



**PHIPSI19**  
BINP, Novosibirsk



## R measurement at KEDR

*Korneliy Todyshev*  
**KEDR collaboration**

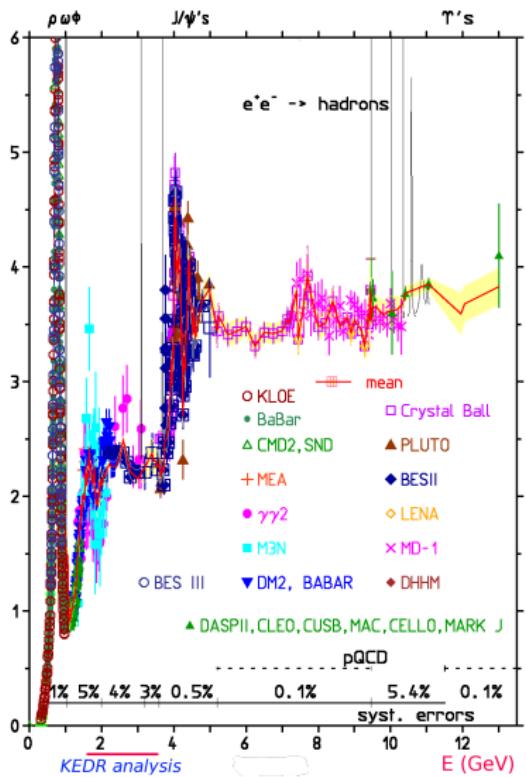
27 February 2019

**International Workshop on  $e^+e^-$  collisions from  $\phi$  to  $\psi$**

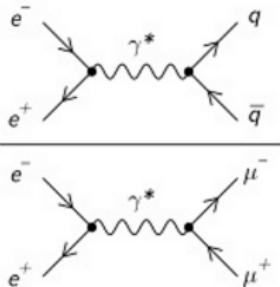


# $R(s)$ measurement. Motivation.

"Priore loco" (In first)



$$R = \frac{\sigma(e^-e^+ \rightarrow \text{hadrons})}{\sigma(e^-e^+ \rightarrow \mu^-\mu^+)} \approx$$



In first approximation:

$$R(s) \simeq 3 \sum e_q^2$$

$R(s)$  is used to determine:

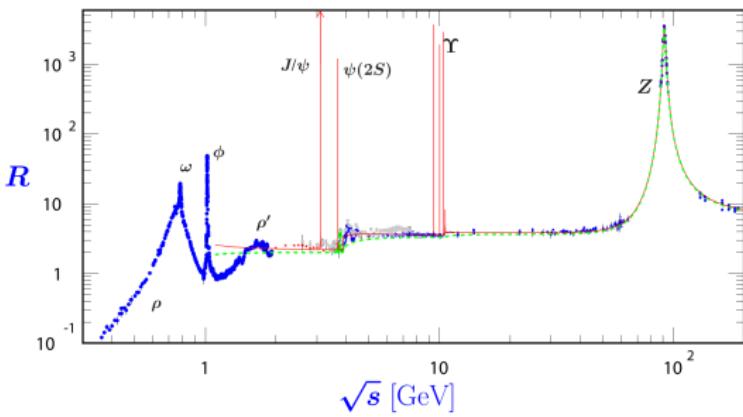
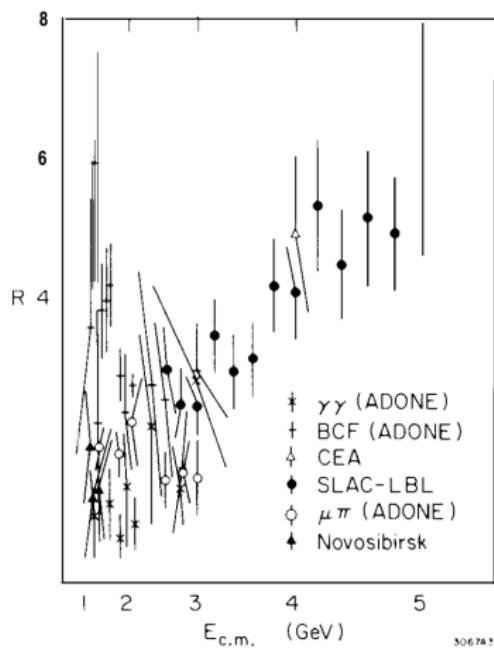
- $\alpha_s(s)$
- $(g_\mu - 2)/2$
- $\alpha(M_Z^2)$

F. Jegerlehner arXiv:1511.0447



# $R(s)$ measurement. Experimental data.

"Confer!" (Compare!)

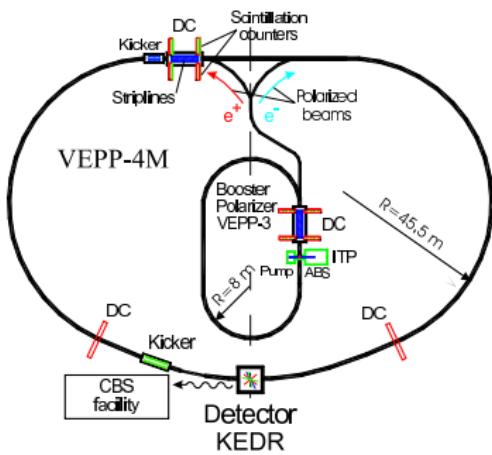


PDG data at the present time.

"The ratio  $R$  as of July 1974"  
Presented by Richter at the  
London conference in July 1974.



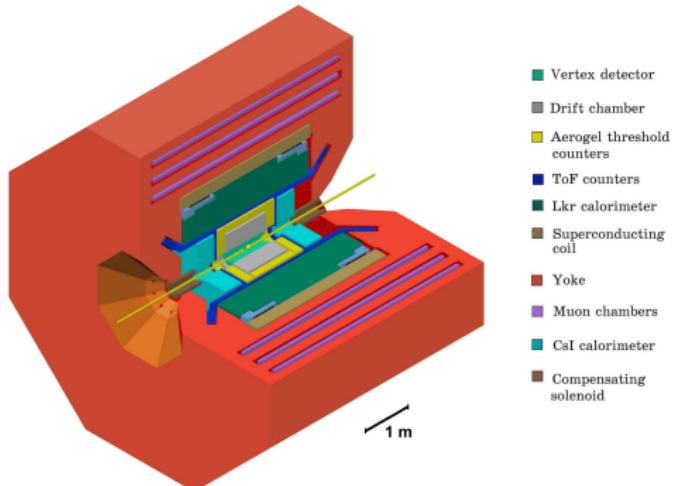
"Actum atque tractatum" (done and discussed)



Beam energy	$1 \div 5 \text{ GeV}$
Number of bunches	$2 \times 2$
Luminosity	$1.8 \text{ GeV}$ $1.5 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$

### Energy measurement:

- Resonant depolarization method:  
Instant measurement accuracy  $\sim 1 \text{ keV}$   
Energy interpolation accuracy  $10 \div 30 \text{ keV}$
- Compton backscattering method  $\sim 100 \text{ keV}$



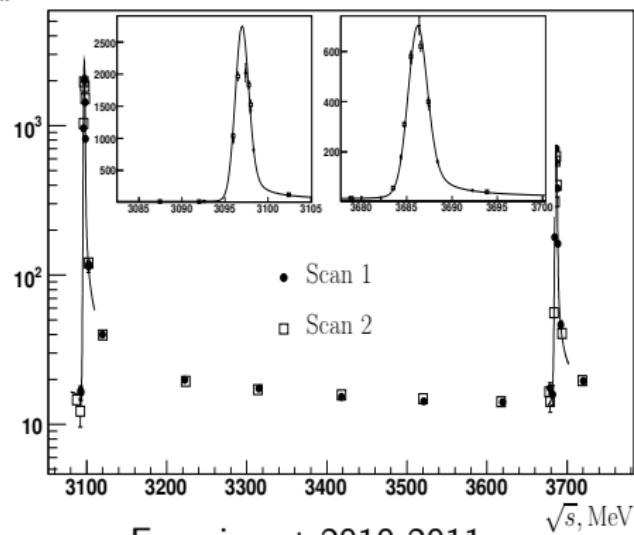


# R measurement between $J/\psi$ and $\psi(2S)$

"Consideratio naturae" (contemplation of nature)

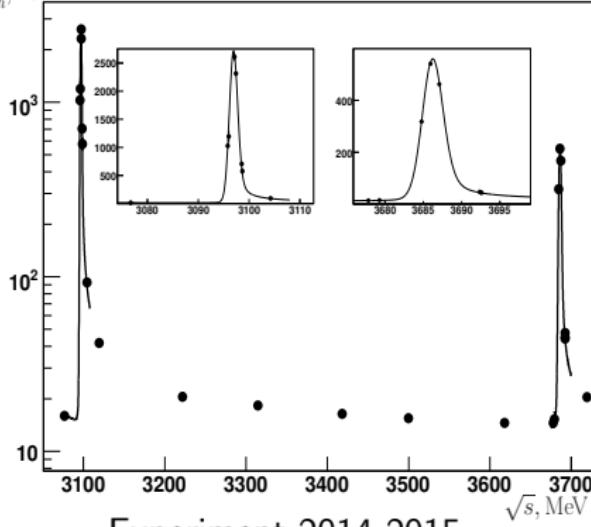
The observed multihadron cross section as a function of the c.m. energy

$\sigma_{mh}^{obs}$ , nb



Experiment 2010-2011

$\sigma_{mh}^{obs}$ , nb



Experiment 2014-2015

- The c.m. energy range between 3.076 and 3.72 GeV studied
- An integrated luminosity of  $2.7 \text{ pb}^{-1}$  collected at 9 energies 3.077, 3.120, 3.223, 3.315, 3.418, 3.500, 3.521, 3.618, 3.719 GeV
- $\sim (2 - 6) \times 10^3$  m.h. events per point,  $\sim 38 \times 10^3$  in total



"Modus operandi" (procedure; method of operating)

The way that we are measuring  $R$ :

$$R = \frac{\sigma_{obs}(s) - \sum \varepsilon_{\psi}^{tail}(s)\sigma_{\psi}^{tail}(s) - \sum \varepsilon_{bg}^i(s)\sigma_{bg}^i(s)}{\varepsilon(s)(1 + \delta(s))\sigma_{\mu\mu}^0}$$

with  $\sigma_{obs}(s) = \frac{N_{mh} - N_{res.bg.}}{\int \mathcal{L} dt}$ . where  $N_{mh}$  represent all events pass hadronic selection criteria,  $N_{res.bg.}$  – residual machine background

$\sum \varepsilon_{\psi}^{tail}(s)\sigma_{\psi}^{tail}(s)$  is contribution from  $J/\psi$  and  $\psi(2S)$  resonances

$\sum \varepsilon_{bg}^i(s)\sigma_{bg}^i(s)$  is contribution from physical processes:  $e^+e^- \rightarrow l^+l^-$ ,  $\gamma\gamma$ -processes.

$\varepsilon(s)$  – multihadron efficiency.

$$1 + \delta(s) = \int dx \frac{1}{1-x} \frac{\mathcal{F}(s,x)}{|1 - \tilde{\Pi}(s(1-x))|^2} \frac{\tilde{R}(s(1-x))\varepsilon(s(1-x))}{R(s)\varepsilon(s)}$$

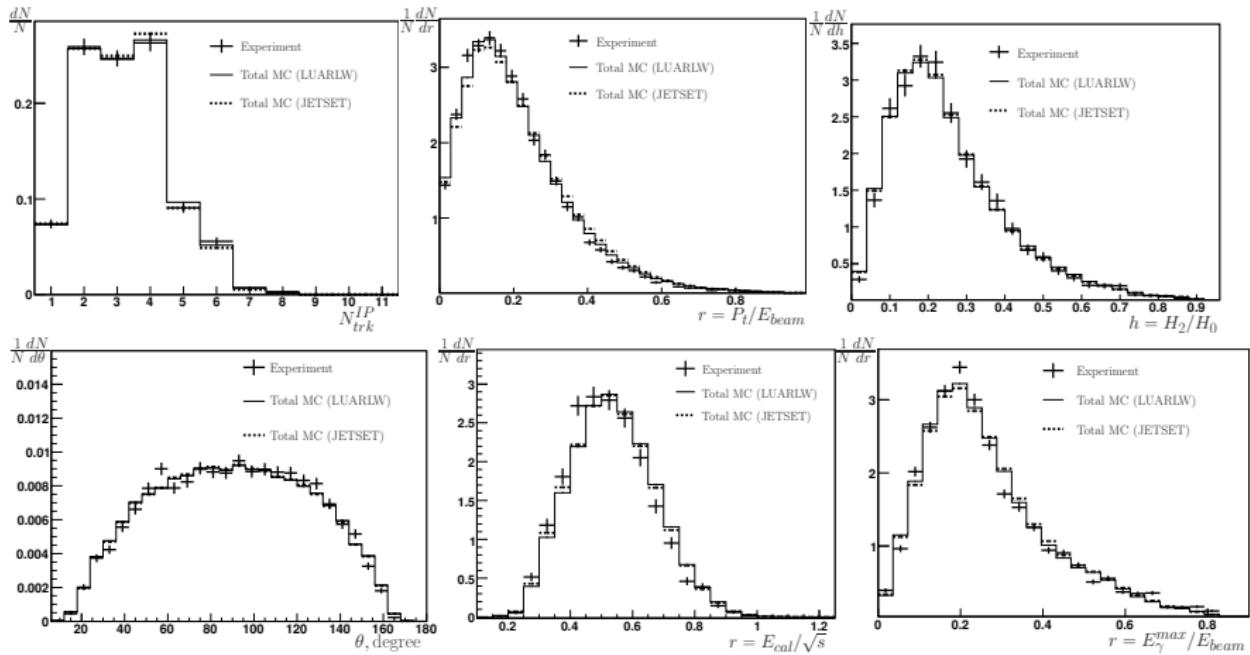
$\mathcal{F}(s,x)$  – radiative correction kernel ([E.A.Kuraev, V.S.Fadin](#)

[Sov.J.Nucl.Phys.41\(466-472\)1985](#)) Here  $\tilde{\Pi}$  and  $\tilde{R}$  does not includes  $J/\psi$  and  $\psi(2S)$  resonances. To determine the contributions of the  $J/\psi$  and  $\psi(2S)$  without external data, the additional data samples of about  $0.4 \text{ pb}^{-1}$ (2010-2011) and  $0.34 \text{ pb}^{-1}$ (2014-2015) were collected in the vicinity of peak regions.



# Simulation: JETSET and LUARLW

*"Punctum saliens" (The most important thing)*



Properties of hadronic events produced in the uds continuum at 3.119 GeV (2014-2015).

Here  $N$  is the number of events,  $N_{IP}^{trk}$  is the number of tracks originated from IP,  $P_t$  is a transverse momentum of the track,  $H_2$  and  $H_0$  are Fox-Wolfram moments,  $\theta$  is a polar angle of the track,  $E_{cal}$  is energy deposited in the calorimeter,  $E_\gamma^{\max}$  is energy of the most energetic photon.



# Systematic uncertainties

*"Satius est supervacua discere quam nihil"*  
(Better to learn more than necessary than nothing at all)

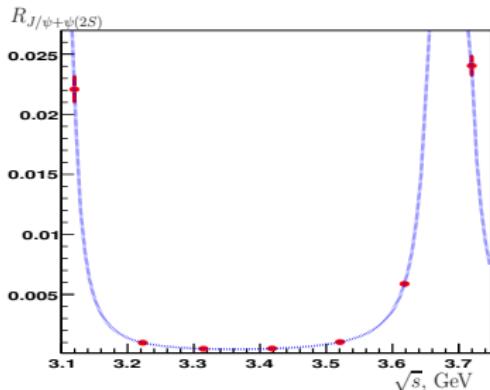
Seneca

Source	Syst. uncertainty, %		
	Scan 1 and 2 (2010-2011)	Scan 2014-2015	Correlated
Luminosity	1.1	0.9	0.4
Rad. corr.	$0.4 \div 0.6$	$0.5 \div 0.8$	$0.2 \div 0.4$
<i>uds</i> simulation	$1.3 \div 2.0$	1.1	0.9
Track reconstruction	0.5	0.4	–
$J/\psi$	$0.1 \div 2.7$	$0.1 \div 1.8$	–
$\psi(2S)$ (at 3.72 GeV)	1.4	1.1	–
$I^+I^-$	$0.1 \div 0.2$	$0.3 \div 0.4$	$0.1 \div 0.2$
$e^+e^-X$	$0.1 \div 0.2$	0.1	0.1
Trigger	0.2	0.2	0.2
Nuclear interaction	0.2	0.2	0.2
Machine background	$0.5 \div 1.1$	$0.4 \div 0.8$	–
Cuts	0.6	0.6	–
Total	$2.1 \div 3.6$ (correlated $1.8 \div 2.5$ )	$1.9 \div 2.7$	1.1



# $R$ for $\sqrt{s} = 3.12 - 3.72$ GeV

"Restitutio in integrum" (Restoring the thing to its original state)



Using  $J/\psi$  and  $\psi(2S)$  parameters, we obtain  $R_{uds}(s) + R_{J/\psi + \psi(2S)} \implies R(s)$

Data 2010-2011		Data 2014-2015		Combination	
$\sqrt{s}$ , MeV	$R_{uds}(s)$	$\sqrt{s}$ , MeV	$R_{uds}(s)$	$\sqrt{s}$ , MeV	$R_{uds}(s)\{R(s)\}$
-	-	$3076.7 \pm 0.2$	$2.188 \pm 0.056 \pm 0.042$	$3076.7 \pm 0.2$	$2.188 \pm 0.056 \pm 0.042$
$3119.9 \pm 0.2$	$2.215 \pm 0.089 \pm 0.066$	$3119.2 \pm 0.2$	$2.211 \pm 0.046 \pm 0.060$	$3119.6 \pm 0.4$	$2.212\{2.235\} \pm 0.042 \pm 0.049$
$3223.0 \pm 0.6$	$2.172 \pm 0.057 \pm 0.045$	$3221.8 \pm 0.2$	$2.214 \pm 0.055 \pm 0.042$	$3222.5 \pm 0.8$	$2.194\{2.195\} \pm 0.040 \pm 0.035$
$3314.7 \pm 0.7$	$2.200 \pm 0.056 \pm 0.043$	$3314.7 \pm 0.4$	$2.233 \pm 0.044 \pm 0.042$	$3314.7 \pm 0.6$	$2.219\{2.219\} \pm 0.035 \pm 0.035$
$3418.2 \pm 0.2$	$2.168 \pm 0.050 \pm 0.042$	$3418.3 \pm 0.4$	$2.197 \pm 0.047 \pm 0.040$	$3418.3 \pm 0.3$	$2.185\{2.185\} \pm 0.032 \pm 0.035$
-	-	$3499.6 \pm 0.4$	$2.224 \pm 0.054 \pm 0.040$	$3499.6 \pm 0.4$	$2.224\{2.224\} \pm 0.054 \pm 0.040$
$3520.8 \pm 0.4$	$2.200 \pm 0.050 \pm 0.044$	-	-	$3520.8 \pm 0.4$	$2.200\{2.201\} \pm 0.050 \pm 0.044$
$3618.2 \pm 1.0$	$2.201 \pm 0.059 \pm 0.044$	$3618.1 \pm 0.4$	$2.220 \pm 0.049 \pm 0.042$	$3618.2 \pm 0.7$	$2.212\{2.218\} \pm 0.038 \pm 0.035$
$3719.4 \pm 0.7$	$2.187 \pm 0.068 \pm 0.060$	$3719.6 \pm 0.2$	$2.213 \pm 0.047 \pm 0.049$	$3719.5 \pm 0.5$	$2.204\{2.228\} \pm 0.039 \pm 0.042$

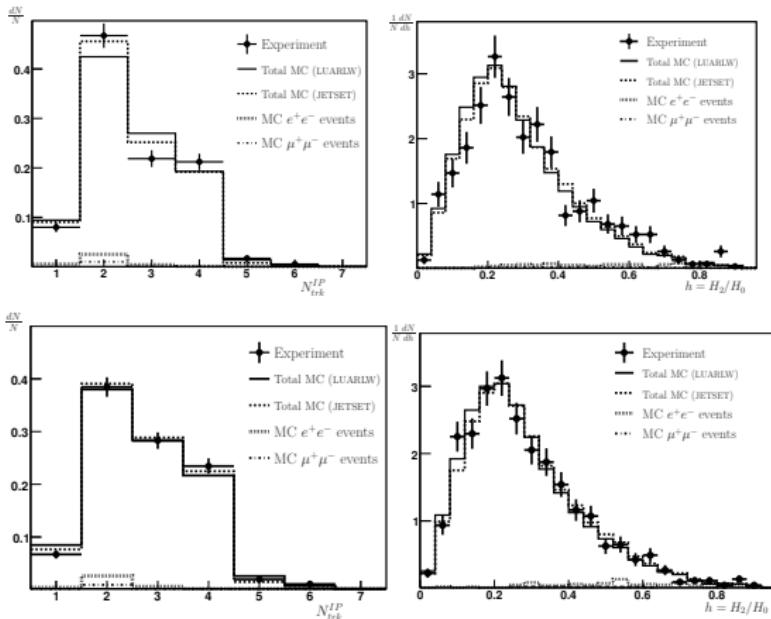
V.V.Anashin et al., Phys.Lett. B 753, 533-541 (2016).[arXiv:1510.02667]

V.V.Anashin et al., Phys.Lett. B 788, 42-51 (2019).[arXiv:1805.06235]

"Fac simile" (to make alike; reproduction exacte)

- An integrated luminosity  $0.66 \text{ pb}^{-1}$  collected at 13 equidistant points with a step  $\sim 0.1 \text{ GeV}$ :  $1.841, 1.937 \dots 3.048 \text{ GeV}$
- $\sim 10^3$  hadronic events per point,  $14.8 \times 10^3$  events in total
- Simulation of the  $uds$  continuum based on the LUARLW generator, tuned JETSET alternatively used at 6 points for a cross-check.

Experimental distribution and two variants of MC simulation based on LUARLW and tuned JETSET are plotted ( $\sqrt{s} = 1.94 \text{ GeV}$  and  $\sqrt{s} = 2.14 \text{ GeV}$ ).





"*Ut supra*" (as (described) above)

Measured value of  $R = \frac{\sigma_{obs}(s) - \sum \varepsilon_{bg}^i(s) \sigma_{bg}^i(s)}{\varepsilon(s)(1+\delta(s))\sigma_{\mu\mu}^0}$

$\sqrt{s}$ , MeV	$R(s)$
1841.0	$2.226 \pm 0.139 \pm 0.158$
1937.0	$2.141 \pm 0.081 \pm 0.073$
2037.3	$2.238 \pm 0.068 \pm 0.072$
2135.7	$2.275 \pm 0.072 \pm 0.055$
2239.2	$2.208 \pm 0.069 \pm 0.053$
2339.5	$2.194 \pm 0.064 \pm 0.048$
2444.1	$2.175 \pm 0.067 \pm 0.048$
2542.6	$2.222 \pm 0.070 \pm 0.047$
2644.8	$2.220 \pm 0.069 \pm 0.049$
2744.6	$2.269 \pm 0.065 \pm 0.050$
2849.7	$2.223 \pm 0.065 \pm 0.047$
2948.9	$2.234 \pm 0.064 \pm 0.051$
3048.1	$2.278 \pm 0.075 \pm 0.048$

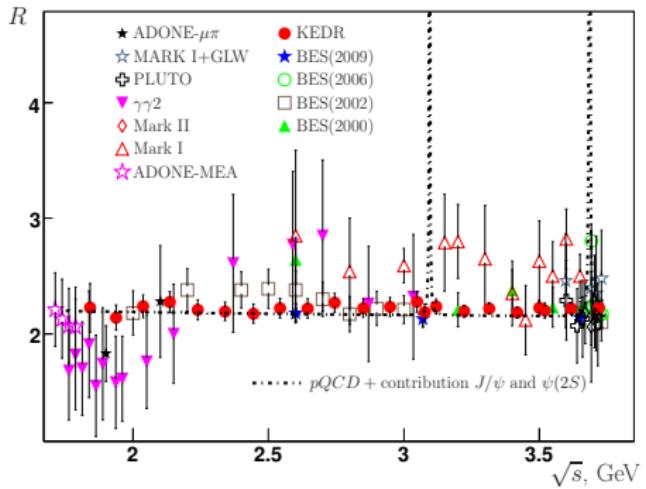
The main systematic  
uncertainties in the  $R$ :

Source	Error, %
Luminosity	1.2
Rad. corr.	$0.5 \div 2.0$
<i>uds</i> simulation	$1.2 \div 6.6$
$I^+I^-$	$0.3 \div 0.6$
$e^+e^-X$	0.2
Trigger	0.3
Nuclear interaction	0.4
Machine background	$0.4 \div 0.9$
Cuts	0.7
Total	$2.1 \div 7.1$



# Comparison with others experiments

*"De omnibus dubitandum"*(All is to be doubted)  
René Descartes



The quantity  $R$  versus the c.m. energy and the sum of the prediction of perturbative QCD and a contribution of narrow resonances.

In the c.m.energy range 3.08-3.72 GeV the weighted average

$\bar{R}_{uds} = 2.204 \pm 0.014 \pm 0.026$  is approximately one sigma higher than that theoretically expected,  $R_{uds}^{pQCD} = 2.16 \pm 0.01$  calculated according to the pQCD In the lower c.m.energy range 1.84-3.05 GeV the weighted average is  $2.225 \pm 0.020 \pm 0.047$  (the pQCD prediction of  $2.18 \pm 0.02$ ).



# An application of the $R(s)$

*"Natura appetit perfectum, ita est lex"*

(Nature desires perfection, so also does the law)

Correlated uncertainties of  $R_{uds}$  in %

Source	Uncertainty in %	
	Data 2010	Data 2010 / 2011,2014
<b>Luminosity</b>		
Cross section calc.	0.5	0.4
Calorimeter response	0.7	-
Calorimeter alignment	0.2	0.2
<b>Rad. correction</b>		
$\bar{n}$ approx.	0.3	0.1
$\delta R_{uds}(s)$	0.2	0.2
$\delta \epsilon(s)$	0.3	0.2
Continuum simulation	1.2	$0.4 \div 0.8$
Track reconstr.	0.5	0.4
$e^+e^-X$ contribution	0.2	0.1
$j/\psi$ contribution	0.3	0.2
Trigger efficiency	0.3	0.2
Nuclear interaction	0.4	0.2
<b>Sum in quadrature</b>	<b>1.8</b>	$0.8 \div 1.1$

$$R_{uds}(s) \simeq 2 \times \left( 1 + \frac{\alpha_s}{\pi} + \frac{\alpha_s^2}{\pi^2} \times \left( \frac{365}{24} - 9\zeta_3 - \frac{11}{4} \right) \right)$$

where  $\zeta$  is the Euler-Riemann zeta function,

$$\begin{aligned} \alpha_s(s) = & \frac{1}{b_0 t} \left( 1 - \frac{b_1 l}{b_0^2 t} + \frac{b_1(l^2 - l - 1) + b_0 b_2}{b_0^4 t^2} \right. \\ & \left. + \frac{b_1^3(-2l^3 + 5l^2 + 4l - 1) - 6b_0 b_2 b_1 l + b_0^2 b_3^2}{2b_0^6 t^3} \right) \end{aligned}$$

with  $t = \ln \frac{s}{\Lambda^2}$ ,  $l = \ln t$  parametrized in terms of the QCD scale parameter  $\Lambda$  and coefficients  $b_0, b_1, b_3$  (can be found in PDG). To determine  $\Lambda$ , we minimise the  $\chi^2$  function

$$\chi^2 = \sum_i \sum_j \left( R_{uds}^{\text{meas}}(s_i) - R_{uds}^{\text{calc}}(s_i) \right) C_{ij}^{-1} \left( R_{uds}^{\text{meas}}(s_j) - R_{uds}^{\text{calc}}(s_j) \right),$$

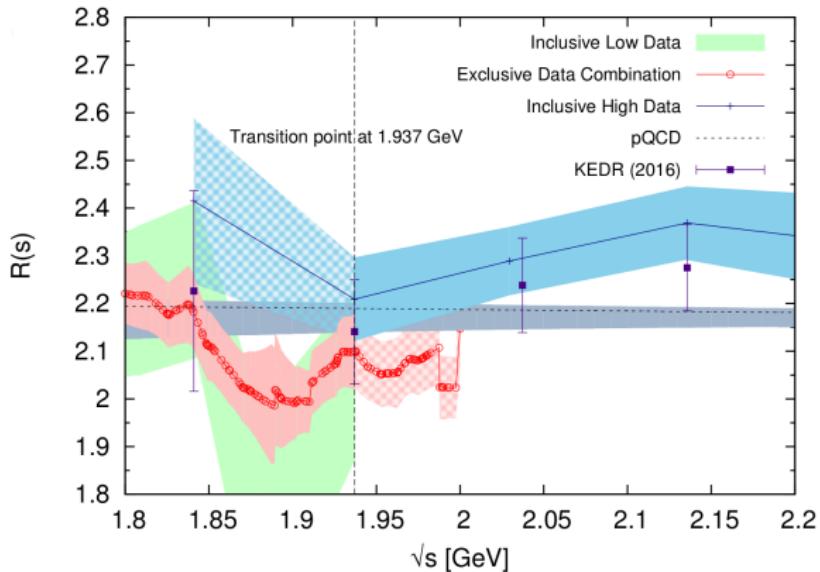
The obtained value of  $\Lambda = 0.361^{+0.155}_{-0.174}$  GeV corresponds to  $\alpha_s(m_\tau) = 0.332^{+0.100}_{-0.092}$ . If the next order of pQCD is included in the expansion of  $R_{uds}$ , the fitting results are as follows:  $\Lambda = 0.437^{+0.210}_{-0.215}$  GeV and  $\alpha_s(m_\tau) = 0.378^{+0.173}_{-0.120}$ .

$\alpha_s(m_\tau)$  determined from our  $R(s)$  results is consistent with obtained in semileptonic  $\tau$  decays ( $\alpha_s(m_\tau) = 0.331 \pm 0.013$ )



# Comparison with exclusive data

*"Ex uno discet omnes"* (From one thing you can discern all)



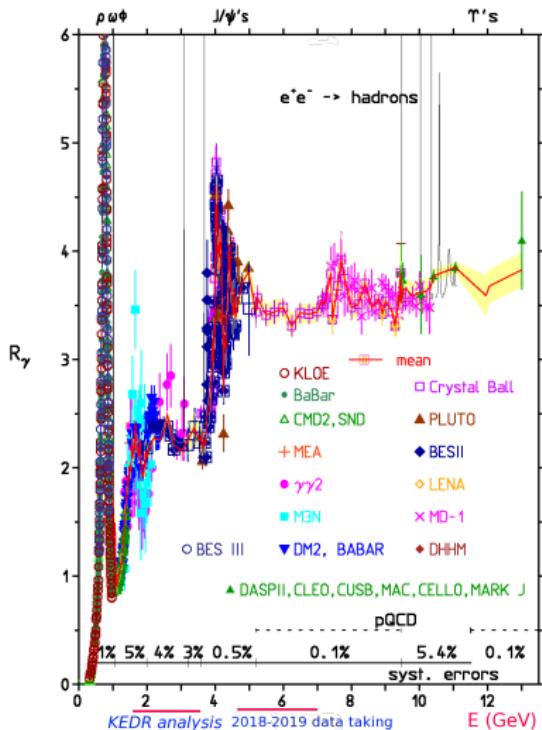
A. Keshavarzi, D. Nomura and T. Teubner.  
The muon  $g - 2$  and  $\alpha(M_Z^2)$ : a new data-based analysis.  
Phys. Rev. D **97**, 114025 (2018).[arXiv:1802.02995].



"Omne futurum incertum" (Every future thing is uncertain)

R measurement in the energy range  
4.56-6.96 GeV.

- First scan finished in 2018. An integrated luminosity  $\sim 4 \text{ pb}^{-1}$  collected at 8 equidistant points with a step  $\sim 0.3 \text{ GeV}$  from 4.71 to 6.81 GeV
- In 2019 we plan to start the second scan (10 equidistant points in the energy range  $4.56 \div 6.96 \text{ GeV}$ ).



F. Jegerlehner arXiv:1511.0447



*"Jucundi acti labores"* (past labors are pleasant)  
Cicero

- KEDR measured the  $R$  values at 22 center-of-mass energies between 1.84 and 3.72 GeV.  
In the energy range between 1.84 and 3.05 GeV the achieved accuracy is about or better than 3.9% at most of the energy points with a systematic uncertainty less than 2.4%.  
For the energies above  $J/\psi$  resonance the total error is about or better than 2.6% and a systematic uncertainty of about 1.9%.
- We plan to take data at the energy range from 4.56 to 6.96 GeV

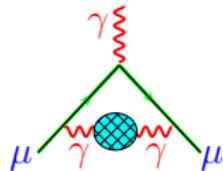
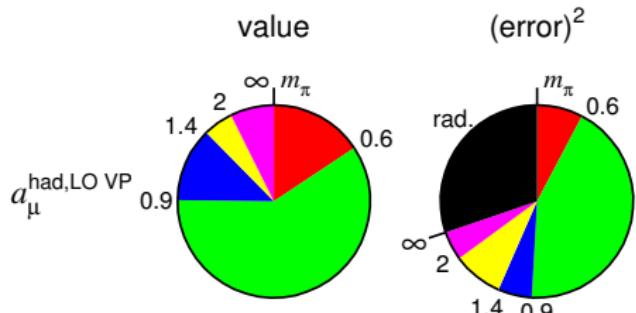
*"Gratia gratiam parit"*(thanks begets thanks)

Thank you for your time and  
attention

# BACKUP SLIDES

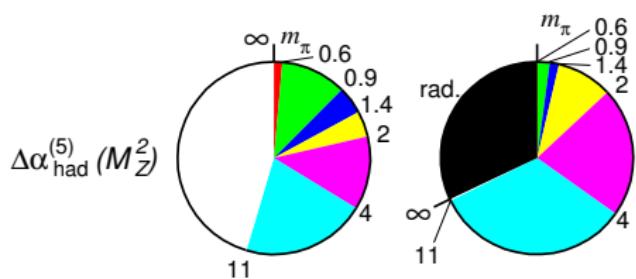
# $R$ contribution in $a_\mu$ and $\alpha(M_Z^2)$

$$a_\mu^{\text{exp}} = (g_\mu - 2)/2$$



$$a_\mu^{\text{LO VP}} = \frac{\alpha^2}{3\pi^2} \int_{m_\pi^2}^\infty \frac{K(s)R(s)}{s} ds$$

Low energy contributions dominate



$$\alpha(s) = \frac{\alpha}{1 - \Delta\alpha(s)}$$

$$\Delta\alpha = \sum_f \text{---} \gamma \text{---} \circlearrowleft \text{---} \gamma \text{---} = \Delta\alpha_{\text{lep}}(s) + \Delta\alpha_{\text{had}}(s)$$

$$\Delta\alpha^{(5)}(M_Z^2) = -\frac{\alpha M_Z^2}{3\pi} \operatorname{Re} \int_{m_\pi^2}^\infty \frac{R(s)ds}{s(s - M_Z^2 - i\epsilon)}$$

K.Hagiwara et al. arxiv:1105.3149

$\sigma^{e^+ e^- \rightarrow \text{hadrons}}$  and  $\sigma^{e^+ e^- \rightarrow e^+ e^-}$  nearby a narrow resonance

In the soft photon approximation analytical expression for the annihilation cross section nearby a narrow resonance.

Ya.I. Azimov et al. JETP Lett. 21 (1975) 172. With up-today modifications one has

$$\sigma^{e^+ e^- \rightarrow \text{hadr}}(s) = \sigma_{\text{continuum}}^{e^+ e^- \rightarrow \text{hadr}} + \frac{12\pi}{s} (1 + \delta_{sf}) \left[ \frac{\Gamma_{ee} \tilde{\Gamma}_h}{\Gamma M} \text{Im } f(s) - \frac{2\alpha \sqrt{R \Gamma_{ee} \tilde{\Gamma}_h}}{3\sqrt{s}} \lambda \text{Re } \frac{f^*(s)}{1 - \Pi_0} \right],$$

$$\begin{aligned} \left( \frac{d\sigma}{d\Omega} \right)^{ee \rightarrow ee} &= \left( \frac{d\sigma}{d\Omega} \right)_{\text{QED}}^{ee \rightarrow ee} + \frac{1}{s} (1 + \delta_{sf}) \left\{ \frac{9}{4} \frac{\Gamma_{ee}^2}{\Gamma M} (1 + \cos^2 \theta) \text{Im } f - \right. \\ &\quad \left. \frac{3\alpha}{2} \frac{\Gamma_{ee}}{M} \left[ (1 + \cos^2 \theta) \text{Re } \frac{f^*}{1 - \Pi_0(s)} - \frac{(1 + \cos \theta)^2}{(1 - \cos \theta)} \text{Re } \frac{f^*}{1 - \Pi_0(t)} \right] \right\}, \end{aligned}$$

Recently it was verified in the work X. Y. Zhou, Y. D. Wang and L. G. Xia, Chin. Phys. C 41 (2017) no.8, 083001

$$\delta = \frac{3}{4}\beta + \frac{\alpha}{\pi} \left( \frac{\pi^2}{3} - \frac{1}{2} \right) + \beta^2 \left( \frac{37}{96} - \frac{\pi^2}{12} - \frac{L}{72} \right), \quad L = \ln \left( s/m_e^2 \right), \quad \beta = \frac{2\alpha}{\pi} (L - 1),$$

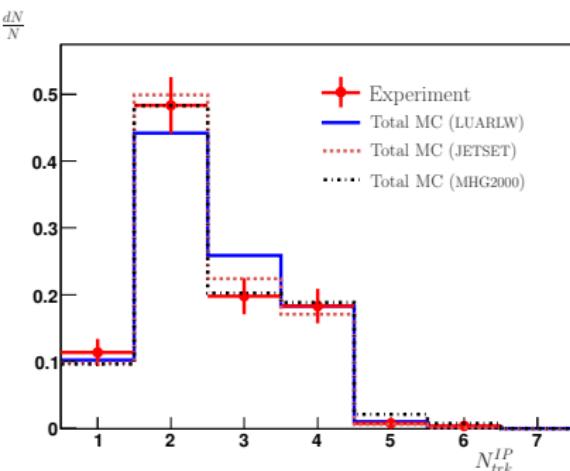
$$f(s) = \frac{\pi\beta}{\sin \pi\beta} \left( \frac{s}{M^2 - s - iM\Gamma} \right)^{1-\beta}$$

$\Gamma_{ee}$ ,  $\Gamma$ ,  $M$  – 'dressed' parameters including corrections to the vacuum polarization,  
 $\Gamma_{ee} = \Gamma_{ee}^{(0)}/|1 - \Pi_0|^2$ ,  $\lambda$ -parameter controls the resonance–continuum interference,  $\tilde{\Gamma}_h \neq \Gamma_h$

Numerical convolution with the collision energy distribution is used to fit resonance.

# Detection efficiency uncertainty in the energy range $\sqrt{s} = 1.84 \div 3.05$ GeV

- Used two essentially different MC generators (LUARLW and tuned JETSET)
- We validated our estimate of the systematic uncertainty related to simulation of the  $uds$  continuum using an unfolding method (Chinese Physics C Vol. 37, No. 6 (2013) 063001).
- The estimate at the most problematic energy point 1.84 GeV was additionally verified using the exclusive generator MHG2000.



Detection efficiency uncertainties obtained by different methods

Energy, MeV	$\delta\epsilon/\epsilon$		
	LUARLW	Unfolding JETSET method	LUARLW MHG2000
1841.0	6.6%	3.6%	3.8%
1937.0 $\div$ 2135.7	2.5%	1.9%	-
2135.7 $\div$ 3048.1	1.2%	0.5%	-

# Unfolding method

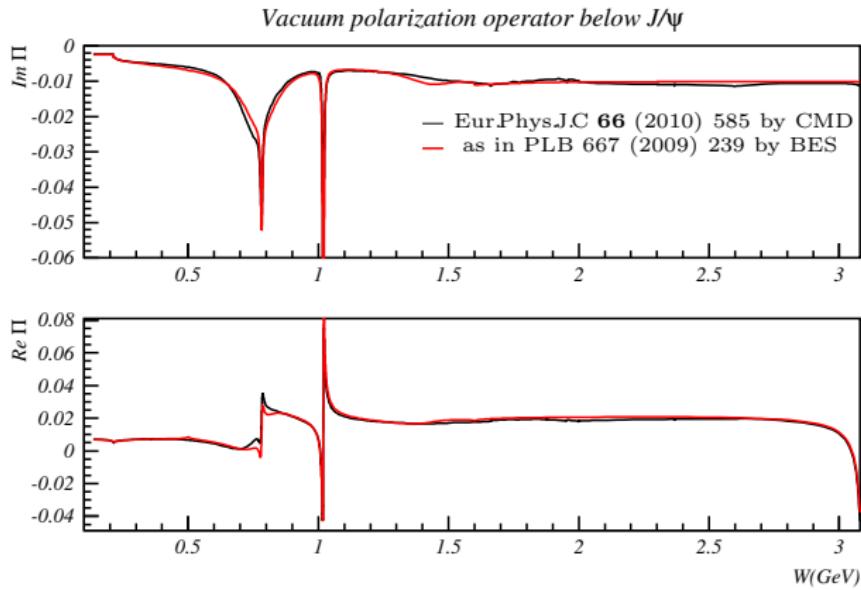
- An efficiency matrix  $\epsilon_{ij}$  describes the efficiency of an event generated with  $j$  charged tracks to be reconstructed with  $i$  charged tracks.
- The distribution of the number of observed charged track events in data,  $N_i^{obs}$ , is known. The true multiplicity distribution in data can be estimated from the observed multiplicity distribution in data and the efficiency matrix by minimizing the  $\chi^2$ .
- 

$$\chi^2 = \sum_{i=1}^{i=8} \frac{N_i^{obs} - \sum_{j=1}^{j=8} \epsilon_{ij} \times N_j}{N_i^{obs}}$$

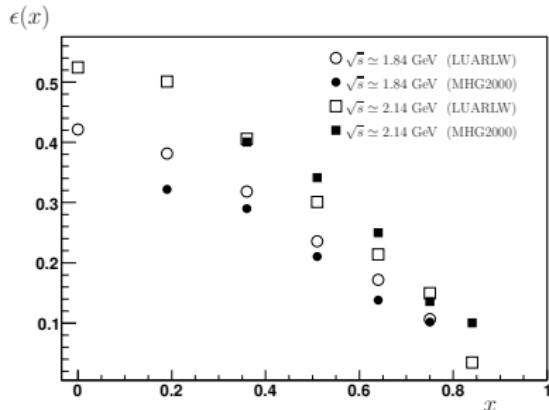
where the  $N_j$  ( $j = 0, 2, 4, 6, 8$ ) describe the true multiplicity distribution in data and are taken as floating parameters in the fit.

- The total «true» number of events in data can be obtained by summing all fitted  $N_j$ .

# $\Pi(s)$ calculation



# Radiation correction calculation in the energy range 1.84 – 3.05 GeV



Detection efficiency vs variable  $x$  at 1.84 and 2.14 GeV.

$$\mathbf{1} + \delta(s) = \int \frac{dx}{\mathbf{1} - x} \frac{\mathcal{F}(s, x)}{|\mathbf{1} - \Pi((\mathbf{1} - x)s)|^2} \frac{R((\mathbf{1} - x)s)\epsilon((\mathbf{1} - x)s)}{R(s)\epsilon(s)}$$

$$R(s) = -\frac{3}{\alpha} \operatorname{Im} \Pi_{\text{hadr}}(s)$$

Vacuum polarization according to  
CMD-2 data compilation:  
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Radiative correction factor  $1 + \delta$

$\sqrt{s}$ , MeV	$1 + \delta$	$\sqrt{s}$ , MeV	$1 + \delta$
1841.0	$1.0423 \pm 0.0208$	2542.6	$1.0739 \pm 0.0054$
1937.0	$1.0429 \pm 0.0156$	2644.8	$1.0796 \pm 0.0054$
2037.3	$1.0515 \pm 0.0126$	2744.6	$1.0809 \pm 0.0054$
2135.7	$1.0634 \pm 0.0106$	2849.7	$1.0823 \pm 0.0054$
2239.2	$1.0645 \pm 0.0096$	2948.9	$1.0774 \pm 0.0054$
2339.5	$1.0664 \pm 0.0075$	3048.1	$1.0584 \pm 0.0053$
2444.1	$1.0684 \pm 0.0064$		

# Selection criteria

Selection criteria for hadronic events which were used by AND.

Variable	Allowed range	
	3.12-3.72 GeV (2010-2011)	1.84 - 3.05 GeV
$N_{\text{track}}^{\text{IP}}$	$\geq 1$	$\geq 1$
$E_{\text{obs}}$	$> 1.6 \text{ GeV}$	$> 1.4 \text{ GeV} \left( > 1.3 \text{ GeV if } E_{\text{beam}} < 1.05 \text{ GeV} \right)$
$E_{\gamma}^{\max}/E_{\text{beam}}$	$< 0.8$	$< 0.8$
$E_{\text{obs}} - E_{\gamma}^{\max}$		$> 1.2 \text{ GeV} \left( > 1.1 \text{ GeV if } E_{\text{beam}} < 1.05 \text{ GeV} \right)$
$E_{\text{cal}}$	$> 0.75 \text{ GeV}$	$> 0.55 \text{ GeV}$
$H_2/H_0$	$< 0.85$	$< 0.9$
$ P_z^{\text{miss}}/E_{\text{obs}} $	$< 0.6$	$< 0.7$
$E_{\text{LKr}}/E_{\text{cal}}^{\text{tot}}$	$> 0.15$	$> 0.15$
$ Z_{\text{vertex}} $	$< 20.0 \text{ cm}$	$< 15.0 \text{ cm}$
	$N_{\text{particles}} \geq 4 \text{ or } \tilde{N}_{\text{track}}^{\text{IP}} \geq 2$	$N_{\text{particles}} \geq 3 \text{ or } \tilde{N}_{\text{track}}^{\text{IP}} \geq 2$

# The correlation matrix for systematic uncertainties of the R value obtained in the KEDR experiments

Point Correlation Matrix

1	1	0.139	0.143	0.193	0.192	0.212	0.212	0.216	0.207	0.211	0.216	0.201	0.222	0.096	0.046	0.096	0.105	0.110	0.098	0.089	0.114	0.071
2	1	0.309	0.418	0.408	0.445	0.437	0.466	0.446	0.457	0.467	0.434	0.480	0.200	0.097	0.201	0.225	0.229	0.212	0.189	0.244	0.151	
3	1	0.423	0.425	0.470	0.470	0.480	0.460	0.463	0.480	0.442	0.486	0.212	0.101	0.212	0.232	0.243	0.218	0.198	0.253	0.158		
4	1	0.575	0.635	0.635	0.649	0.622	0.610	0.649	0.598	0.637	0.287	0.137	0.286	0.314	0.329	0.295	0.268	0.342	0.213			
5	1	0.621	0.621	0.642	0.615	0.629	0.643	0.598	0.661	0.280	0.134	0.280	0.310	0.322	0.293	0.262	0.336	0.208				
6	1	0.677	0.709	0.679	0.695	0.710	0.661	0.730	0.306	0.148	0.305	0.342	0.351	0.323	0.287	0.371	0.229					
7	1	0.709	0.679	0.695	0.710	0.661	0.730	0.304	0.148	0.305	0.342	0.348	0.323	0.287	0.371	0.229						
8	1	0.695	0.710	0.725	0.675	0.745	0.320	0.153	0.320	0.351	0.368	0.330	0.299	0.382	0.238							
9	1	0.681	0.695	0.647	0.715	0.307	0.146	0.306	0.336	0.352	0.316	0.287	0.366	0.228								
10	1	0.710	0.654	0.701	0.314	0.150	0.313	0.344	0.360	0.323	0.293	0.374	0.233									
11	1	0.675	0.745	0.321	0.153	0.320	0.351	0.368	0.330	0.300	0.382	0.238										
12	1	0.687	0.298	0.142	0.298	0.327	0.342	0.307	0.279	0.356	0.222											
13	1	0.330	0.157	0.329	0.361	0.378	0.339	0.308	0.393	0.245												
14	1	0.288	0.396	0.405	0.394	0.356	0.317	0.403	0.333													
15	1	0.345	0.347	0.345	0.305	0.275	0.345	0.288														
16	1	0.486	0.475	0.427	0.380	0.483	0.400															
17	1	0.486	0.427	0.387	0.486	0.405																
18	1	0.427	0.380	0.483	0.400																	
19	1	0.340	0.427	0.356																		
20	1	0.384	0.318																			
21	1	0.403																				
22																						