

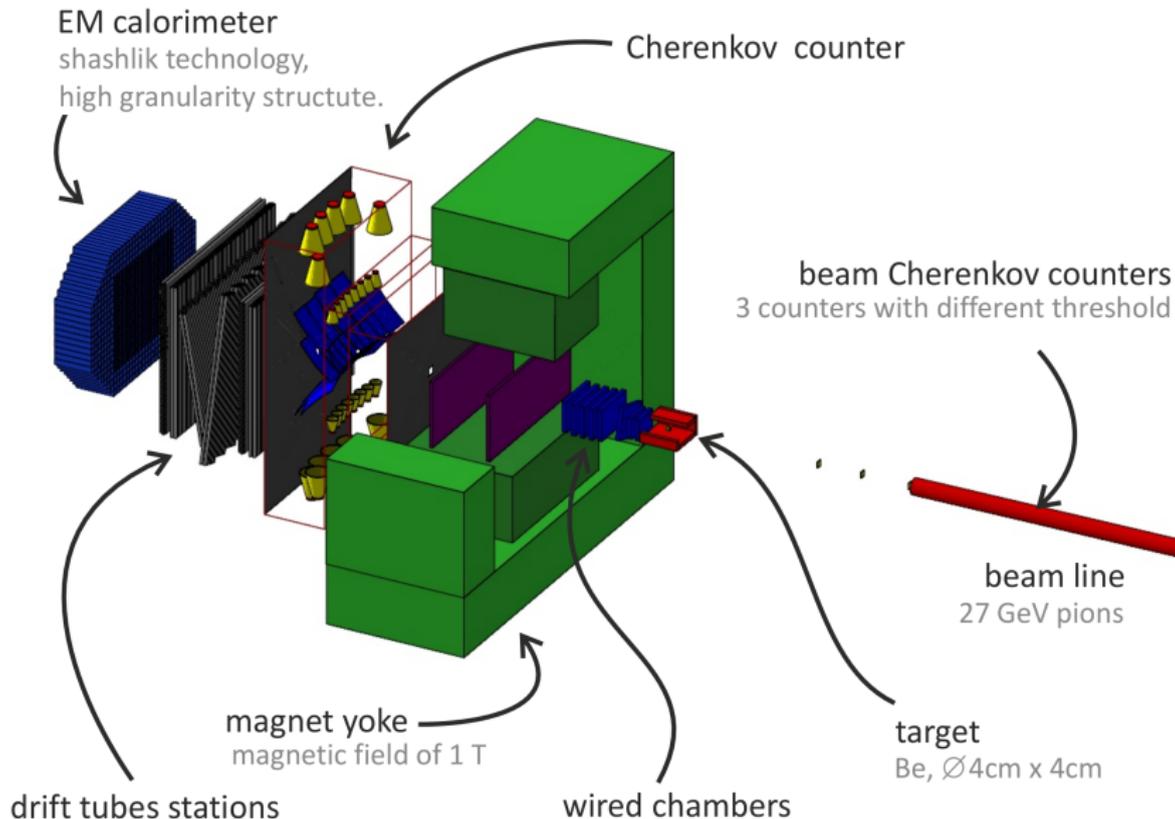
Meson spectroscopy at VES and COMPASS

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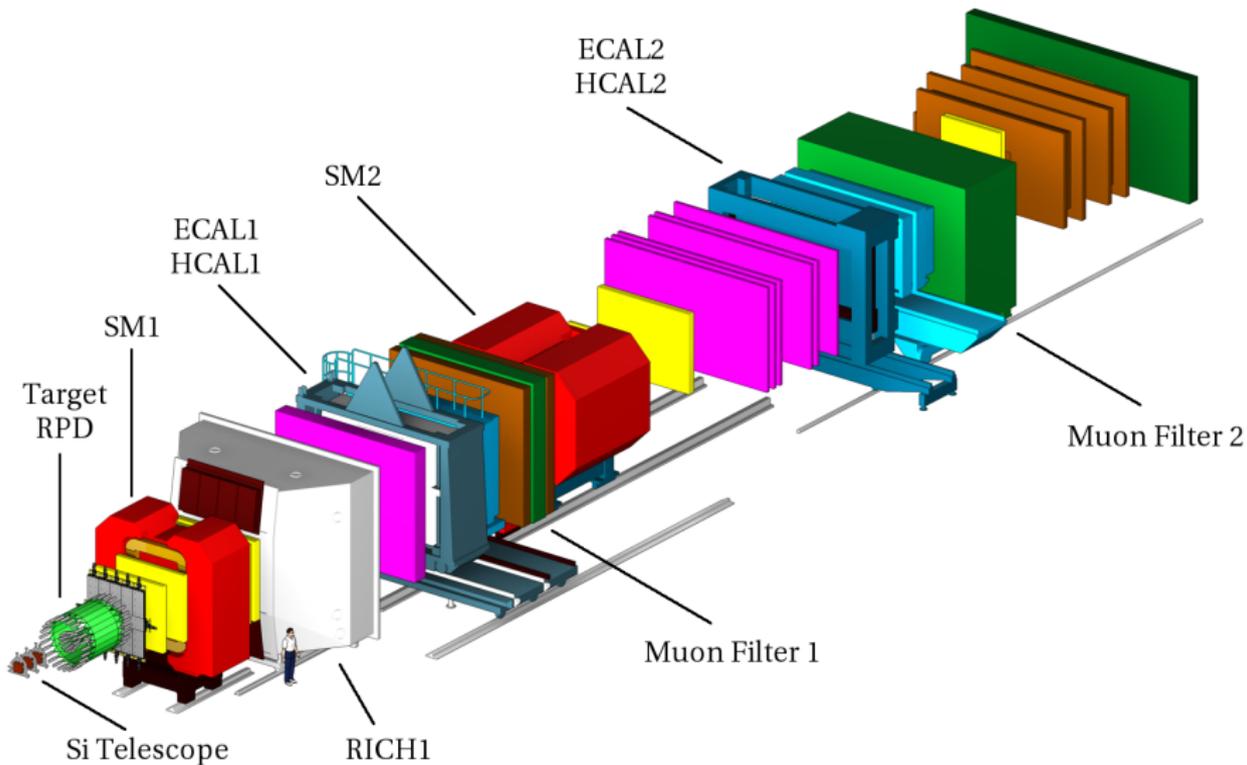
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- Two flagship reactions for VES and COMPASS experiments:
 $\pi^- N \rightarrow \pi^- \pi^- \pi^+ N$ and $\pi^- N \rightarrow \pi^- \pi^0 \pi^0 N$
- Methods of the analysis:
 - The mass-independent PWA with established isobars
 - Isospin relations between $\pi^- \pi^- \pi^+$ and $\pi^- \pi^0 \pi^0$ amplitudes
 - Two parametrizations of PWA density matrix: rank=1 and unlimited rank with extracting the Largest-Eigenvalue-Eigenvector (LEV)
 - Resonance-model fits
 - Analysis with free parametrization of $\pi\pi$ -isobars
- Results of the analysis:
 - Mass-independent PWA for VES with established shapes of $\pi\pi$ -isobars, comparison of isospin relations between $\pi^- \pi^- \pi^+$ and $\pi^- \pi^0 \pi^0$ (VES and COMPASS)
 - Selected results for resonance model fits (COMPASS)
 - Comparison of PWA with rank=1 and unlimited rank with extracting the Largest-Eigenvalue-Eigenvector (VES)
 - Selected results for analysis with free parametrization of $\pi\pi$ -isobars (COMPASS)

The VES detector



The COMPASS detector



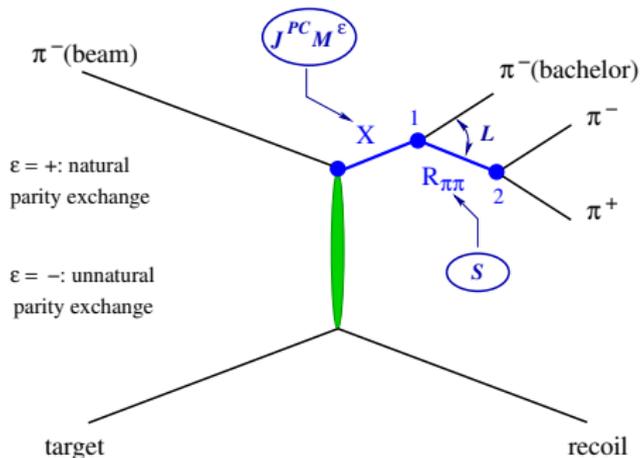
- VES

- $p_{beam} = 29 \text{ GeV}/c$
- Nucleus **Be** target, no detection of the recoil particle
- Momentum transferred squared $0 < t' < 1 \text{ GeV}^2/c^2$
- $\pi^- Be \rightarrow \pi^- \pi^- \pi^+ Be$ 87×10^6 events
- $\pi^- Be \rightarrow \pi^- \pi^0 \pi^0 Be$ 32×10^6 events

- COMPASS

- $p_{beam} = 190 \text{ GeV}/c$
- **LiH** target, Recoil Proton Detector
- Momentum transferred squared $0.1 < t' < 1 \text{ GeV}^2/c^2$
- $\pi^- p \rightarrow \pi^- \pi^- \pi^+ p$ 46×10^6 events
- $\pi^- p \rightarrow \pi^- \pi^0 \pi^0 p$ 3.5×10^6 events

π^- -beam dissociation on nucleon or nucleus target



- Reggeon exchange, naturality $\eta = P_R(-1)^J$
- Gottfried-Jackson frame: SCM of X: $Z_{GJ} \parallel \vec{p}_{beam}^*$, $Y_{GJ} = [\vec{p}_{recoil}^* \times \vec{p}_{beam}^*]$
- Reflectivity basis for system of mesons:
 $|JM^\epsilon\rangle = |JM\rangle - \epsilon P(-1)^{J-M} |J-M\rangle$
- At high beam energies: reflectivity ϵ equal to naturality η
- unpolarised target: $\epsilon = \pm 1$ states do not interfere

Brief introduction to mass-independent PWA

Mass-independent PWA events density:

$$\mathcal{I}(m, t', \tau) = \sum_{\varepsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{i,r}^{\varepsilon}(m, t') \psi_i^{\varepsilon}(m, \tau) \right|^2 + FLAT$$

The decay amplitudes $\psi_i^{\varepsilon}(\tau)$ are enumerated by their quantum numbers $i, \varepsilon = J^{PC} M^{\varepsilon}$ [isobar] πL and have no free parameters

Transition amplitudes $T_{i,r}^{\varepsilon}(m, t')$ fitted independently in each (m, t') - bin

Events density expressed through spin-density matrix:

$$\mathcal{I}(m, t', \tau) = \sum_{\varepsilon=\pm 1} \sum_{i,j} \rho_{i,j}^{\varepsilon}(m, t') \psi_i^{\varepsilon}(m, \tau) \psi_j^{\varepsilon*}(m, \tau), \quad \rho_{i,j}^{\varepsilon} = \sum_{r=1}^{N_r} T_{i,r}^{\varepsilon} T_{j,r}^{\varepsilon*}$$

For COMPASS data $N_r = 1$ is chosen, for VES - two models are tried: $N_r = 1$ and unlimited rank and extracting Eigenvector with Largest Eigenvalue

Isospin relations between $\pi^- \pi^- \pi^+$ and $\pi^- \pi^0 \pi^0$

Decay amplitudes for $I=1$ of 2π isobar (i.e. $\rho\pi$) in case of $I(3\pi) = 1$ are **connected**:

$$\sqrt{\frac{1}{2}} \left(\frac{1}{\sqrt{2}} [(\pi_{(1)}^- \pi^+) \pi_{(2)}^- + (\pi_{(2)}^- \pi^+) \pi_{(1)}^-] \right) \leftrightarrow -\sqrt{\frac{1}{2}} \left(\frac{1}{\sqrt{2}} [(\pi^- \pi_{(1)}^0) \pi_{(2)}^0 + (\pi^- \pi_{(2)}^0) \pi_{(1)}^0] \right)$$

Same dalitz-plot structure $\rightarrow N(\pi^- \pi^0 \pi^0):N(\pi^- \pi^- \pi^+) = 1:1$

Decay amplitudes for $I=0$ of 2π isobar (i.e. $f_0\pi$ or $f_2\pi$) channels are **always connected** :

$$\sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} [(\pi_{(1)}^- \pi^+) \pi_{(2)}^- + (\pi_{(2)}^- \pi^+) \pi_{(1)}^-] \right) \leftrightarrow -\sqrt{\frac{1}{3}} \left((\pi_{(1)}^0 \pi_{(2)}^0) \pi^- \right)$$

Different Dalitz-plot structure.

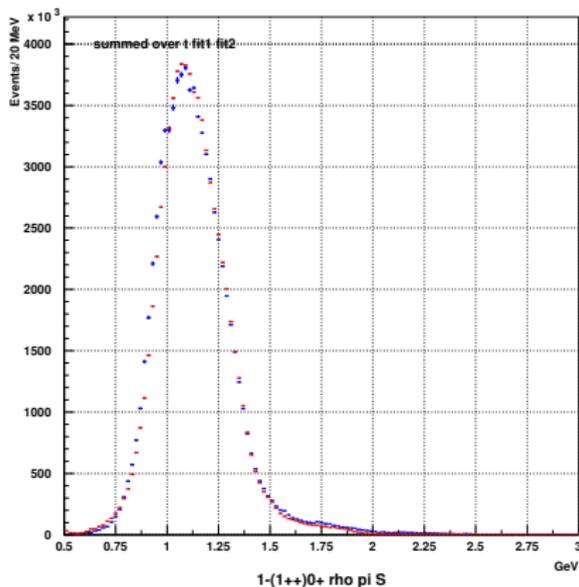
Narrow isobar and $m(3\pi) \gg m_{isob} + m_\pi \rightarrow N(\pi^- \pi^0 \pi^0):N(\pi^- \pi^- \pi^+) = 0.5:1$

Case of broad, overlapping isobars for $\pi^- \pi^- \pi^+$ - ratio can be significantly larger or smaller than **0.5**

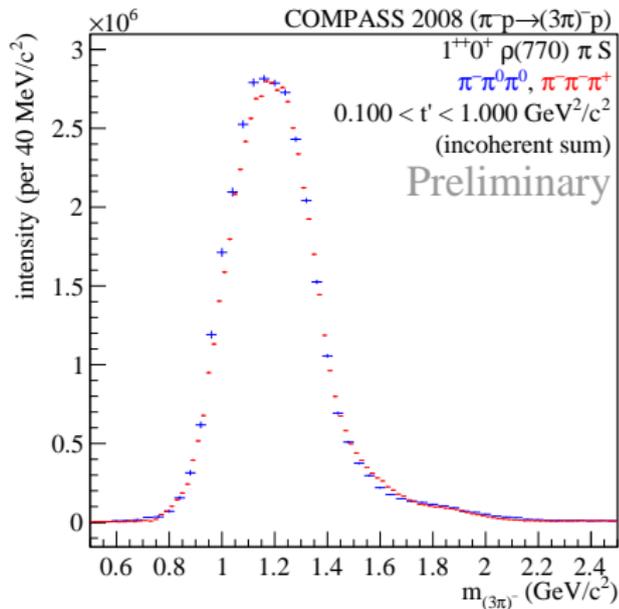
All corresponding **relative phases** in $\pi^- \pi^0 \pi^0$ and $\pi^- \pi^- \pi^+$ are equal - in case of appropriate choosing of directions of "spin analyzers" - π^- - direction in 2π center-of-mass for both systems

$J^{PC} M^{\epsilon} = 1^{++} 0^{+} \rho(770) \pi S$ -intensity

VES

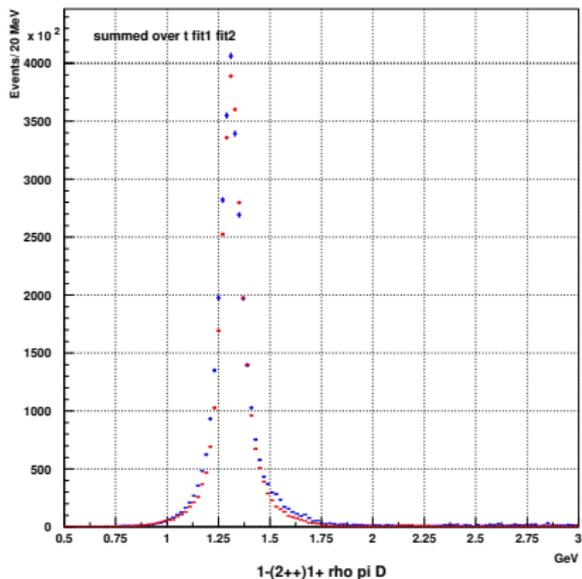


COMPASS

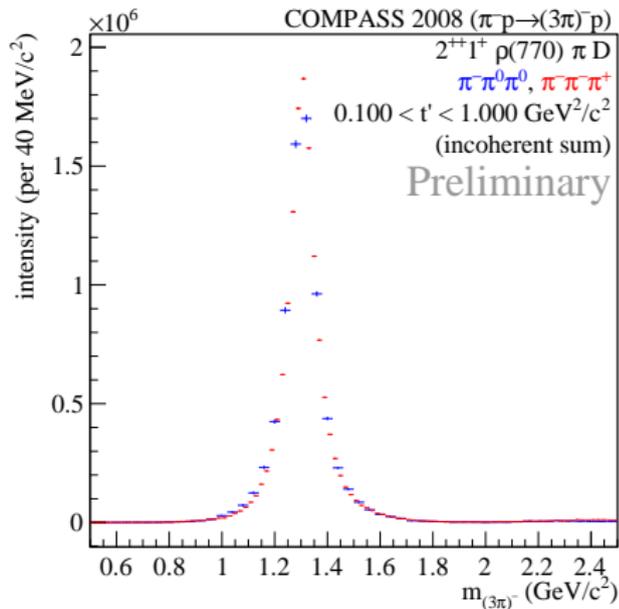


$J^{PC} M^{\epsilon} = 2^{++} 1^{+} \rho(770) \pi D$ -intensity

VES

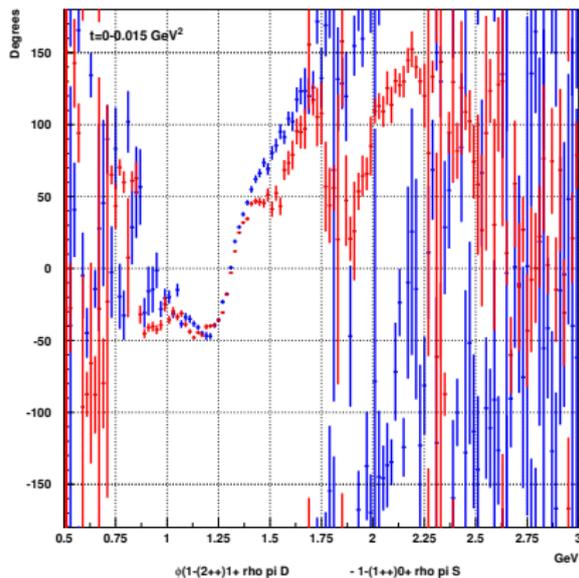


COMPASS

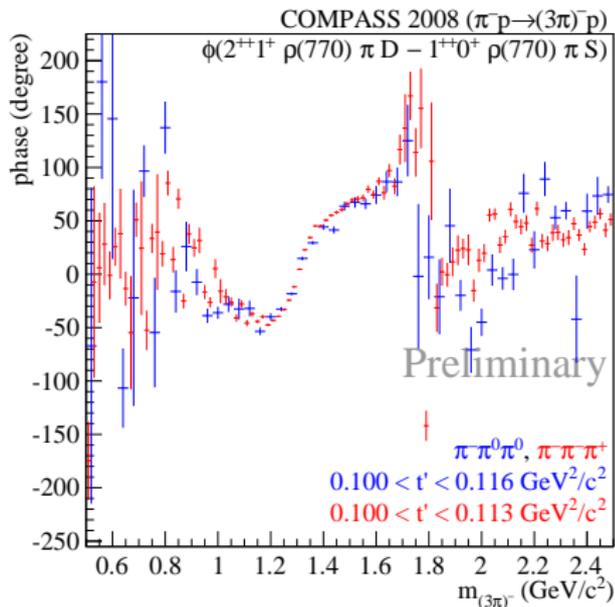


$J^{PC} M^{\epsilon} = 2^{++} 1^{+} \rho(770) \pi D$ -phase relative to $1^{++} 0^{+} \rho \pi S$

VES



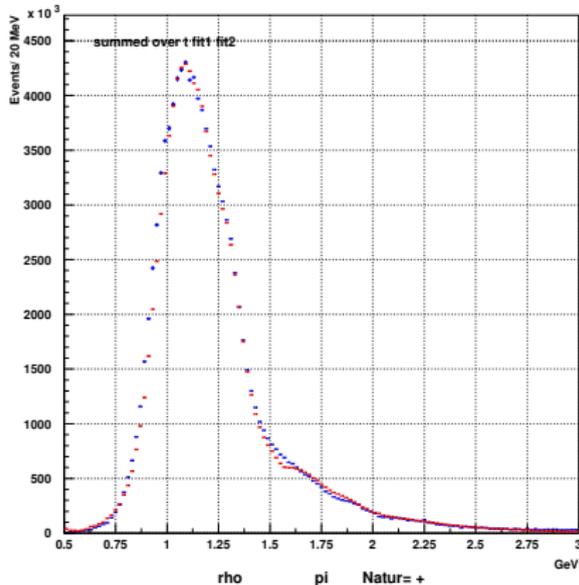
COMPASS



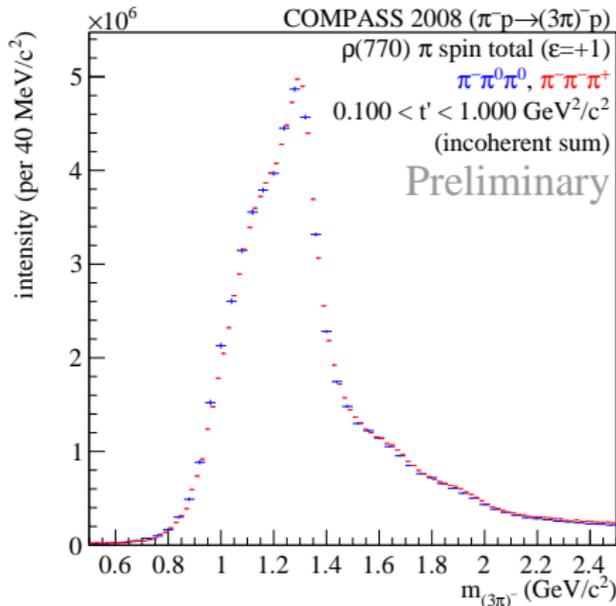
The relative phases between waves containing 3π -resonances do not depend on beam energy and t'

The total intensity of $\rho(770)\pi$ with $\varepsilon = +1$

VES



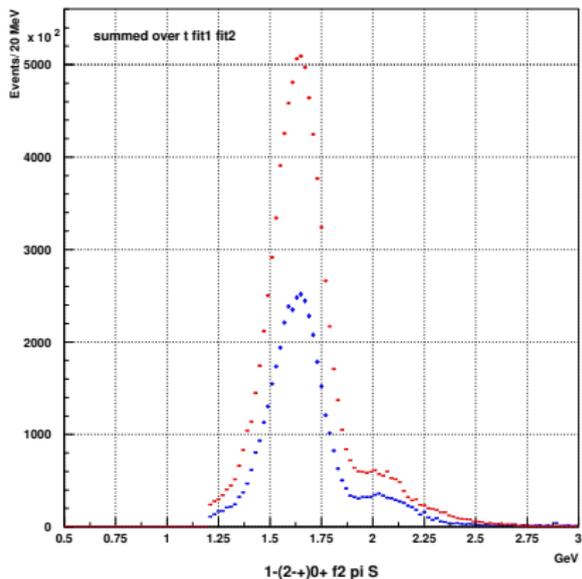
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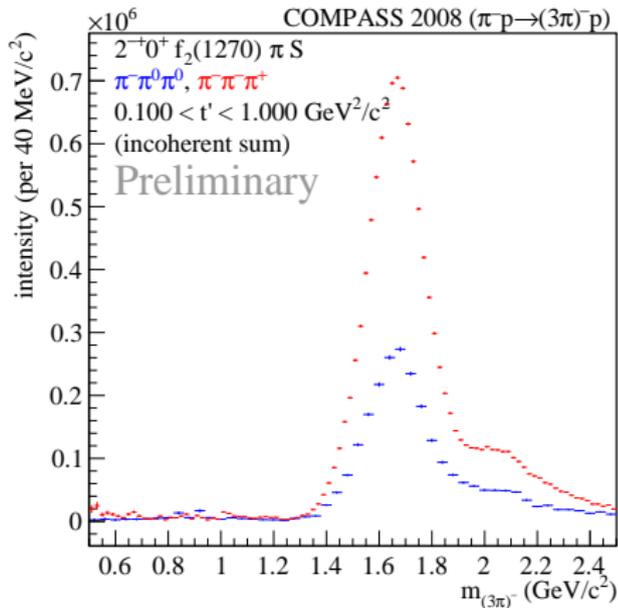
The $\varepsilon = +1$ dominates over $\varepsilon = -1$ (latest also do not show isospin relations)

$J^{PC} M^{\epsilon} = 2^{-+} 0^{+} f_2(1270) \pi S$ -intensity

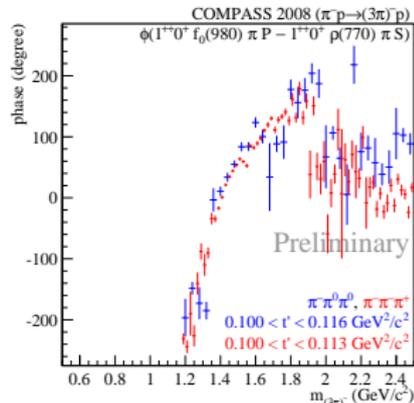
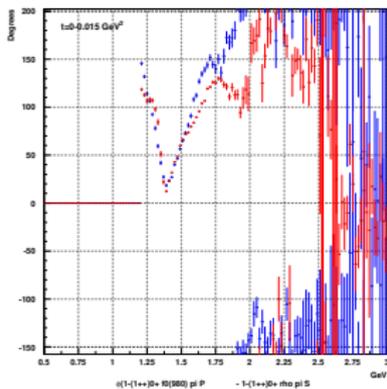
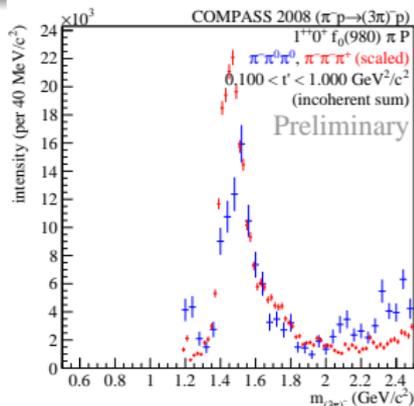
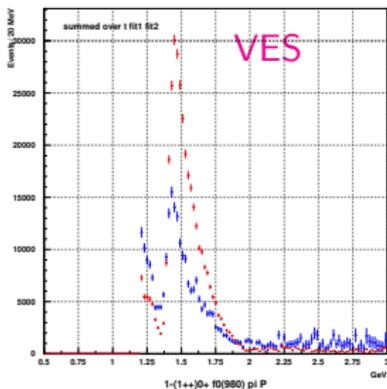
VES



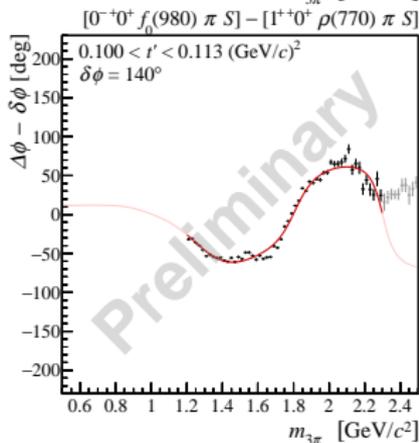
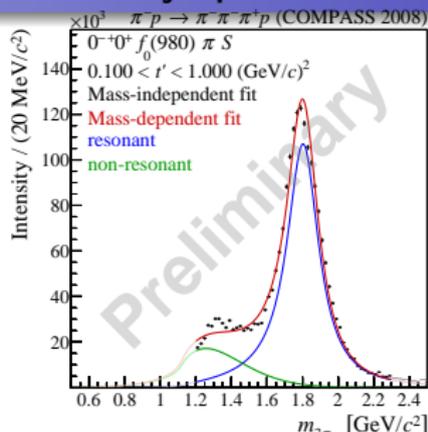
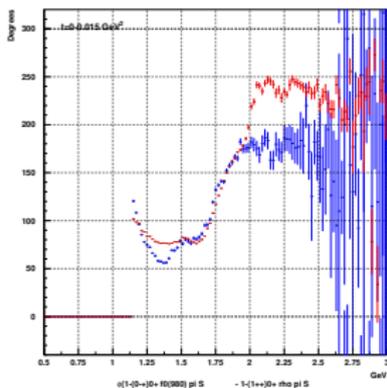
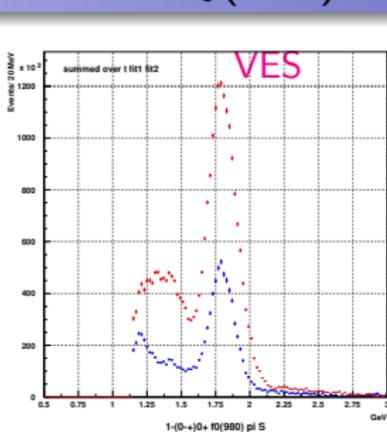
COMPASS



$J^{PC} M^{\epsilon} = 1^{++} 0^{+} f_0(980) \pi P$ -intensity, phase

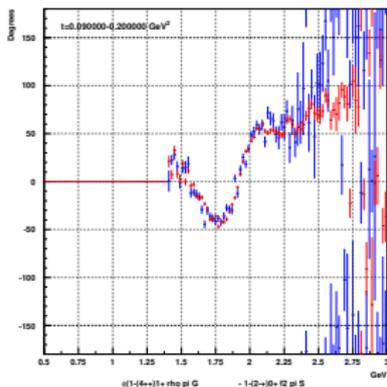
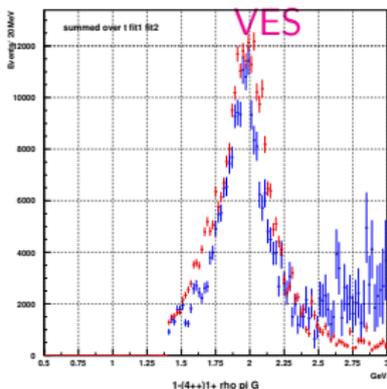


$J^{PC} M^{\epsilon} = 0^{-+}0^{+} f_0(980) \pi S$ -intensity, phase

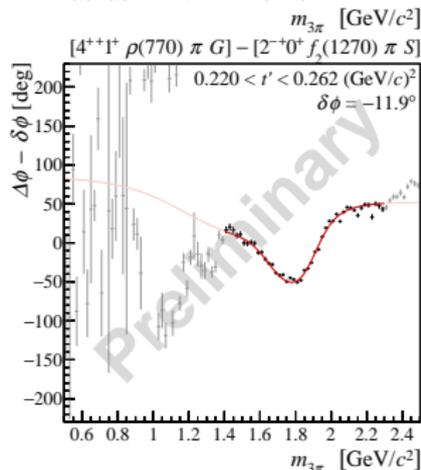
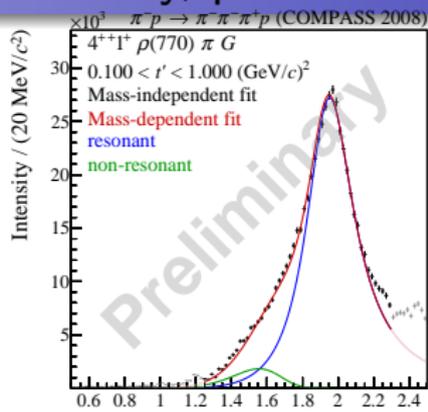


Clear $\pi(1800)$ signal, no $\pi(1300)$ -signal

$J^{PC} M^{\epsilon} = 4^{++}1^{+} \rho(770) \pi G$ -intensity, phase



$$M = 1980 \pm 10, \Gamma = 300 \pm 40 \text{ MeV}/c^2$$



$$M = 1935^{+11}_{-13}, \Gamma = 0.333^{+10}_{-21} \text{ MeV}/c^2$$

Unlimited rank density matrix approach

Coherent part of the density matrix R is the largest part of the matrix which has rank 1 and behaves like vector of amplitudes. Let

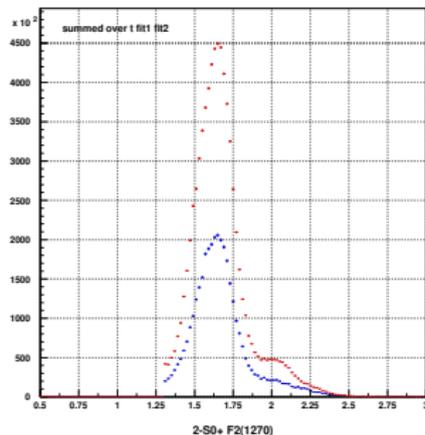
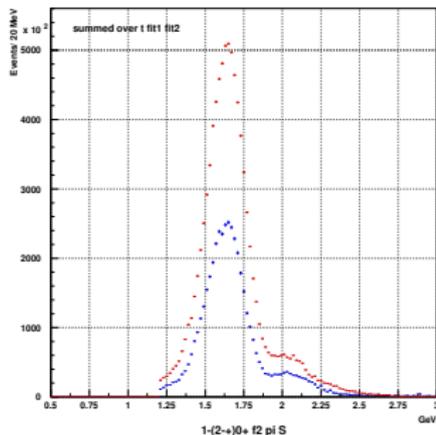
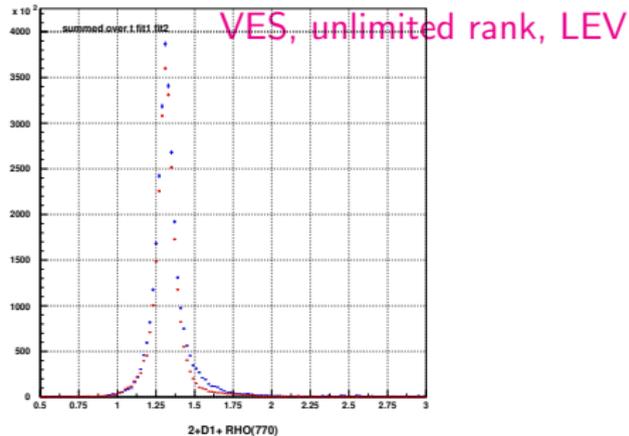
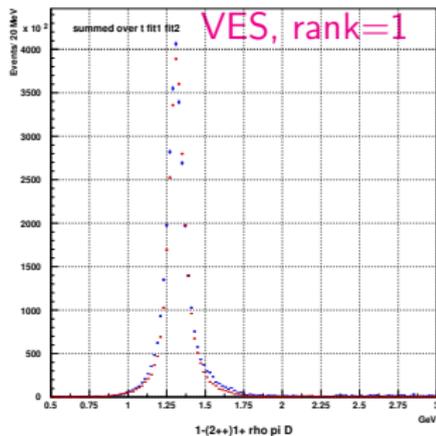
$$R = \sum_{k=1}^d e_k * V_k * V_k^+ \quad \text{where} \quad \begin{cases} e_k \text{ is } k\text{-th eigenvalue} \\ V_k \text{ is } k\text{-th eigenvector} \end{cases}$$

Let $e_1 \gg e_2 > \dots > e_d > 0$. Leading term R_L is coherent part of density matrix and R_S is the rest (incoherent part). This decomposition is stable w.r.t. variations of R matrix elements.

$$R = R_L + R_S, \quad R_L = e_1 * V_1 * V_1^+, \quad R_S = \sum_{k=2}^d e_k * V_k * V_k^+$$

Experience shows that resonances tend to concentrate in R_L .

VES: rank=1 (left) and unlimited rank, LEV (right)



PWA with fixed shapes of isobars vs. freed isobares

Decay amplitude with established isobars:

The decay amplitude $\psi_i^\varepsilon(\tau)$ contains angular part and $\pi^- \pi^+$ isobar Breit-Wigner function and is bose-symmetrized by swapping (1) \leftrightarrow (3) in $\pi_{(1)}^- \pi_{(2)}^+ \pi_{(3)}^-$ system:

$$\psi_i^\varepsilon(\tau) = A_i^\varepsilon(\Omega_{12}, \Omega_1^*) BW_{j(i)}(m_{12}) + A_i^\varepsilon(\Omega_{32}, \Omega_3^*) BW_{j(i)}(m_{32})$$

Decay amplitudes with freed isobars:

The fixed amplitude of $\pi^- \pi^+$ isobar is replaced by sum of step-like functions with complex coefficients: $BW(m)_j \rightarrow \sum_\beta \omega_{j,\beta} \Pi_\beta(m)$

In that case [isobar] $\rightarrow (\pi\pi)_s$ and wave notation is $J^{PC} M^\varepsilon (\pi\pi)_s \pi L$

The full free-isobarred amplitude for $J^{PC} M^\varepsilon$ sector:

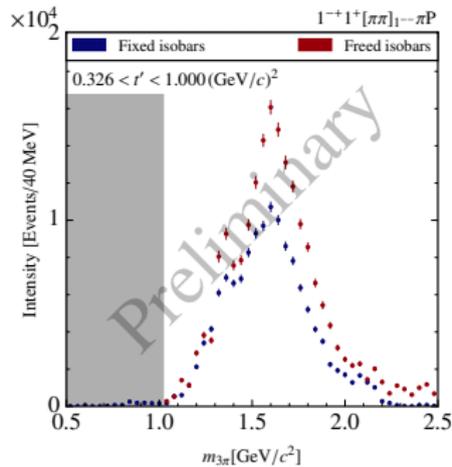
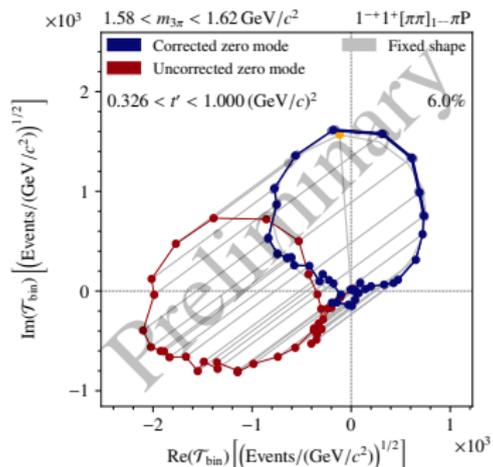
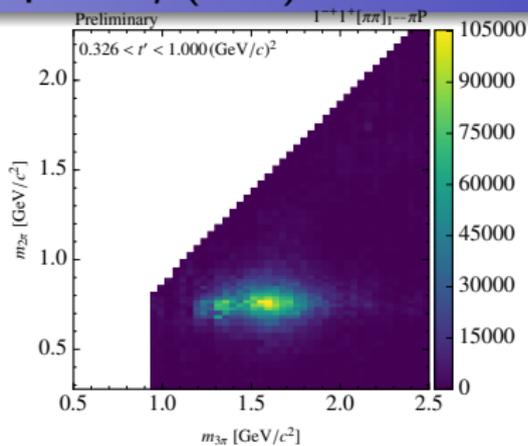
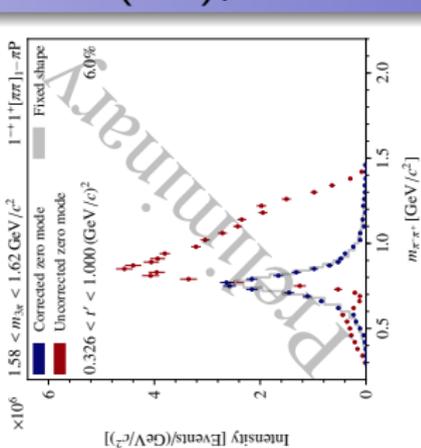
$$\begin{aligned} F_{J^P M^\varepsilon}(\tau) &= \sum_k \sum_\beta \omega_{k,\beta} [A_{J^P M^\varepsilon, k}^\varepsilon(\Omega_{12}, \Omega_1^*) \Pi_\beta(m_{12}) + A_{J^P M^\varepsilon, k}^\varepsilon(\Omega_{32}, \Omega_3^*) \Pi_\beta(m_{32})] \\ &= \sum_k \sum_\beta \omega_{k,\beta} \hat{\Psi}_{J^P M^\varepsilon, k, \beta}(\tau) \end{aligned}$$

where k sums over different L, s for fixed $J^P M^\varepsilon$

We found linear dependences inside the set of free-isobaric decay amplitudes $\hat{\Psi}_{J^P M^\varepsilon, k, \beta}(\tau)$, called **zero modes**.

For two amplitudes: $0^{-+}(\pi\pi)_S \pi S$ and $0^{-+}(\pi\pi)_P \pi P$ - one real-valued function of zero mode found. For one amplitude: $1^{-+}(\pi\pi)_P \pi P$ - one zero mode.

$J^{PC} = 1^{-+}(\pi\pi)_{\rho}\pi$ - free shape of $\rho(770)$ COMPASS



CONCLUSIONS, OUTLOOK

- Analysis of 3π states in VES and COMPASS shows the dominance of diffractive production mechanism for both beam energies and different t' -ranges:
 - The positive reflectivity dominates for both beam energies
 - The isospin relations between $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$ demonstrate 1:1 and for $f\pi$ have 0.5:1
 - Relative phases match for corresponding pairs of waves in $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$
 - The relative phases between resonant 3π waves do not depend on beam energy and on t'
- VES has compatible statistics of $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$ and enhanced production of $J^{PC}M^E = 1^{++}0^+$ states:
 - perspective study of $a_1(1420)$ -phenomenon in both 3π -final states
 - perspective to perform "free-isobarred" analysis in both 3π -final states
- COMPASS analysis with freed isobars was first time performed for $J^{PC}M^E = 1^{-+}1^+(\pi\pi)_P\pi P$:
 - The continuous ambiguities were resolved by applying Breit-Wigner model for $(\pi\pi)_P$ freed amplitude
 - The obtained $(\pi\pi)_P$ model-independent amplitude is well described by $\rho(770)$ Breit-Wigner