

Proton Size Puzzle: Thick or Thin?

“Thick and Thin” famous Anton Chekhov novel

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Introduction

Experimental data

On 2-photon contributions

Conclusions

1) Phys.Part.Nucl.Lett. 14 (2017) 857

3) EPJA 54 (2018) 131

2) Phys.Lett. B776 (2018) 105;

4)PRA98 (2018) 042501

In 2010 the CREMA (Charge Radius Experiments with Muonic Atoms) Collaboration measured very precisely the Lamb shift of muonic hydrogen. It has opened the new era of the precise investigation of the spectrum of simple muonic atoms.

In the new experiments by this Collaboration with muonic deuterium and ions of muonic helium a charge radii of light nuclei were obtained with very high precision.

For muonic hydrogen and muonic deuterium it was shown that obtained values of the charged radii are significantly different from those which were extracted from spectra of electronic atoms and in the scattering of the electrons with nucleon and were recommended for using by CODATA, so-called, “PROTON CHARGE RADIUS PUZZLE”.

Several experimental groups plan to measure the hyperfine structure of various muonic and electronic atoms with more high precision.

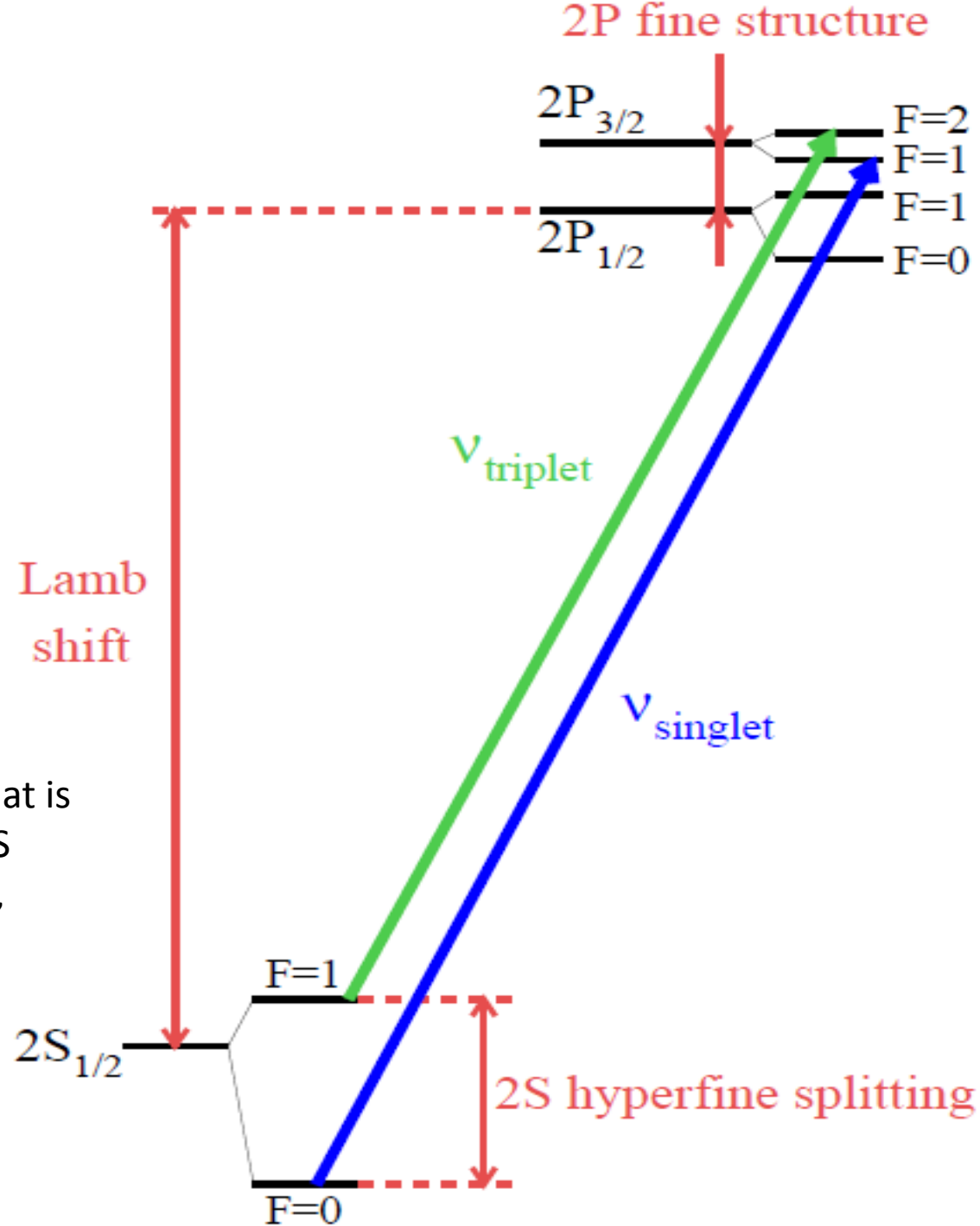
One can consider experiments with **muonic atoms** as a **smoking gun** for:

Precise measurements of the **proton charge radius and other parameters**

Test of the **Standard Model** with greater accuracy

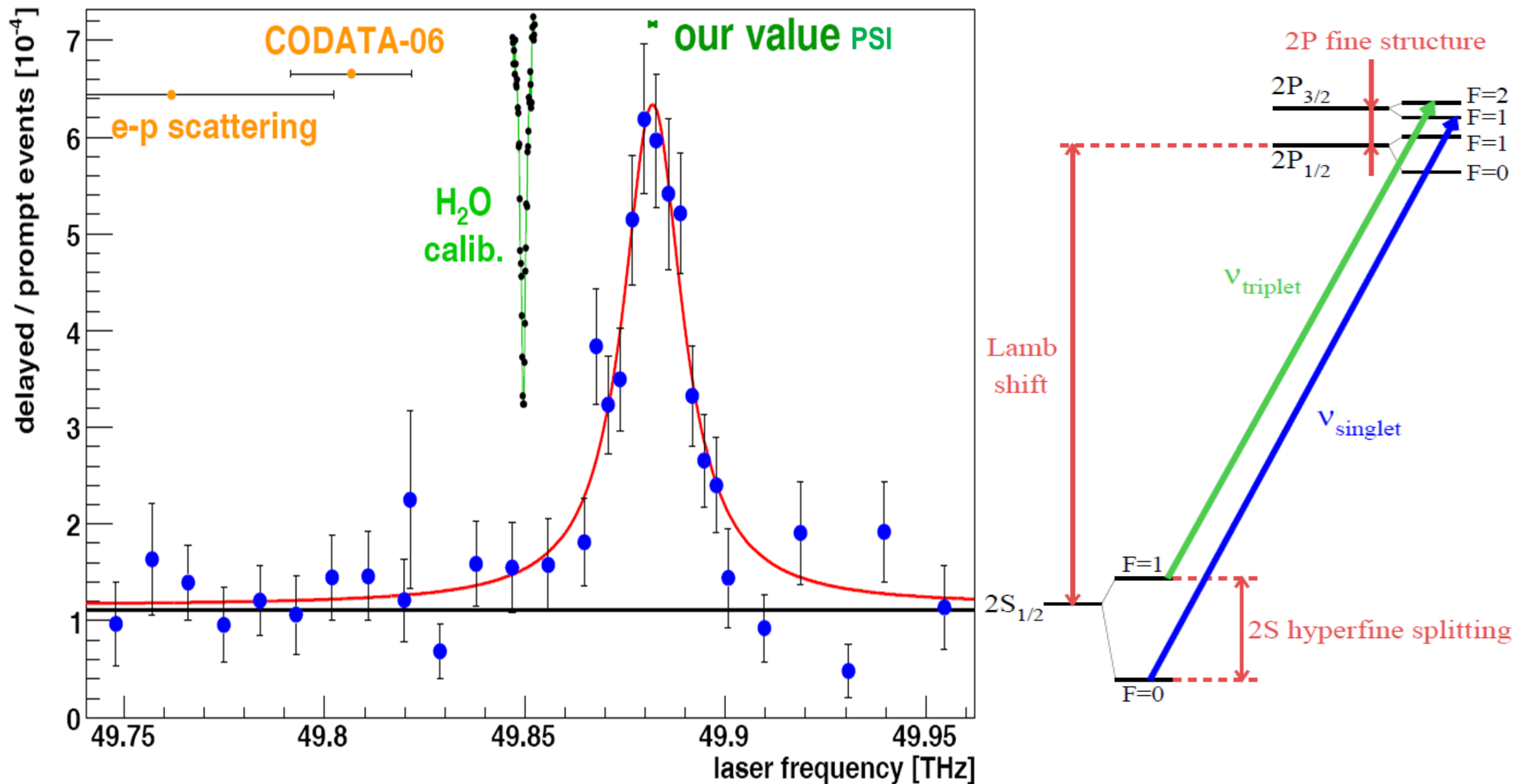
and, possibly, to reveal the source of previously unaccounted interactions between the particles forming the **bound state in QED**.

The HFS requires the spin-spin coupling that is the interaction between the nuclear spin S and the lepton total angular momentum j , where $F = j + S$ is the atom total angular momentum

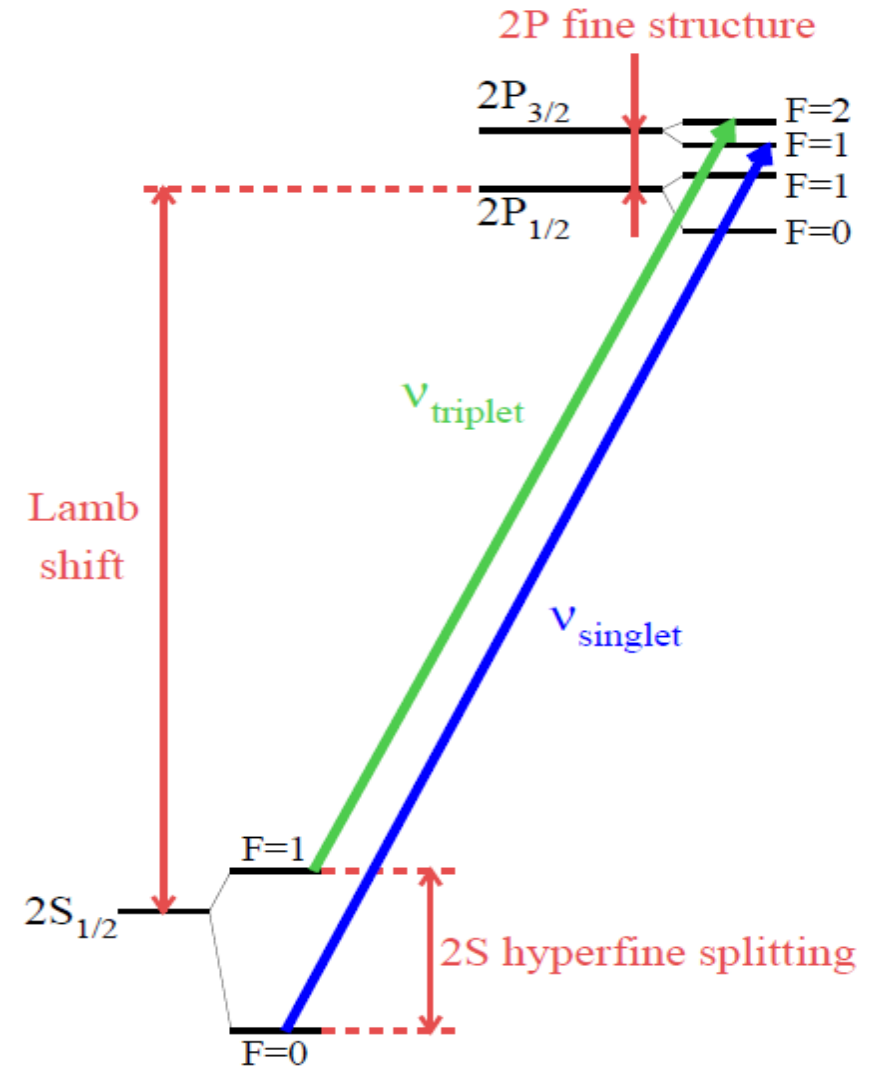
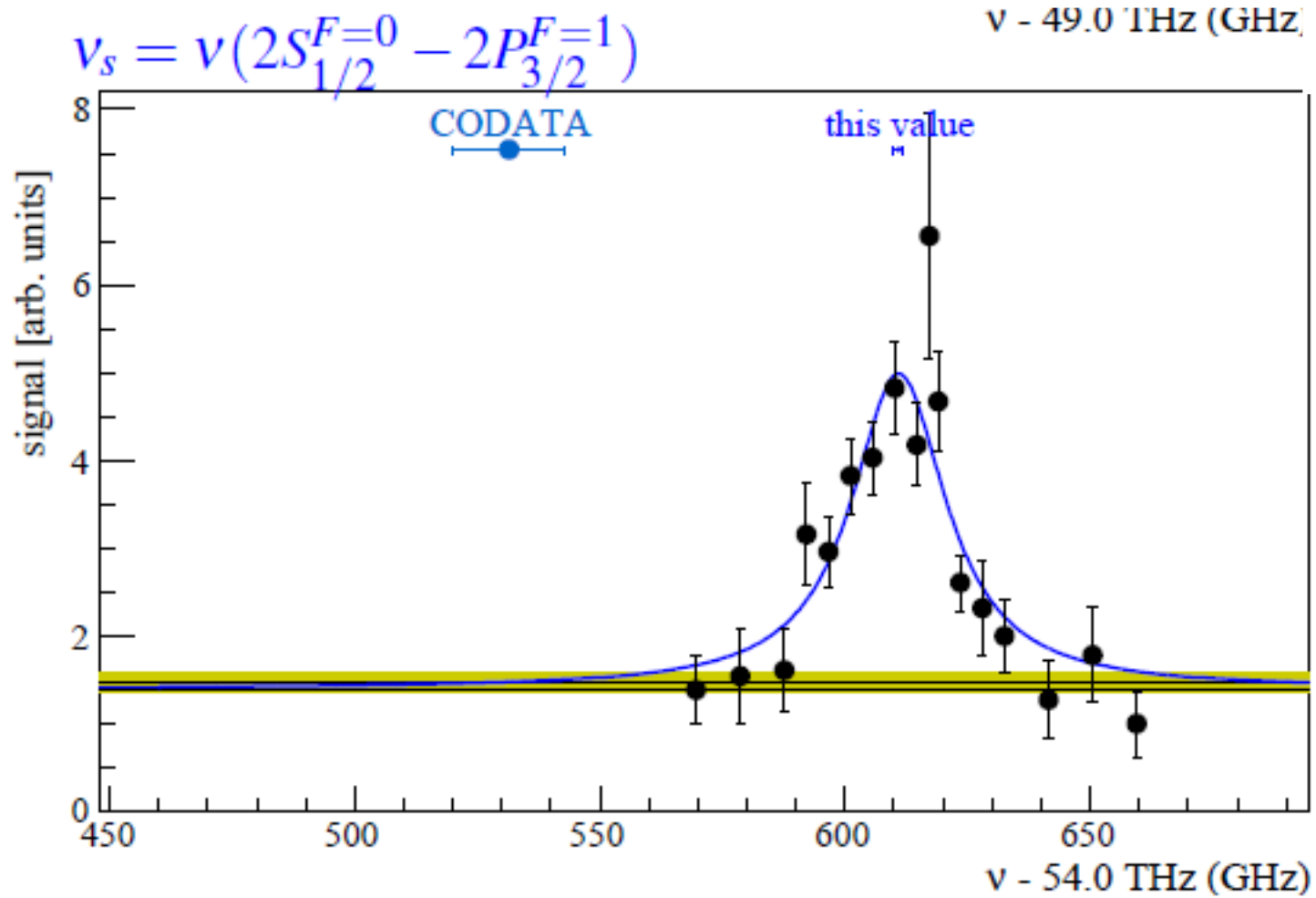


$$\nu_t = \nu(2S_{1/2}^{F=1} - 2P_{3/2}^{F=2})$$

CREMA-10-13



CREMA-13



Two transitions measured

$$\nu_T = 49881.35(65) \text{ GHz}$$

$$\nu_S = 54611.16(1.05) \text{ GHz}$$

From these two transition measurements, one can independently deduce both the Lamb shift $\Delta E_L = \Delta E(2P_{1/2} - 2S_{1/2})$ and the 2S-HFS splitting (ΔE_{HFS}) by the linear combinations

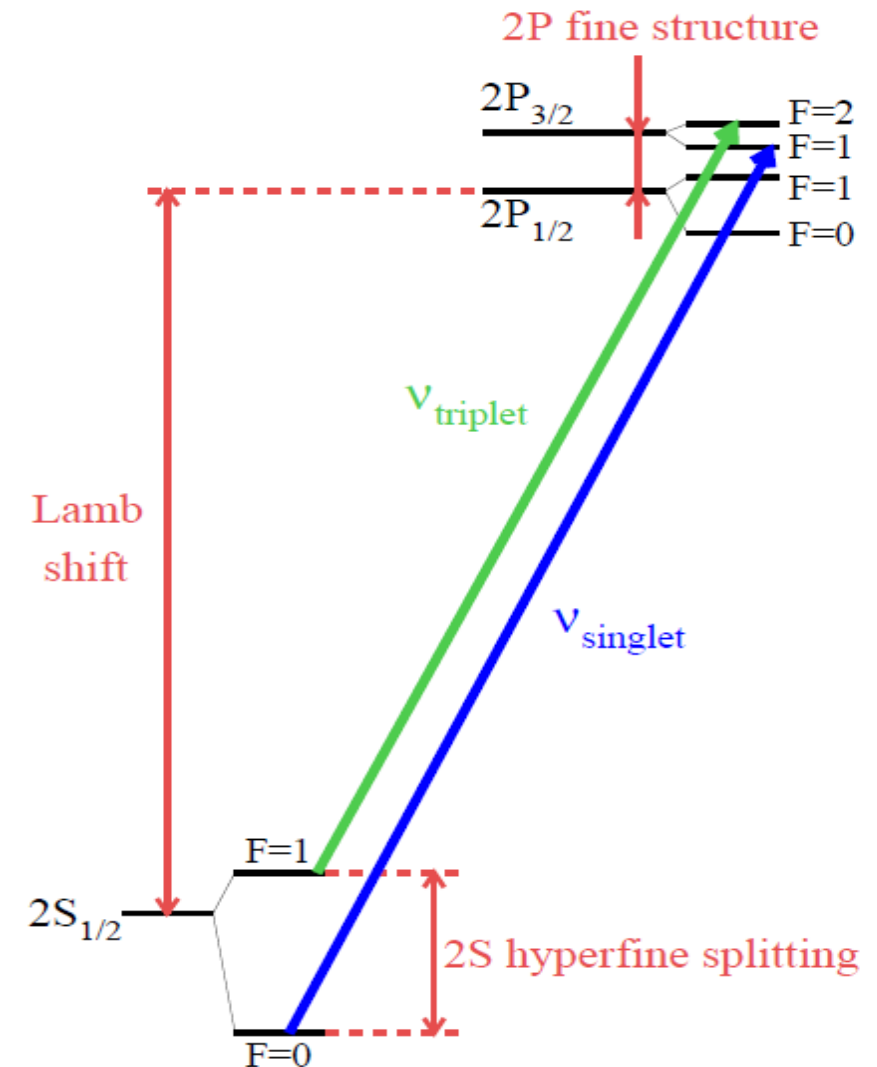
$$\frac{1}{4}h\nu_s + \frac{3}{4}h\nu_t = \Delta E_L + 8.8123(2)\text{meV}$$

$$h\nu_s - h\nu_t = \Delta E_{\text{HFS}} - 3.2480(2)\text{meV}$$

Then one get

$$\Delta E_{\text{LS}}^{\text{exp}} = 202.3706(23) \text{ meV}$$

$$\Delta E_{\text{HFS}}^{\text{exp}} = 22.8089(51) \text{ meV}$$



Proton Size Puzzle

$$r_p = 0.8751(61) \text{ fm (CODATA-2014)}$$

From spectrum of electronic atoms and
the scattering of the electrons with nuclei

Lamb shift in muonic Hydrogen (μp) (a proton orbited by a negative muon)

muon mass $m_\mu \approx 200 \times m_e$

Bohr radius $r_\mu \approx 1/200 \times r_e$

μp has much smaller Bohr radius compared to electronic hydrogen and
so is **much more sensitive** to the finite size of the proton

μ inside the proton: $200^3 \approx 10^7$

$$\Delta E_{\text{LS}}^{\text{exp}} = 202.3706(23) \text{ meV}$$

CREMA experiment **Nature 2010; Science 2013**
(0.05% precision)

$$\Delta E_{\text{LS}} = 206.0668(25) - 5.2275(10) \langle r_p^2 \rangle + \mathcal{O}(\langle r_p^2 \rangle^{3/2}) [\text{meV}]$$

Theory summary: Antognini et al. AnnPhys 2013
(2% effect)

$$r_p = 0.84087(39) \text{ fm (CREMA coll. Antognini et al., 2013)}$$

5.6 σ deviation!

$$r_E^2 = \int d^3r r^2 \rho_E(r)$$

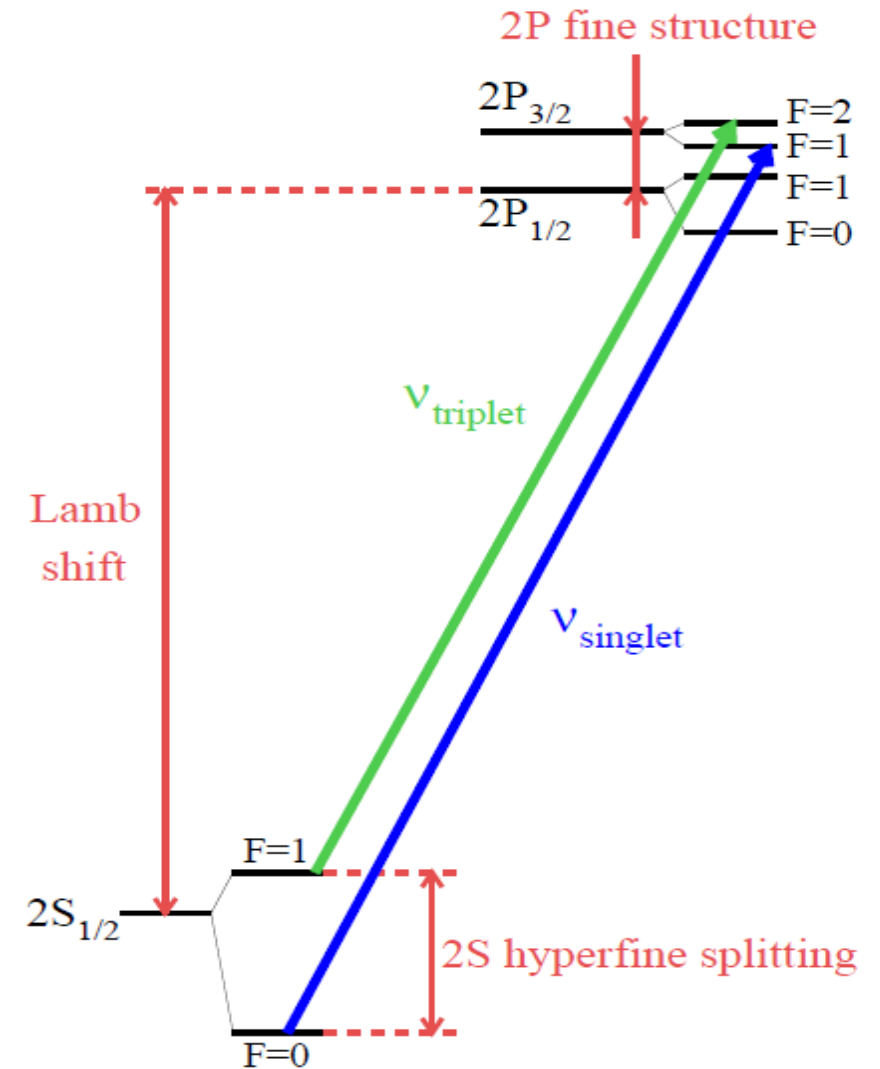
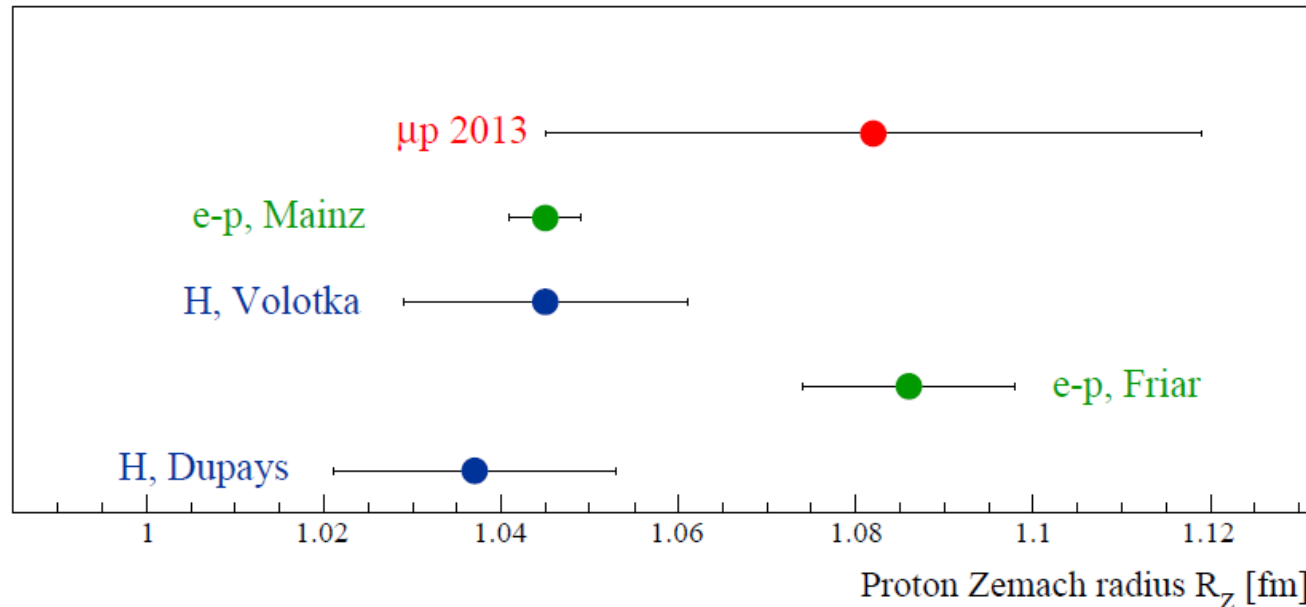
Zemach radius

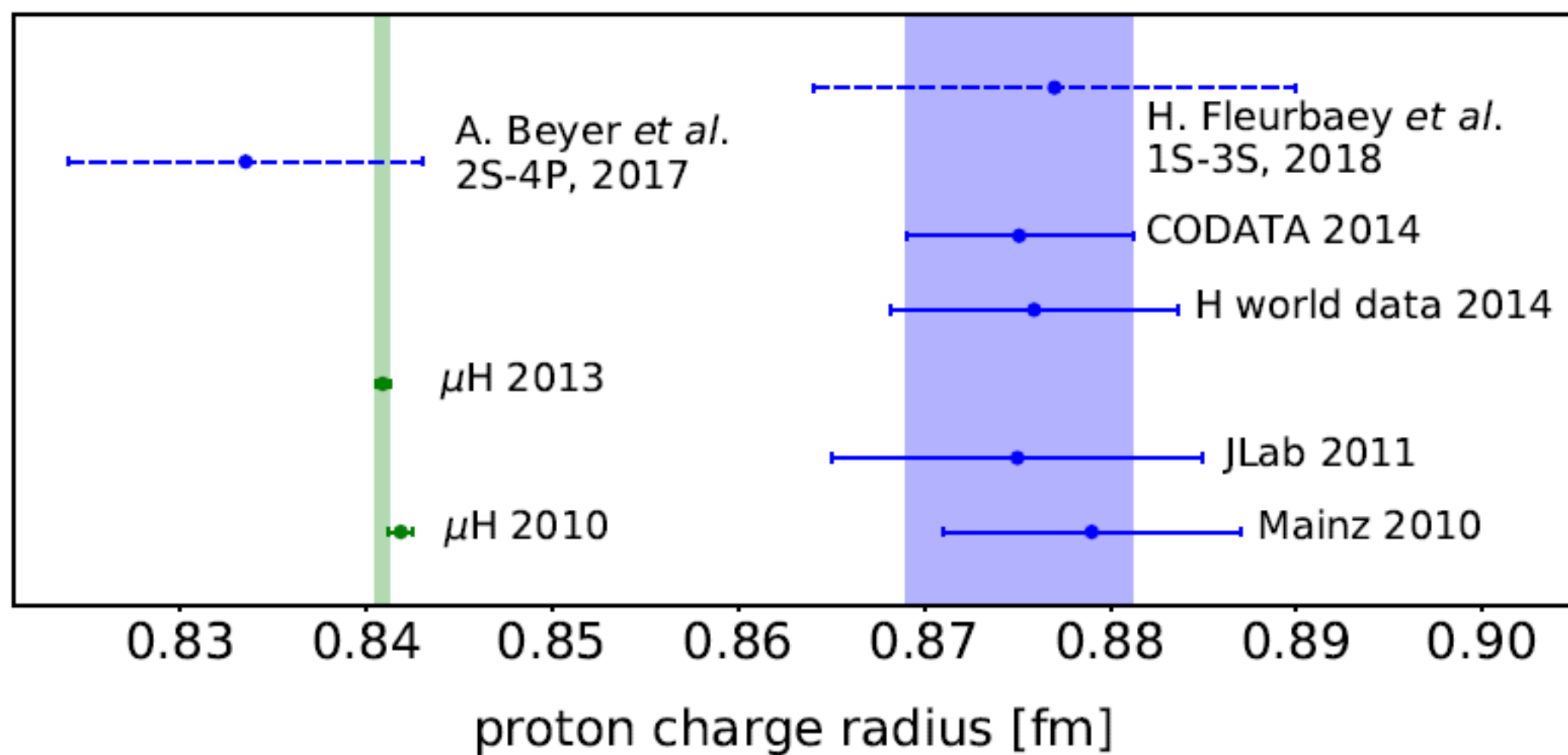
$$\Delta E_{2S-HFS}^{\text{exp}} = 22.8089(51) \text{ meV}$$

$$\Delta E_{2S-HFS}^{\text{th}} = 22.9843(30) - 0.1621(10)r_Z \text{ meV}$$

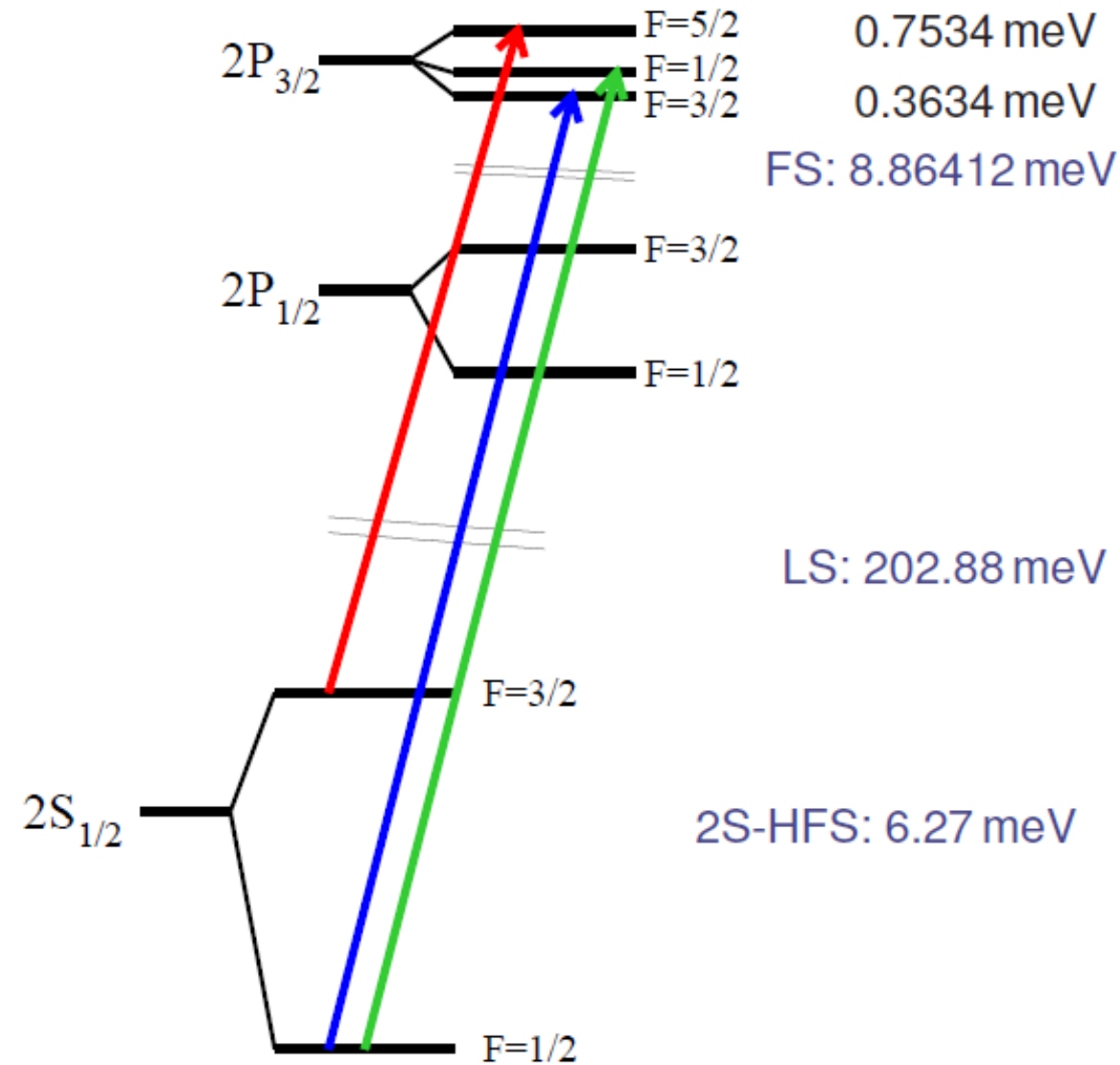
$$r_Z = \int d^3r \int d^3r' r \rho_E(r) \rho_M(r-r')$$

$$r_Z = 1.082(37) \text{ fm (CREMA coll. Antognini et al., 2013)}$$





Muonic deuterium CREMA-16



Deuteron charge radius

$$\Delta E_{\text{LS}}^{\text{exp}} = 202.8785(34) \text{ meV}$$

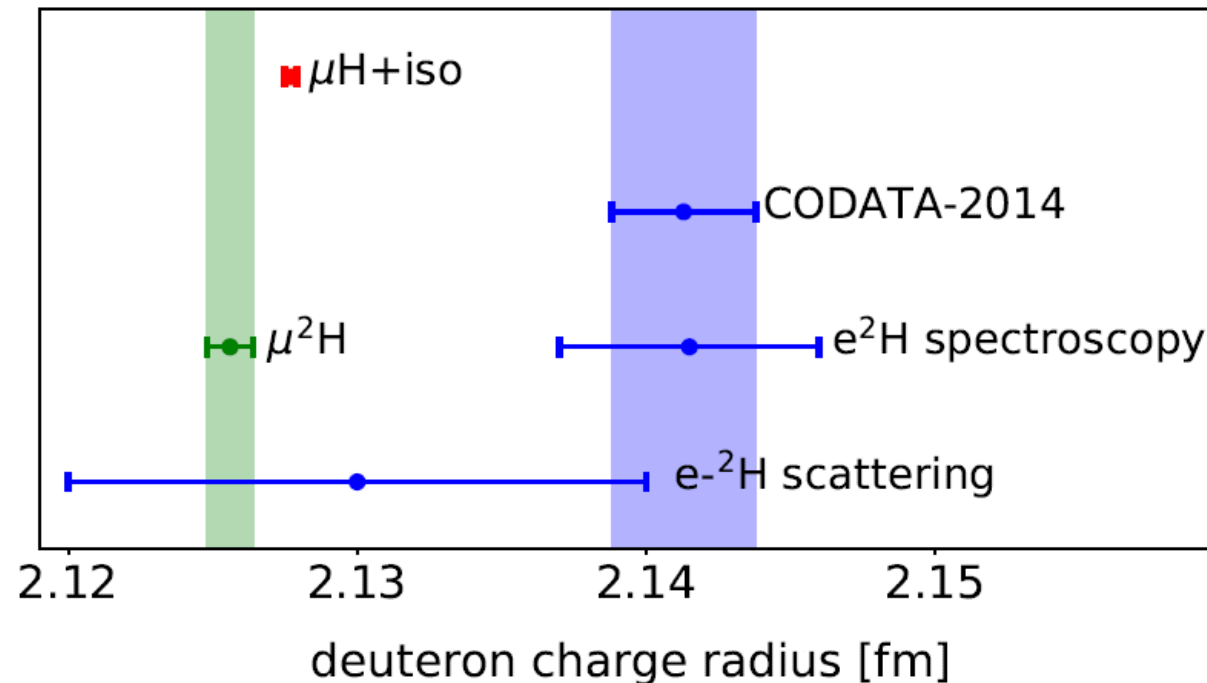
$$\begin{aligned} \Delta E_{\text{LS}}^{\text{th}} &= 228.7766(10) \text{ meV (QED)} \\ &+ 1.7096(200) \text{ meV (TPE)} \\ &- 6.1103(3) r_d^2 \text{ meV/fm}^2 \end{aligned}$$

$$\mu\text{D} \Rightarrow r_d = 2.12562(78) \text{ meV (CREMA 2016)}$$

H/D isotope shift: $r_d^2 - r_p^2 = 3.82007(65) \text{ fm}^2$

$$r_p \text{ from } \mu\text{H} \text{ gives } r_d = 2.12771(22) \text{ fm} \leftarrow 5.4\sigma \text{ from } r_p$$

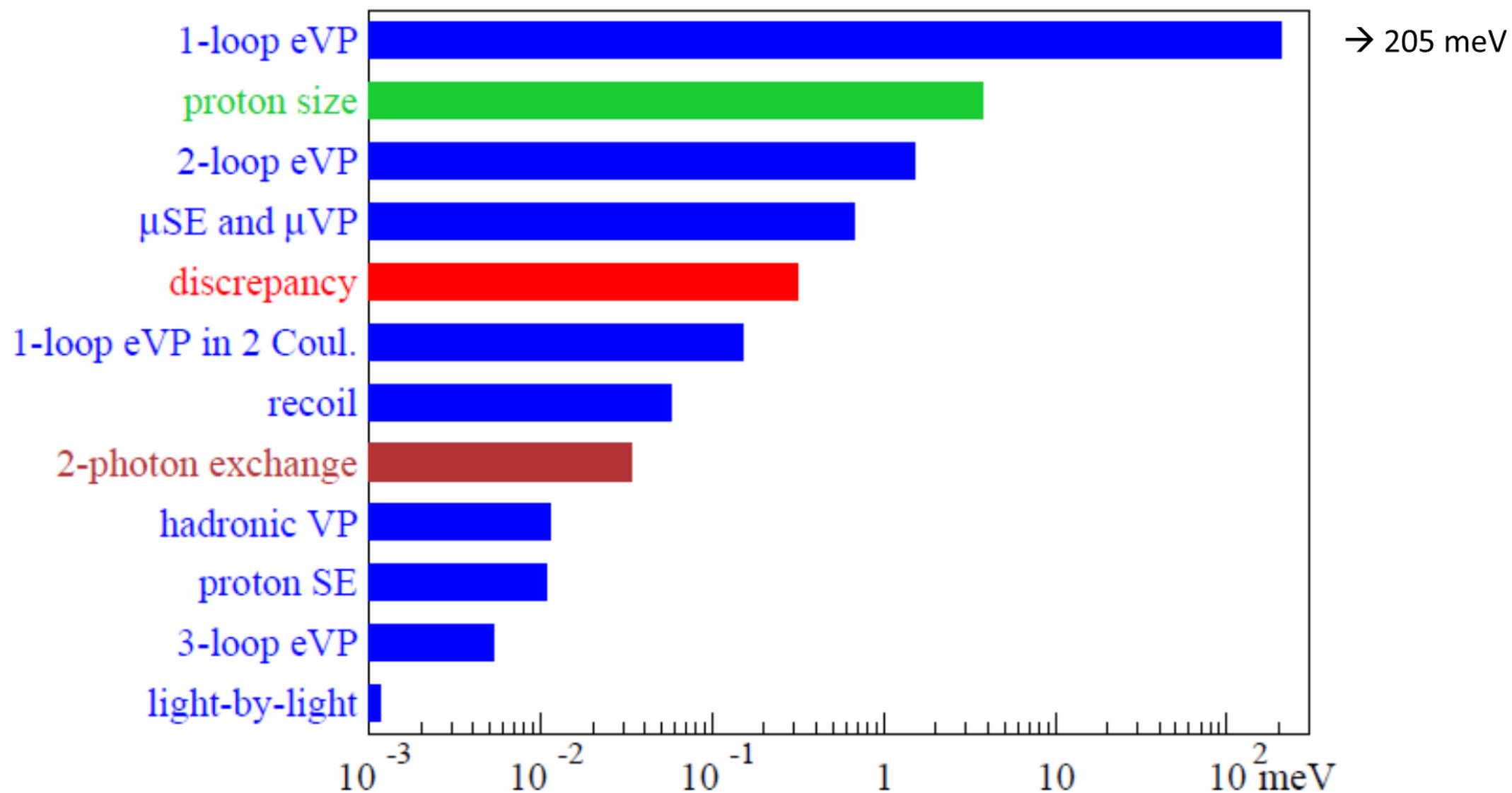
$$\text{CODATA 2014 } r_d = 2.14130(250) \text{ fm}$$



r_d is 7.5σ smaller than CODATA-2010
(99% correlated with r_p !)

3.5σ smaller than r_d (D spectroscopy)

Hierarchy of contributions to Lamb shift



The leading contribution to the hyperfine splitting (HFS) in muonic hydrogen is coming from one-photon exchange
 [AP Martynenko, RN Faustov 2004]

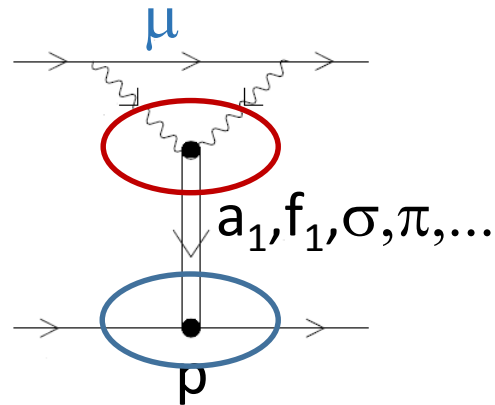
$$\Delta V_B^{hfs} = \frac{8\pi\alpha\mu_p}{3m_\mu m_p} (\mathbf{S}_p \mathbf{S}_\mu) \delta(\mathbf{r}) - \frac{\alpha\mu_p(1+a_\mu)}{m_\mu m_p r^3} [(\mathbf{S}_p \mathbf{S}_\mu) - 3(\mathbf{S}_p \mathbf{n})(\mathbf{S}_p \mathbf{n})] +$$

$$\frac{\alpha\mu_p}{m_\mu m_p r^3} \left[1 + \frac{m_\mu}{m_p} - \frac{m_\mu}{2m_p \mu_p} \right] (\mathbf{L} \mathbf{S}_p)$$

We calculate further the contribution to HFS coming from light pseudoscalar and axial-vector meson exchanges

New theoretical contributions to the Lamb shift and HF structure of muonic hydrogen μP :

Meson exchanges through the anomalous two-photon vertex



Pseudosclar mesons

The effective vertex of the interaction of the **PS mesons** and virtual photons can be expressed in terms of the transition form factors

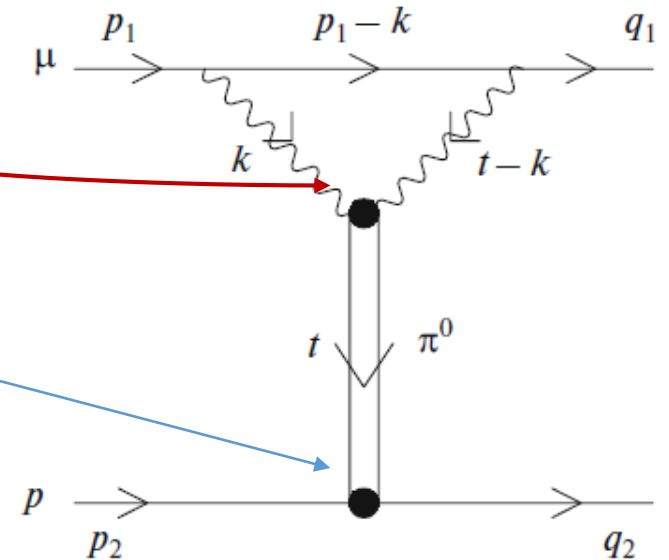
$$A \left(\gamma_{(q_1, \epsilon_1)}^* \gamma_{(q_2, \epsilon_2)}^* \rightarrow \pi_{(p)} \right) = ie^2 \varepsilon_{\mu\nu\rho\sigma} \epsilon_1^\mu \epsilon_2^\nu q_1^\rho q_2^\sigma \frac{1}{4\pi^2 F_\pi} F_{\pi\gamma^*\gamma^*} (p^2; q_1^2, q_2^2)$$

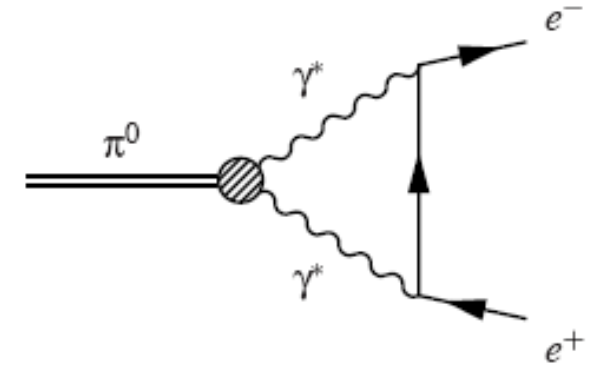
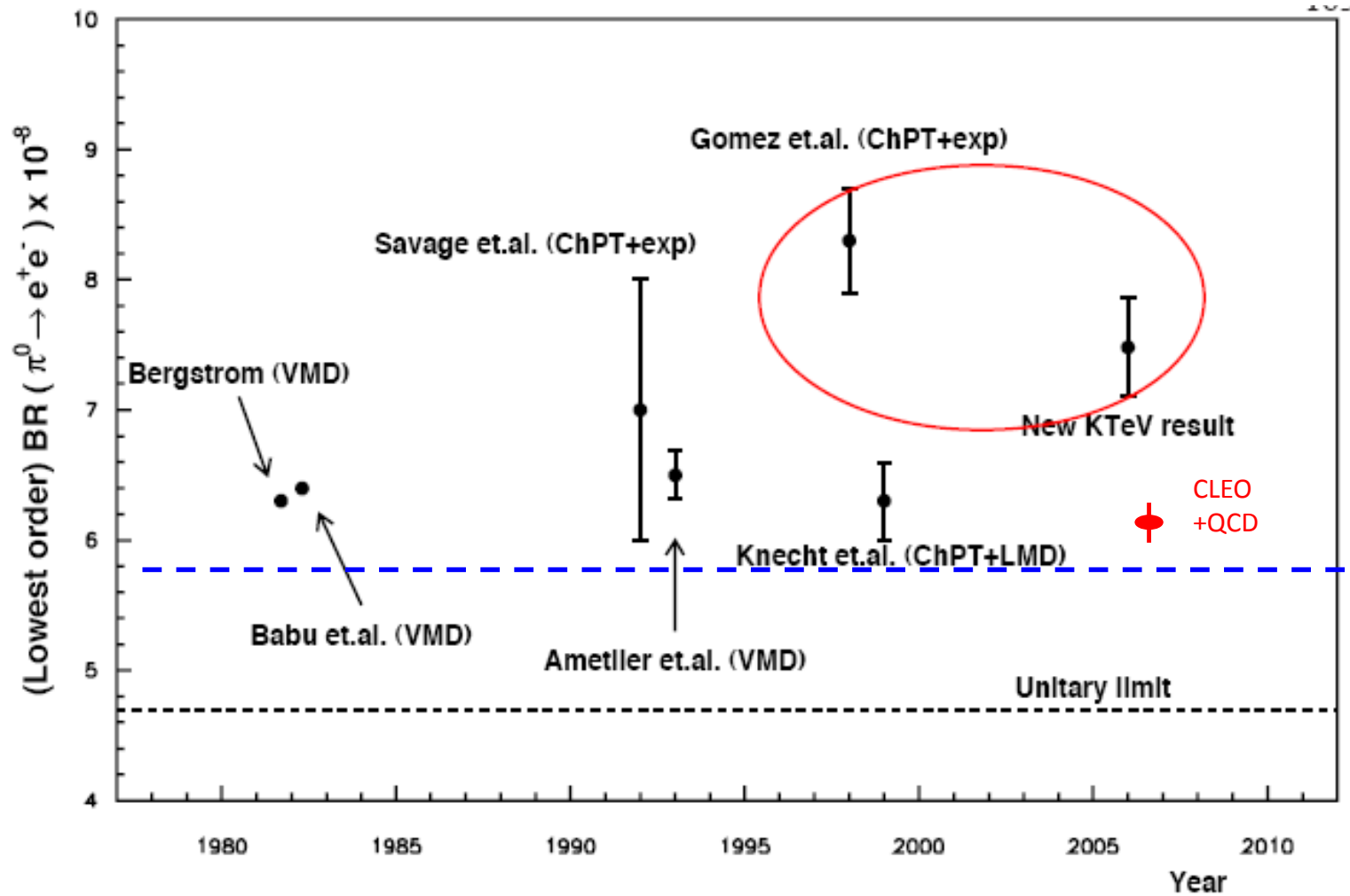
As a result, the hyperfine part of the potential of the one-pion interaction of a muon and a proton in the S-state takes the form

$$\Delta V^{hfs}(\mathbf{p}, \mathbf{q}) = \frac{\alpha^2}{6\pi^2} \frac{g_p}{m_p F_\pi} \frac{(\mathbf{p} - \mathbf{q})^2}{(\mathbf{p} - \mathbf{q})^2 + m_\pi^2} \mathcal{A}(t^2),$$

where

$$\mathcal{A}(t^2) = \frac{2i}{\pi^2 t^2} \int d^4 k \frac{t^2 k^2 - (tk)^2}{k^2 (k-t)^2 (k^2 - 2kp_1)} F_{\pi\gamma^*\gamma^*}(k^2, (k-t)^2).$$





} **3σ diff**

$$A_{M \rightarrow ll}(0) = -3 \ln \left(\frac{m_l^2}{\Lambda_{M \rightarrow \gamma\gamma}^2} \right) - \frac{3}{2} + O \left(\frac{m_l^2}{\Lambda_{M \rightarrow \gamma\gamma}^2} \right)$$

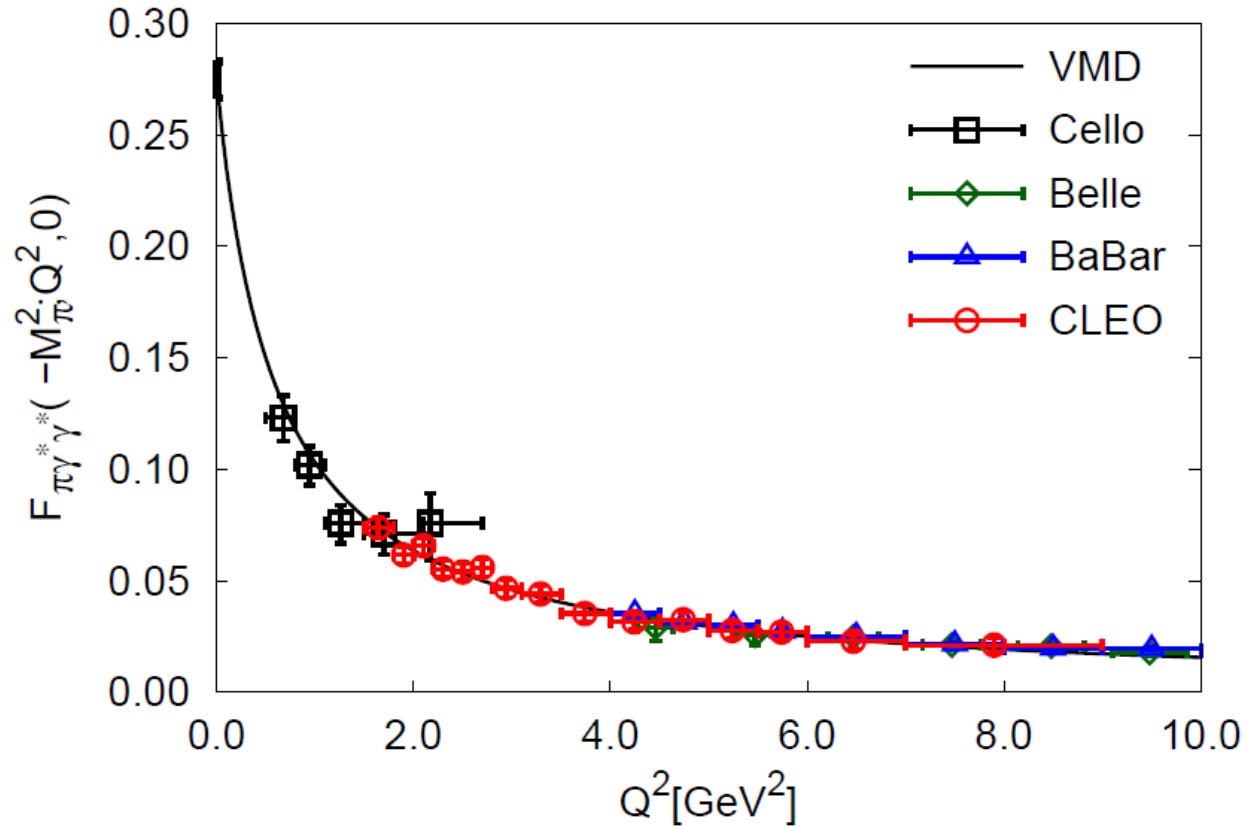
$$B_{\pi \rightarrow e^+e^-}^{\text{theory}} = (6.2 \pm 0.1) \cdot 10^{-8}$$

Which is **3.3σ below data!!**

$$B_{\pi \rightarrow e^+e^-}^{\text{KTeV}} = (7.48 \pm 0.38) \cdot 10^{-8}$$

A. E. Dorokhov and M. A. Ivanov, Phys. Rev. D **75** (2007) 114007

Upper block: two-gamma transition FF into PS meson



$$F_{\pi\gamma^*}(Q^2, 0) = \frac{\Lambda_\pi^2}{\Lambda_\pi^2 + k^2}, \quad \Lambda_\pi^{CLEO} = 776 \pm 22 \text{ MeV},$$

$$F_{\pi\gamma^*\gamma^*}(Q^2, Q^2) = \frac{\overline{\Lambda}_\pi^2}{\overline{\Lambda}_\pi^2 + k^2}, \quad \overline{\Lambda}_\pi^{CLEO+QCD} = 499 \pm 50 \text{ MeV}$$

Calculating the matrix elements with wave functions of 1S , 2S and 2P_{1/2} states, we obtain the corresponding contributions to the HFS spectrum (**we use data from CLEO on TFF**)

$$\Delta E^{hfs}(1S) = \frac{\mu^3 \alpha^5 g_A}{6F_\pi^2 \pi^3} \left\{ \mathcal{A}(0) \frac{4W(1 + \frac{W}{m_\pi})}{m_\pi(1 + \frac{2W}{m_\pi})^2} - \frac{1}{\pi} \int_0^\infty \frac{ds}{s} \text{Im} \mathcal{A}(s) \left[1 + \frac{1}{4W^2(s - m_\pi^2)} \left(\frac{m_\pi^4}{(1 + \frac{m_\pi}{2W})^2} - \frac{s^2}{(1 + \frac{\sqrt{s}}{2W})^2} \right) \right] \right\} = -0.0017 \text{ meV},$$

$$\Delta E^{hfs}(2S) = \frac{\mu^3 \alpha^5 g_A}{48F_\pi^2 \pi^3} \left\{ \mathcal{A}(0) \frac{W(8 + 11\frac{W}{m_\pi} + 8\frac{W^2}{m_\pi^2} + 2\frac{W^3}{m_\pi^3})}{2m_\pi(1 + \frac{W}{m_\pi})^4} - \frac{1}{\pi} \int_0^\infty \frac{ds}{s} \text{Im} \mathcal{A}(s) \left[1 + \frac{1}{(s - m_\pi^2)} \left(\frac{m_\pi^2(2 + \frac{W^2}{m_\pi^2})}{2(1 + \frac{W}{m_\pi})^4} - \frac{s(2 + \frac{W^2}{s})}{2(1 + \frac{W}{\sqrt{s}})^4} \right) \right] \right\} = -0.0002 \text{ meV},$$

$$\Delta E_{2P_{1/2}}^{hfs} = \frac{\alpha^7 \mu^5 g_A}{288\pi^3 F_\pi^2 m_\pi^2} \mathcal{A}(0) \frac{\left(9 + 8\frac{W}{m_\pi} + 2\frac{W^2}{m_\pi^2}\right)}{\left(1 + \frac{W}{m_\pi}\right)^4} = 0.0004 \text{ } \mu\text{eV}.$$

Axial-Vector Mesons

The transition from initial state of two virtual photons with four-momenta k_1, k_2 to an **axial vector meson** A ($JPC = 1^{++}$) with the mass M_A for the small values of the relative momenta of particles in the initial and final states and small value of transfer momentum t between muon and proton is

$$T^{\mu\nu\alpha} = 8\pi i \alpha \varepsilon_{\mu\nu\alpha\tau} (k_1^\tau k_2^2 - k_2^\tau k_1^2) F_{AV\gamma^*\gamma^*}^2(M_A^2; k_1^2, k_2^2)$$

The final result for the HFS potential is

$$\Delta V_{AV}^{hfs}(\mathbf{p} - \mathbf{q}) = -\frac{32\alpha^2 g_{AVPP}}{3\pi^2(t^2 + M_A^2)} \int id^4k \frac{(2k^2 + k_0^2)}{k^2(k^2 - 2m_\mu k_0)} F_{AV\gamma^*\gamma^*}(0, k^2, k^2)$$

By using L3 Collaboration data (production of $f_1(1285)$ and $f_1(1420)$), we can parameterize the transition form factor for the case of two photons with equal virtualities

$$F_{AV\gamma^*\gamma^*}(M_A^2; k^2, k^2) = F_{AV\gamma^*\gamma^*}(M_A^2; 0, 0) F_{AV}^2(k^2) \quad F_{AV}(k^2) = \frac{\Lambda_A^2}{\Lambda_A^2 + k^2}$$

with the normalization fixed from the decay width of axial-vector meson measured by **L3 Collaboration**. (Expect improvements From BESIII and BELLEII)

$$\Lambda_{f_1(1285)}^{L3} = 1040 \pm 78 \text{ MeV}, \quad \Lambda_{f_1(1420)}^{L3} = 926 \pm 78 \text{ MeV}$$

Meson-Nucleon Couplings

Within the NJL model from Triangle diagram one has

$$F_{f_1(1285)\gamma^*\gamma^*}(M_{f_1(1285)}^2; 0, 0) = \frac{5g_{f_1(1285)qq}}{72\pi^2 M_q^2}$$

Another couplings are related to each others by using chiral symmetry and SU(6)-model for wave function of the proton

$$g_{f_1(1260)qq} = g_{f_1(1285)qq},$$

$$g_{f_1(1285)pp} = g_{f_1(1285)qq}, \quad g_{f_1(1260)pp} = \frac{5}{3} g_{f_1(1285)qq}$$

With $M_q=300$ MeV one gets

$$g_{f_1(1285)pp} = g_{f_1(1285)qq} = g_{f_1(1260)qq} = 3.40 \pm 1.19,$$

$$g_{f_1(1260)pp} = 5.67 \pm 1.98, \quad g_{f_1(1420)pp} = 1.51 \pm 0.19$$

Our result for HFS interaction can be rewritten ($a=2m_\mu/\Lambda_A$)

$$\Delta V_{AV}^{hfs}(\mathbf{p} - \mathbf{q}) = -\frac{32\alpha^2 g_{AVpp} F_{AV\gamma^*\gamma^*}^{(0)}(0, 0, 0)}{3\pi^2(\mathbf{t}^2 + M_A^2)} I\left(\frac{m_\mu}{\Lambda_A}\right)$$

$$I\left(\frac{m_\mu}{\Lambda_A}\right) = -\frac{\pi^2 \Lambda_A^2}{4(1 - a_\mu^2)^{5/2}} \left[3\sqrt{1 - a_\mu^2} - a_\mu^2(5 - 2a_\mu^2) \ln \frac{1 + \sqrt{1 - a_\mu^2}}{a_\mu} \right]$$

Making the Fourier transform and averaging the obtained expression with the wave functions for 1S and 2S states, we obtain the following contribution to hyperfine splitting

$$\Delta E_{AV}^{hfs}(1S) = \frac{32\alpha^5 \mu^3 \Lambda^2 g_{AVNN} F_{AV\gamma^*\gamma^*}^{(0)}(0, 0, 0)}{3M_A^2 \pi^3 \left(1 + \frac{2W}{M_A}\right)^2} I\left(\frac{m_\mu}{\Lambda}\right),$$

$$\Delta E_{AV}^{hfs}(2S) = \frac{2\alpha^5 \mu^3 \Lambda^2 g_{AVNN} F_{AV\gamma^*\gamma^*}^{(0)}(0, 0, 0) \left(2 + \frac{W^2}{M_A^2}\right)}{3M_A^2 \pi^3 \left(1 + \frac{W}{M_A}\right)^4} I\left(\frac{m_\mu}{\Lambda}\right)$$

where $W = \mu \alpha$ and μ is the reduced mass.

Axial-Vector and Pseudoscalar mesons contribution to HFS of muonic hydrogen

mesons	$I^G(J^{PC})$	Λ_A in MeV	$F_{AV\gamma^*\gamma^*}^{(0)}(0,0)$ in GeV^{-2}	$\Delta E^{hfs}(1S)$ in meV	$\Delta E^{hfs}(2S)$ in meV
$f_1(1285)$	$0^+(1^{++})$	1040	0.266	-0.0093 ± 0.0033	-0.0012 ± 0.0004
$a_1(1260)$	$1^-(1^{++})$	1040	0.591	-0.0437 ± 0.0175	-0.0055 ± 0.0022
$f_1(1420)$	$0^+(1^{++})$	926	0.193	-0.0013 ± 0.0008	-0.0002 ± 0.0001
π^0	$1^-(0^{-+})$	776		-0.0017 ± 0.0001	-0.0002 ± 0.00002
Sum				-0.0560 ± 0.0178	-0.0071 ± 0.0024

$$\Delta E_{2S-HFS}^{\text{exp}} = 22.8089(51) \text{ meV}$$

$$\Delta E_{2S-HFS}^{\text{th}} = 22.9843(30) - 0.1621(10)r_Z \text{ meV}$$

$$r_Z = 1.082(37) \text{ fm (CREMA coll. Antognini et al., 2013)}$$

$$0.001 \text{ meV} \rightarrow 0.006 \text{ fm}$$

$$r_Z = 1.040(37) \text{ fm}$$

Scalar mesons contribution to the Lamb shift of muonic hydrogen

Scalar meson	$I^G(J^{PC})$	$\Gamma_{\gamma\gamma}^{exp}$ in keV	Λ_S in MeV	A_S in GeV^{-1}	$\Delta E^{Ls}(2P - 2S)$ in μeV
$\sigma(550)$	$0^+(0^{++})$	4.5	2000	-0.299	17.05
$f_0(980)$	$0^+(0^{++})$	0.33	2000	-0.034	0.72
$a_0(980)$	$1^-(0^{++})$	0.30	2000	-0.032	0.53
$f_0(1370)$	$0^+(0^{++})$	4.5	2000	-0.075	0.81

$$\Delta E_{LS}^{exp} = 202.3706(23) \text{ meV}$$

$$\Delta E_{LS} = 206.0668(25) - 5.2275(10) \langle r_p^2 \rangle + O\left(\langle r_p^2 \rangle^{3/2}\right) [\text{meV}]$$

$$r_p = 0.84087(39) \text{ fm (CREMA coll. Antognini et al., 2013)}$$

Conclusions

From **CREMA** experiment side:

“Proton radius puzzle” is in fact “**Z=1 radius puzzle**”

$$r_p = 0.84087(39) \text{ fm}$$

$$r_Z = 1.082(37) \text{ fm}$$

Muonic **helium-3 and -4** ions show (preliminary) no big discrepancy

New projects are ongoing, one of them FAMU (Fisica Atomica MUonica) with accuracy 2 ppm

In our theoretical work:

A new **large** contribution to the HFS of muonic hydrogen is found, that
Induced by **pseudoscalar, axial-vector and scalar** couplings to two photon state.

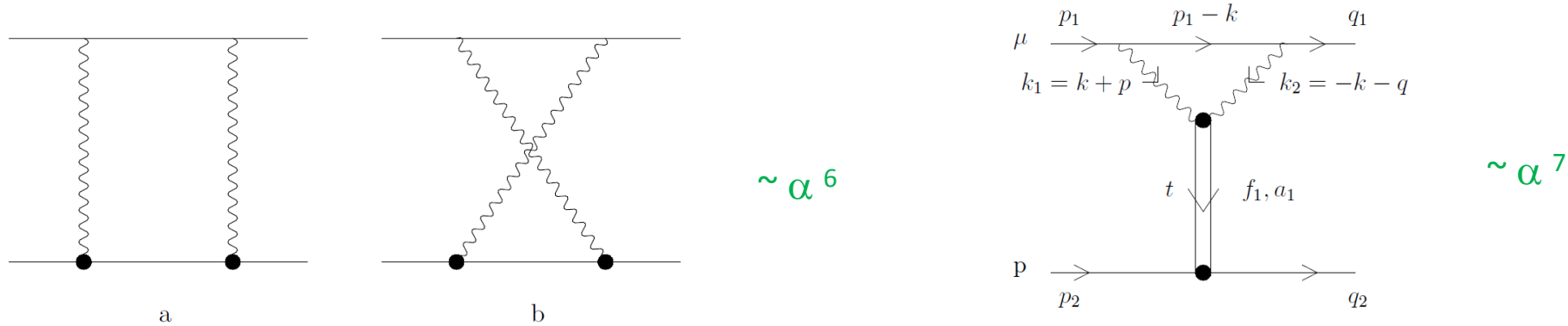
In particular these results provide diminishing of the Zemach radius
(compare with $r_Z = 1.040(37) \text{ fm}$ $r_Z = 1.082(37) \text{ fm}$)

It should be taking into account for the interpretation of the new data on HFS in this atom.

There are still a number of uncertainties in phenomenological input used in our calculations and some other new effects unaccounted by us. (Work is in progress)

This mechanism is very important to obtain the total value of the HFS in μP with high precision and needs to be account in current experiments (CREMA (PSI), FAMU (GB), ...)

Corrections of 2-photon interaction to the fine and hyperfine structure of P-energy levels of mu-hydrogen



State	Correction of 2γ exchange amplitudes, μeV	Correction of axial meson exchange amplitudes, μeV
$2^1P_{1/2}$	-0.2027	0.0005
$2^3P_{1/2}$	0.0761	-0.0002
$2^3P_{3/2}$	0	-0.00005
$2^5P_{3/2}$	0	0.00003

The corrections from two-photon exchange amplitudes are more significant for a precise comparison with the experimental data because their numerical values are of the order of 0.0001 meV.

Recall that to explain the proton radius puzzle , a contribution of about 0.3 meV is needed.