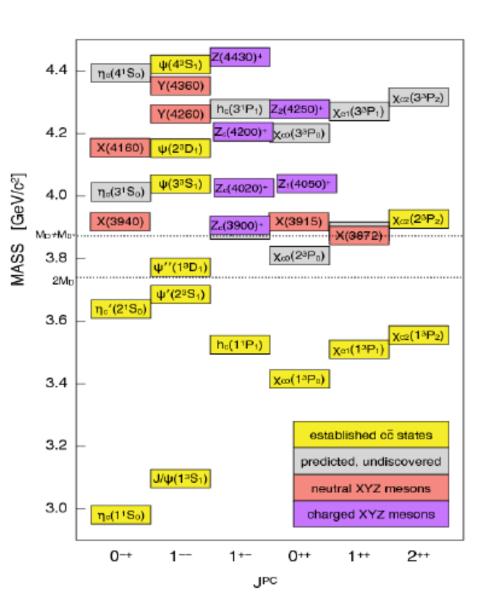
Charm spectroscopy- tetraquark states decay constants, form-factors, mixing parameters...

QCD (lattice) in action!



#### QCD in action:

# Charmonium and Exotic Spectroscopy with Charm Quarks in Lattice QCD



- Plethora of unexpected charmonium-like
   (X, Y, Z) states discovered experimentally
- Masses and widths of some D<sub>s</sub> states significantly lower than those expected from quark model.
- Tetraquarks? Molecules? Cusps? Hybrids?
- First principles calculations using lattice QCD to understand these states.

# Search for New Physics

Solution by New Physics

B meson puzzles

Tests of Lepton flavour universality

 $(g-2)_{\mu}$  discrepancy SM prediction and experimental result

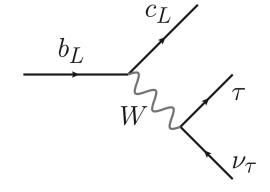
How about charm?

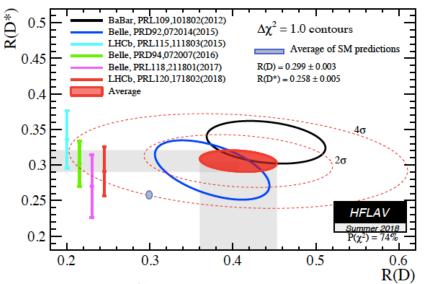
- CP violation in the up sector;
- Charm offers tests of possible NP in up sector at low-energies;
- If NP couples to weak doublets of quarks, CKM connects it with charm sector.
- Can one see NP in charm decays not being present in B meson?

#### B physics anomalies: experimental results ≠ SM predictions!

#### charged current (SM tree level)

$$R_{D^{(*)}} = rac{BR(B o D^{(*)} au 
u_{ au})}{BR(B o D^{(*)} \mu 
u_{\mu})}$$
 3.8 $\sigma$ 





Freytsis, et al., 1506.08896, S.F. et al., 1206.1872;

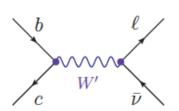
Di Luzio & Nardecchia, 1706.01868, Bernlochner et al., 1703.05330, F. Feruglio et al., 1806.10155, 1606.00524.

$$\mathcal{L}_{eff} = -\frac{4G_F}{\sqrt{2}} V_{cb} [(1 + g_{V_L})(\bar{c}_L \gamma_\mu b_L)(\bar{l}_L \gamma^\mu \nu_L) + g_{V_R}(\bar{c}_R \gamma_\mu b_R)(\bar{l}_L \gamma^\mu \nu_L) + g_{S_R}(\bar{c}_L b_R)(\bar{l}_R \nu_L) + g_{T_R}(\bar{c}_L \sigma_{\mu\nu} b_R)(\bar{l}_R \sigma^{\mu\nu} \nu_L)]$$

Assuming NP at scale  $\Lambda_{NP}$  (Di Luzio, Nardecchia, 1706.0!868)

$$\frac{4G_F}{\sqrt{2}}V_{cb}\,g_V \to \frac{2}{\Lambda_{NP}^2}$$

What is the scale of New Physics?



$$\Lambda_{NP} \simeq 3 \, {\rm TeV}$$

b  $\ell$   $\Delta$ 

Perturbativity of NP

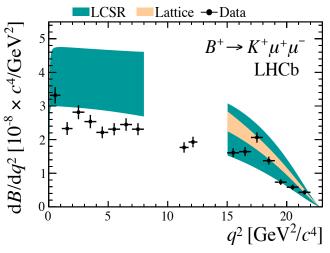
$$\mathcal{L}_{NP} \supset \frac{C_D}{\Lambda_{NP}^2} (\bar{c}_L \Gamma_\mu b_L) (\tau_L \gamma^\mu \nu_L)$$

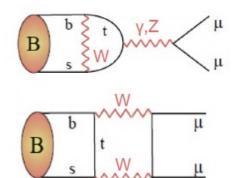
V-A form of NP

(current)(current) operators are invariant under QCD running

$$\Lambda_{NP} > 3 \, {
m TeV}$$
 C<sub>D</sub> becomes non-perturbative!

# FCNC - SM loop process





$$R_K = \frac{\mathcal{B}(B \to K\mu\mu)_{q^2 \in [1,6]\text{GeV}^2}}{\mathcal{B}(B \to Kee)_{q^2 \in [1,6]\text{GeV}^2}} = 0.745 \pm_{0.074}^{0.090} \pm 0.036$$

 $2.4\sigma$ 

$$R_{K^*}^{\text{central}} = \frac{\mathcal{B}(B \to K\mu\mu)_{q^2 \in [1.1,6] \text{GeV}^2}}{\mathcal{B}(B \to Kee)_{q^2 \in [1.1,6] \text{GeV}^2}} = 0.685 \pm_{0.069}^{0.113} \pm 0.047,$$

What is the scale of New Physics?

$$\mathcal{L}_{NP} = \frac{1}{\Lambda_{NP}^2} \bar{s}_L \gamma^\alpha b_L \bar{\mu}_L \gamma_\alpha \mu_L$$

 $\Lambda_{NP} \simeq 30 \, {\rm TeV}$ 

# NP explaining both B anomalies

$$\mathcal{R}_{D^{(*)}}^{exp} > \mathcal{R}_{D^{(*)}}^{SM}$$

$$\mathcal{L}_{NP} = \frac{1}{(\Lambda^D)^2} 2 \, \bar{c}_L \gamma_\mu b_L \bar{\tau} \gamma^\mu \nu_L$$

$$\mathcal{L}_{NP} = \frac{1}{(\Lambda^K)^2} \bar{s}_L \gamma_\mu b_L \bar{\mu}_L \gamma^\mu \mu_L$$

$$\Lambda^D \simeq 3 \, \text{TeV}$$

$$\Lambda^K \simeq 30 \, \text{TeV}$$

$$\Lambda^D \simeq \Lambda^K \equiv \Lambda$$

NP in FCNC 
$$B \to K^{(*)} \mu^+ \mu^-$$
 has to be suppressed

$$\frac{1}{(\Lambda^K)^2} = \frac{C_K}{\Lambda^2} \qquad C_K \simeq 0.01$$

#### Charged current charm meson decays and New Physics

$$\mathcal{L}_{SM} = \frac{4G_F}{\sqrt{2}} V_{cs} \bar{s}_L \gamma^\mu c_L \, \bar{\nu}_l \gamma_\mu l \qquad \mathcal{L}_{NP} = \frac{2}{\Lambda_c^2} \bar{s}_L \gamma^\mu c_L \, \bar{\nu}_l \gamma_\mu l$$

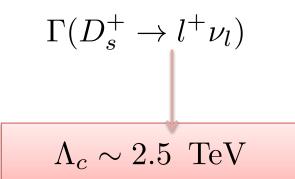
$$\mathcal{L}_{NP} = \frac{2}{\Lambda_c^2} \bar{s}_L \gamma^\mu c_L \, \bar{\nu}_l \gamma_\mu l$$

PDG 2018

$$f_{D^+} = 211.9(1.1) \text{ MeV}$$

$$f_{D_s} = 249.0(1.2) \text{ MeV}$$

$$\frac{f_{Ds}}{f_{D^+}} = 1.173(3) \; . \label{eq:fDs}$$



$$|V_{cs}| = 0.997 \pm 0.017$$

Electro-magnetic correction 1-3%

#### Message:

Even if there is NP at 3 TeV scale the effect on charm leptonic decay can be ~ 1%!

# New Physics in charm processes



Constraints from K, B physics

Constraints from EW physics, oblique corrections,  $Z \to b \bar b$ 

Constraints from LHC

NP in charm

Up quark in weak doublet "talks" to down quark via CKM!

Effects of NP in charm suppressed by V<sub>cb</sub>\* V<sub>ub</sub>.

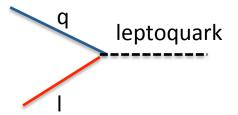
$$Q_{iL} = \begin{bmatrix} V^*_{il} u_J \\ d_i \end{bmatrix}$$

# Models of NP explaining B anomalies

Spin	Color singlet	Color tripet
0	2HDM	Scalar LQ  Parity - sbottom
1	W' ,Z'	Vector LQ
Leptoquarks?		Dark ma <b>sten</b> nmary

2HDMII cannot explain  $R_{D(*)}$ 

New gauge bosons, W', Z'-difficult to construct UV complete theory



Nature of anomaly requires NP in quark and lepton sector! It seems that LQs are ideal candidates to explain all B anomalies at tree level!

 $(SU(3)_c, SU(2)_L, U(1)_Y)$ 

- Is charm physics sensitive on NP explaining B puzzles?
- Can some NP be present in charm and not in beauty mesons?

## LQ and charm charged current

Triplet LQ S<sub>3</sub> in charm leptonic decays decay

$$\mathcal{L}_{\bar{u}^i d^j \bar{\ell} \nu_k} = -\frac{4G_F}{\sqrt{2}} \left[ (V_{ij} U_{\ell k} + g^L_{ij;\ell k}) (\bar{u}^i_L \gamma^\mu d^j_L) (\bar{\ell}_L \gamma_\mu \nu^k_L) \right]$$

Test of lepton flavour universality (LFU)

$$R_{\tau,\mu}^c = \frac{\Gamma(D_s \to \tau \nu)}{\Gamma(D_s \to \mu \nu)}$$

$$\frac{R_{\tau,\mu,LQ}^c}{R_{\tau,\mu,SM}^c} = \left[1 - \frac{v^2}{2M_{S3}^2} \overline{\text{Re}((Vy*)_{c\tau}y_{s\tau} - (Vy*)_{c\mu}y_{s\mu})}\right]$$

$$S_3 = (3,3,-1/3)$$

Comes from the fit of  $R_{K(*)}$  with  $S_3$ 

Doršner, SF, Greljo, Kamenik Košnik, 1603.04993;

m <sub>s3</sub> [TeV]	$1 - R_{\tau,\mu,LQ}^c / R_{\tau,\mu,SM}^c$
1.0 1.2	3.2% 2.4%
1.5	1.5%

#### CHARM quark electric (chromo-electric) dipole moment

$$\mathcal{L}_{\text{eff}} = \frac{d_q}{2} \frac{1}{2} (\bar{q} \sigma_{\mu\nu} i \gamma_5 q) F^{\mu\nu} + \tilde{d}_q \frac{1}{2} (\bar{q} \sigma_{\mu\nu} T^a i \gamma_5 q) g_s G^{\mu\nu}_a + w \frac{1}{6} f^{abc} \epsilon^{\mu\nu\lambda\rho} G^a_{\mu\sigma} G^b_{\nu} G^c_{\lambda\rho}$$
 quark CEDM Weingerg operator 
$$w = \frac{g_s^3}{32\pi^2} \frac{\tilde{d}_q}{m_q}.$$

Sala, 1312.2589 Considered charm quark EDM and CEDM

CEDM threshold correction to w

$$| ilde{d}_c|\lesssim 1.0 imes 10^{-22}{
m cm}$$
 from neutron EDM  $|d_c|\lesssim 4.4 imes 10^{-17}e\,{
m cm}$  from  $B o X_s\gamma$ 

In 1809.09114, Dekens et al, NP from B anomalies creates c-quark EDM, which can be related to neutron (lattice computation of c –bar c content of neutron) or Hg EDM!

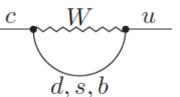
More studies of charm quark EDM(CEDM) – new source of CP violation!

# SM effective Hamiltonian for rare charm decays -FCNC

$$\mathcal{H}_{\text{eff}} = \lambda_d \mathcal{H}^d + \lambda_s \mathcal{H}^s - \frac{4G_F \lambda_b}{\sqrt{2}} \sum_{i=3,\dots,10,S,P,\dots} C_i \mathcal{O}_i$$

$$\lambda_q = V_{uq} V_{cq}^*$$
 Tree-level 4-quark operators

(Short-distance) penguin operators



- 1) At scale m<sub>w</sub> all penguin contributions vanish due to GIM;
- 2) SM contributions to  $C_{7...10}$  at scale  $m_c$  entirely due to mixing of tree-level operators into penguin ones under QCD
- 3) SM values at m<sub>c</sub>

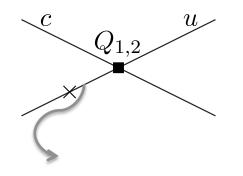
$$C_7 = 0.12, \qquad C_9 = -0.41$$

(recent results: de Boer, Hiller, 1510.00311, 1701.06392, De Boer et al, 1606.05521) 1707.00988)

C. Greub et al., PLB 382 (1996) 415;

$$BR(D \to X_u \gamma) \sim 10^{-8}$$

branching ratio	$D^0 \to \rho^0 \gamma$	$D^0  o \omega \gamma$	$D^0  o \phi \gamma$	$D^0 \to \bar K^{*0} \gamma$
Belle $[24]^{\dagger}$	$(1.77 \pm 0.31) \times 10^{-5}$	_	$(2.76 \pm 0.21) \times 10^{-5}$	$(4.66 \pm 0.30) \times 10^{-4}$
BaBar $[33]^{\dagger a}$	_	_	$(2.81 \pm 0.41) \times 10^{-5}$	$(3.31 \pm 0.34) \times 10^{-4}$
CLEO [34]	_	$< 2.4 \times 10^{-4}$	_	_



Hiller & De Boer 1701.06392

photon emission

Note: all SM th. predictions for BR(D<sup>0</sup>  $\rightarrow$   $\rho^0 \gamma$ ) smaller than exp. rate!

#### previuos works:

SF& Singer, hep-ph/9705327, SF, Prelovsek &hep-ph/9801279
S. F. P. Singer and J. Zupan, EPJC 27(2003) 201 Burdman et al. hep-ph/9502329, Khodjamirian et al, hep-ph/9506242

# CP asymmetry in charm radiative decays

$$A_{CP}(D \to V\gamma) = \frac{\Gamma(D \to V\gamma) - \Gamma(\bar{D} \to \bar{V}\gamma)}{\Gamma(D \to V\gamma) + \Gamma(\bar{D} \to \bar{V}\gamma)}$$

$$|A_{CP}^{\rm SM}| < 2 \cdot 10^{-3}$$

Belle, 1603.03257

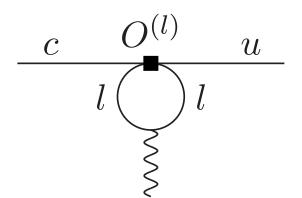
$$A_{CP}(D^0 \to \rho^0 \gamma) = 0.056 \pm 0.152 \pm 0.006$$
,  
 $A_{CP}(D^0 \to \phi \gamma) = -0.094 \pm 0.066 \pm 0.001$   
 $A_{CP}(D^0 \to \bar{K}^{*0} \gamma) = -0.003 \pm 0.020 \pm 0.000$ 

Hiller& de Boer 1701. 06392 LQs give as large contributions as SM

# New Physics in FCNC charm decays

Leptoquarks in  $c \rightarrow u\gamma$ 

Hiller& de Boer 1701. 06392 SF and Košnik, 1510.00965



Constraints from

$$\tau^{-} \to \pi^{-} \nu_{\tau}$$

$$\tau^{-} \to K^{-} \nu_{\tau}$$

$$\Delta m_{D}$$

$$D^{+} \to \tau^{+} \nu_{\tau}$$

$$D_{s}^{+} \to \tau^{+} \nu_{\tau}$$

$$K^{+} \to \pi^{+} \nu \bar{\nu}$$

Even for  $\tau$  in the loop too small contribution!

Masses of  $m_{LQ} \approx 1 \text{ TeV}$ .

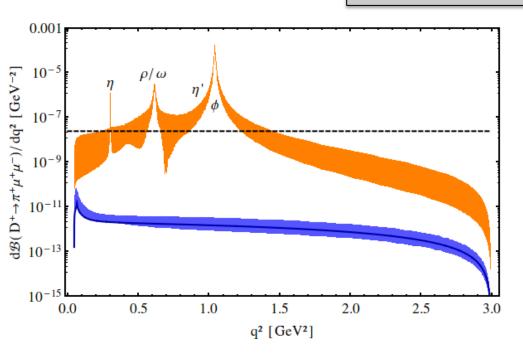
$$S_3 = (3,3,-1/3)$$

Within LQ models the c  $\rightarrow$  u $\gamma$  branching ratios are SM-like with CP asymmetries at O(0.01) for S<sub>1,2</sub> and V $_2$  and SM-like for S<sub>3</sub>.

Vector LQ  $V_{1}A_{CP} \sim O(10\%)$ . The largest effects arise from τ-loops.

 $S_3$  can explain  $R_{K(*)}$ !

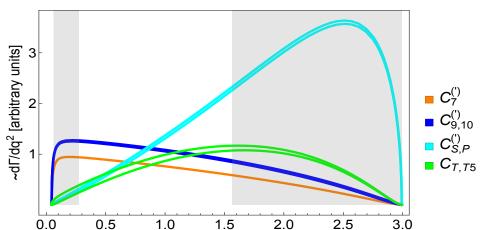
# $D \to \pi l^+ l^-$



SM prediction: Long distance contributions most important!

$$D \to \pi V \to \pi l^+ l^-$$
 peaks at  $\rho, \omega, \phi$  and  $\eta$  resonances

de Boer, Hiller, 1510.00311, SF and Kosnik, 1510.00965



 $q^2[GeV^2]$ 

Maximally allowed values of the Wilson coefficients in the low and high energy bins, according to LHCb 1304.6365:

LHCb 1304.6365

	$\mathrm{BR}(\pi\mu\mu)_{\mathrm{I}}$	$\mathrm{BR}(\pi\mu\mu)_{\mathrm{II}}$	DD(D)
П		$DR(\pi\mu\mu)\Pi$	$\left  \operatorname{BR}(D^0 \to \mu \mu) \right $
$\parallel$ $ ilde{C}_7$ $\parallel$	2.4	1.6	-
$\left\   ilde{C}_{9}  imes  ight\ $	2.1	1.3	
$\left\   ilde{C}_{10}  ilde{ }  ight\ $	1.4	0.92	0.56
$ ilde{C}_S$	4.5	0.38	0.043
$\tilde{C}_P$	3.6	0.37	0.043
$\left\   ilde{C}_T  ilde{ }  ight\ $	4.1	0.76	-
$\left\   ilde{C}_{T5}  imes  ight\ $	4.4	0.74	
$ \begin{array}{c c} \tilde{C}_{7} \\ \tilde{C}_{9} \\ \tilde{C}_{10} \\ \tilde{C}_{S} \\ \tilde{C}_{P} \\ \tilde{C}_{T} \\ \tilde{C}_{T5} \\ \tilde{C}_{9} = \pm \tilde{C}_{10} \end{array} $	1.3	0.81	0.56

from

Best bounds

$$D^0 \to \mu^+ \mu^-$$

$$|\tilde{C}_i| = |V_{ub}V_{cb}^*C_i|$$

0.043

region I

region II

$$q^2 \in [0.0625, 0.276] \, GeV^2$$

$$q^2 \in [1.56, 4.00] \, GeV^2$$

$$BR(D^0 \to \mu^+ \mu^-) < 6.2 \times 10^{-9}$$

Model	Effect	Size of the effect
Scalar leptoquark (3,2,7/6)	C <sub>S</sub> ,C <sub>P</sub> , C <sub>S</sub> ',C <sub>P</sub> ',C <sub>T</sub> ,C <sub>T5</sub> , C <sub>9</sub> ,C <sub>10</sub> ,C <sub>9</sub> ',C <sub>10</sub> '	V <sub>cb</sub> V <sub>ub</sub>  C <sub>9,</sub> C <sub>10</sub>  < 0.34
Vector leptoquark (3,1,5/3)	C <sub>9</sub> ' = C <sub>10</sub> '	V <sub>cb</sub> V <sub>ub</sub>  C <sub>9</sub> ', C <sub>10</sub> ' < 0.24
Two Higgs doublet  Model type III  C <sub>S</sub> ,C <sub>P</sub> , C <sub>S</sub> ',C <sub>P</sub> '		$V_{cb}V_{ub} C_S - C_S'  < 0.005$ $V_{cb}V_{ub} C_P - C_P'  < 0.005$
Z' model	C <sub>9</sub> ',C <sub>10</sub> '	$V_{cb}V_{ub} C_{9',} <0.001$ $V_{cb}V_{ub} C_{10'} <0.014$

# Lepton flavor violation

$$c \to u \mu^{\pm} e^{\mp}$$

1510.00311 (de Beor and Hiller) 1705.02251 (Sahoo and Mohanta)

$$\mathcal{L}_{\text{eff}}^{\text{weak}}(\mu \sim m_c) = \frac{4G_F}{\sqrt{2}} \frac{\alpha_e}{4\pi} \sum_{i} \left( K_i^{(e)} O_i^{(e)} + K_i^{(\mu)} O_i^{(\mu)} \right)$$

$$O_9^{(e)} = (\bar{u}\gamma_\mu P_L c) (\bar{e}\gamma^\mu \mu)$$

$$O_9^{(\mu)} = (\bar{u}\gamma_\mu P_L c) (\bar{\mu}\gamma^\mu e)$$

LHCb bound, 1512.00322

$$BR(D^0 \to e^+\mu^- + e^-\mu^+) < 2.6 \times 10^{-7}$$

$$BR(D^+ \to \pi^+ e^+ \mu^-) < 2.9 \times 10^{-6}$$

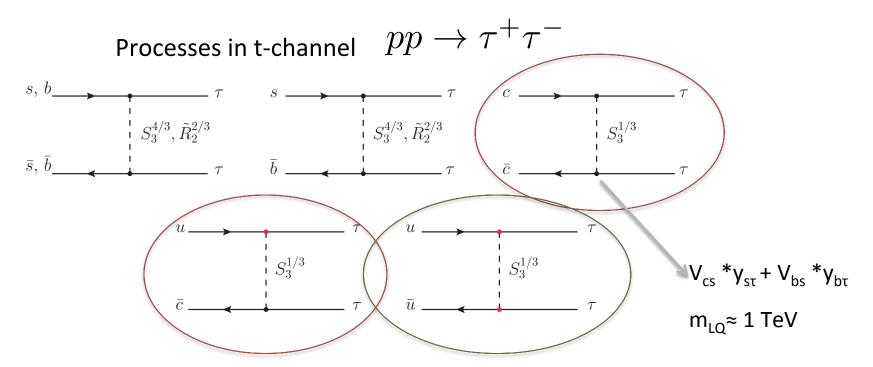
$$BR(D^+ \to \pi^+ e^- \mu^+) < 3.6 \times 10^{-6}$$

$$\begin{aligned}
\left| K_{S,P}^{(l)} - K_{S,P}^{(l)\prime} \right| &\lesssim 0.4, \\
\left| K_{9,10}^{(l)} - K_{9,10}^{(l)\prime} \right| &\lesssim 6, \quad \left| K_{T,T5}^{(l)} \right| &\lesssim 7, \end{aligned}$$

$$BR(D^0 \to e^{\pm} \tau^{\mp}) < 7 \times 10^{-15}$$

$$l = e, \mu$$

# LHC constraints on $S_3$ : high-mass $\tau \tau$ production



Flavour anomalies generate s  $\tau$  , b $\tau$  and c $\tau$  relatively large couplings.

s quark pdf function for protons are ~ 3 times lagrer contribution then for b quark.

1706.07779, Doršner, SF, Faroughy, Košnik

$$\sigma_{s\bar{s}}(y_{s\tau}) = 12.042 y_{st}^4 + 5.126 y_{st}^2,$$

$$\sigma_{s\bar{b}}(y_{s\tau}, y_{b\tau}) = 12.568 y_{s\tau}^2 y_{b\tau}^2,$$

$$\sigma_{b\bar{b}}(y_{b\tau}) = 3.199 y_{b\tau}^4 + 1.385 y_{b\tau}^2,$$

$$\sigma_{c\bar{c},u\bar{u}.u\bar{c}}(y_{s\tau}) = 3.987 y_{s\tau}^4 - 5.189 y_{s\tau}^2.$$

#### Dark Matter in charm decays

Belle collaboration 1611.09455 BR(D<sup>0</sup>  $\rightarrow$  invisible) <9.4 × 10<sup>-5</sup>

SM: BR(D<sup>0</sup> 
$$\rightarrow$$
 vv) = 1.1 × 10<sup>-30</sup>

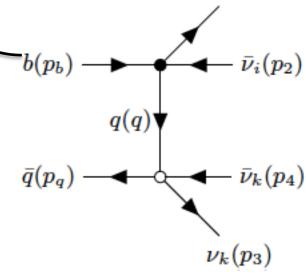
Badin & Petrov 1005.1277 suggested to search for processes with missing energy/£ in

Bhattacharya, Grant and Petrov 1809.04606

$$\mathcal{B}(D \to invisibles) = \mathcal{B}(D \to \nu\bar{\nu}) + \mathcal{B}(D \to \nu\bar{\nu} + \nu\bar{\nu}) + \dots$$
 c instead of b

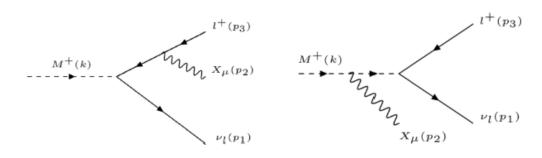
The SM contributions to invisible widths of heavy mesons  $\Gamma(D^0 \to \text{missing energy})$  are completely dominated by the four-neutrino transitions  $D^0 \to v\bar{v}v\bar{v}$ .

$$\mathcal{B}(D^0 \to \nu \bar{\nu} \nu \bar{\nu}) = (2.96 \pm 0.39) \times 10^{-27}$$

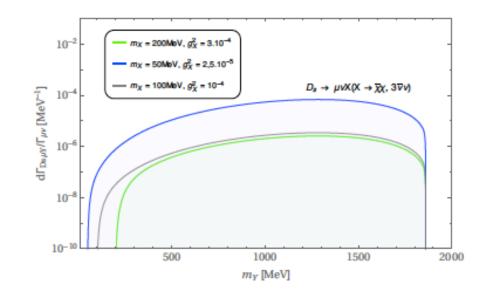


$$A_{\mu}$$
 and  $X_{\mu}$  mix via  $\kappa$ 

$$M^+ \to \mu^+ \not\!\!\! E$$



Radiative - not  $\gamma$  but X



Is it possible to search for decay  $D \rightarrow \mu X$  X is SM  $v_{\mu}$  + DM gauge boson  $\rightarrow$  invisible fermions

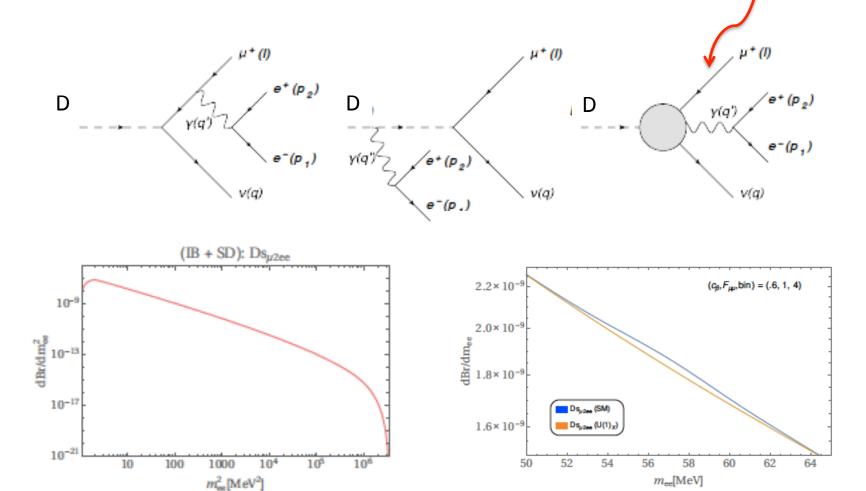
Exp:  $D \rightarrow \tau \bar{\nu}_{\tau} \rightarrow \mu \bar{\nu}_{\mu} \nu_{\tau} \bar{\nu}_{\tau}$ 

Difficult to differentiate

- There is a possibility that X → e<sup>+</sup>e<sup>-</sup>
- Can one see it in the decays P→μ vX → μυ e<sup>+</sup>e

First one should calculate SM values

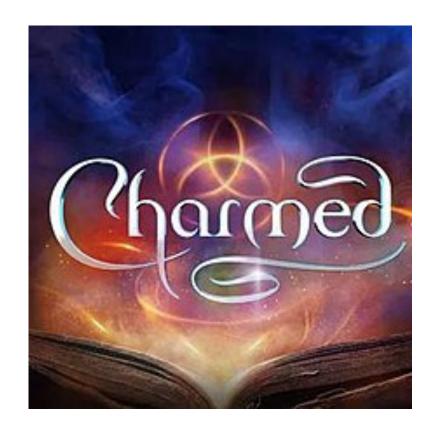
Thanks D. Melikhov for providing us with  $\langle \gamma^* | J_{\mu} | D_s \rangle$ 



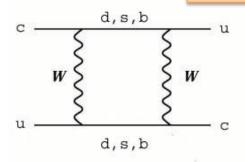
## Summary and outlook

- QCD (lattice) a lot of open issues in Charm spectroscopy!
   Improvement on decay constants and form-factors!
- CP-violation in up sector (NP search) more studies on direct CP violation and (C)EDM of c-quark;
- New physics explaining B anomalies, leads to rather small effects in charge current transitions;
- FCNC transition small contribution of Leptoquarks in charm decays observables;
- To perform all possible test of LFU;
- Few proposals to test DM in charm physics;
- Charm physics complement any search for NP at low energies!

# Thanks!

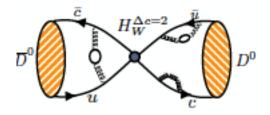


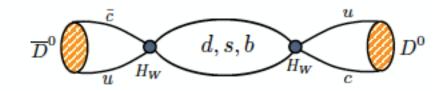
## Mixing and indirect CP violation



- intermediate down-type quarks;
- due to CKM contribution of b quark negligible;
- in the SU(3) limit 0;

$$M_{12} - \frac{i}{2}\Gamma_{12} \propto \langle D^0|H_W^{\Delta c=2}|\overline{D}^0\rangle + \sum_n \frac{\langle D^0|H_W^{\Delta c=1}|n\rangle\langle n|H_W^{\Delta c=1}|\overline{D}^0\rangle}{M_D - E_n + i\epsilon}$$





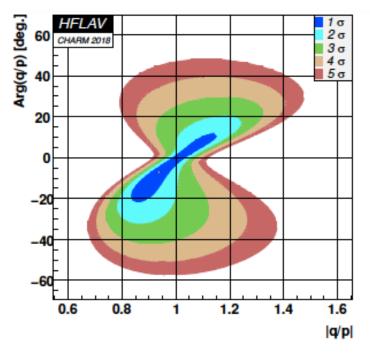
Short distance Lattice QCD helps!

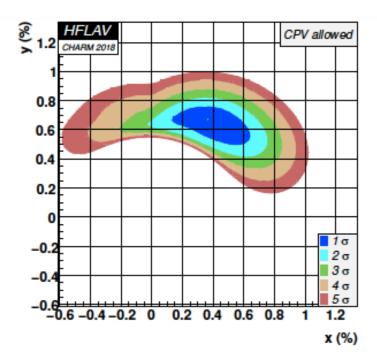
Long distance difficult to determine

Lattice determined

$$\langle \mathcal{O}_i \rangle \equiv \langle D^0 | \mathcal{O}_i | \overline{D}^0 \rangle (\mu) = e_i M_D^2 f_D^2 B_D^{(i)}(\mu)$$

Possible NP effect difficult to isolate!



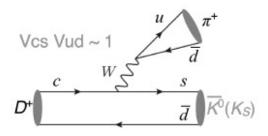


$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\overline{D}^0\rangle$$

- |q/p≠|1 would indicate CPV in mixing.
- $Arg(q/p) \neq 0$  would indicate CPV from interference mixing/decay.
- Mixing parameters  $x = \Delta m/\Gamma$  and  $y = \Delta \Gamma/(2\Gamma)$ .

#### SM features of CPV in D

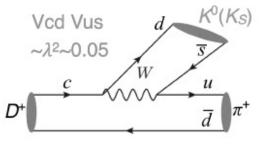
- CPV in D  $\bar{\rm D}$  mixing suppressed due to  ${\cal O}(V_{cb}V_{ub}^*/V_{cs}V_{us}^*)\sim 10^{-3}$
- direct CPV suppressed due to  $\mathcal{O}([V_{cb}V_{ub}^*/V_{cs}V_{us}^*]\alpha_s/\pi) \sim 10^{-4}$



$$A_{CP}^{D^+\to K_S^0\pi^+}=(-0.363\pm0.094\pm0.067)\%$$

Belle, 1203.6409, mainly attributed to the K mixing

1707.09297, Wang, F.S. Yu, and H.N.Li,



Nierste & Schacht, 1708.03572

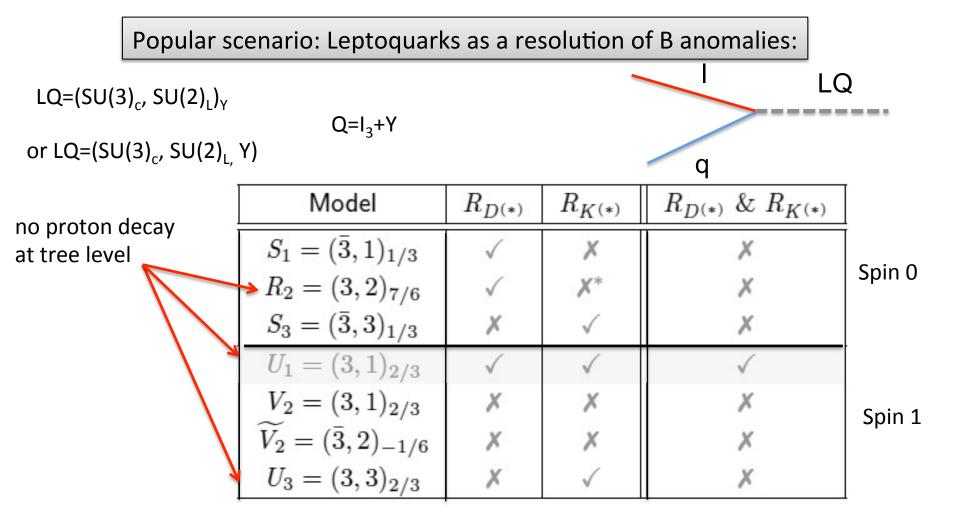
the time-dependent and time-integrated CP asymmetries in 
$$D\to fK_S(\to\pi^+\pi^-)$$
 the interference CF and the DCS amplitudes with the K

mixing, effect of the order 10<sup>-3</sup>.

Proposal: search for the difference of the time-integrated CP asymmetries in the mode with  $\pi$  and K.

NP might be present!

$$A_{CP}^{dir}(D \to K_S K^{0*}) \le 0.3\%$$



No single scalar LQ to solve simultaneously both anomalies!

Scalar LQ simpler UV completion;

Doršner, SF, Greljo, Kamenik, Košnik, 1603.04993

Only  $R_2$  and  $S_1$  might explain  $(g-2)_{\mu}$  (both chiralities are required with the enhancement factor  $m_t/m_{\mu}$ ) Muller 1801.0338.

# Angular distributions in $D \rightarrow P_1 P_2 I^+I^-$

LHCb, 1707.08377

$$\begin{split} \mathcal{B}(D^0 \to \pi^+ \pi^- \mu^+ \mu^-)|_{[0.565-0.950]\,\mathrm{GeV}} &= (40.6 \pm 5.7) \times 10^{-8} \,, \\ \mathcal{B}(D^0 \to \pi^+ \pi^- \mu^+ \mu^-)|_{[0.950-1.100]\,\mathrm{GeV}} &= (45.4 \pm 5.9) \times 10^{-8} \,, \\ \mathcal{B}(D^0 \to K^+ K^- \mu^+ \mu^-)|_{[>0.565]\,\mathrm{GeV}} &= (12.0 \pm 2.7) \times 10^{-8} \,, \end{split}$$

De Beor and Hiller, 1805.08516 Modes sensitive to NP

$$D^0 \to \pi^+ \pi^- l^+ l^- \,, \quad D^0 \to K^+ K^- l^+ l^- \,,$$
 
$$D^+ \to K^+ \bar{K}^0 l^+ l^- \,,$$
 
$$D_s \to K^+ \pi^0 l^+ l^- \,, \quad D_s \to K^0 \pi^+ l^+ l^- \,,$$

- study of angular distributions SM null tests
- simpler then in B decays due to dominance of long distance physics  $R_{\pi\pi}^{D\,{
  m SM}}=1.00\pm\mathcal{O}(\%)$  (resonances)
- NP induced integrated CP asymmetries can reach few percent
- sensitive on C<sub>10</sub><sup>(')</sup>

$$R_{KK}^{D\,\text{SM}} = 1.00 \pm \mathcal{O}(\%)$$

$$A_{2\phi}(D^0 \to \pi^+\pi^-\mu^+\mu^-) = (-0.6 \pm 3.7 \pm 0.6)\%,$$
  
 $A_{CP}(D^0 \to \pi^+\pi^-\mu^+\mu^-) = (-0.6 \pm 3.8 \pm 0.7)\%,$ 

 $A_{\rm FB}(D^0 \to \pi^+\pi^-\mu^+\mu^-) = (3.3 \pm 3.7 \pm 0.6)\%,$ 

$$A_{\rm FB}(D^0 \to K^+ K^- \mu^+ \mu^-) = (0 \pm 11 \pm 2)\%,$$

$$A_{2\phi}(D^0 \to K^+K^-\mu^+\mu^-) = (9 \pm 11 \pm 1)\%,$$

$$A_{CP}(D^0 \to K^+K^-\mu^+\mu^-) = (0 \pm 11 \pm 2)\%,$$

Tests of LFU

$$R_{P_1P_2}^D = \frac{\int_{q_{\rm min}^2}^{q_{\rm max}^2} d\mathcal{B}/dq^2(D \to P_1P_2\mu^+\mu^-)}{\int_{q_{\rm min}^2}^{q_{\rm max}^2} d\mathcal{B}/dq^2(D \to P_1P_2e^+e^-)}$$

LHCb , 1806.10793 consistent with SM