

Light mesons from JPAC+COMPASS analyses

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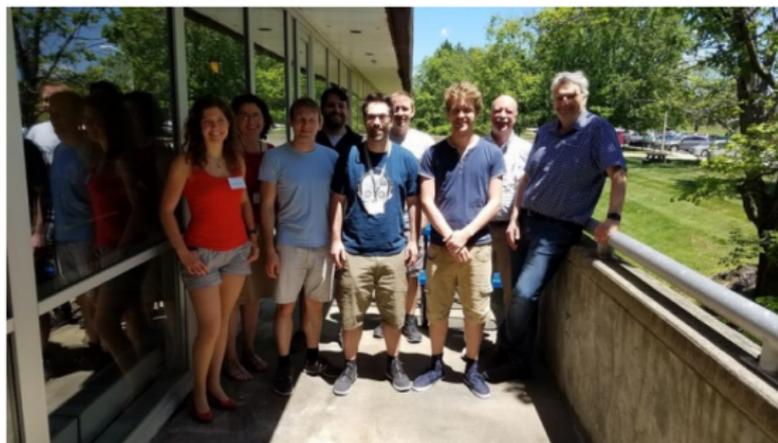
Introduction

Joint Physics Analysis Center

[collage by Vincent Mathieu]

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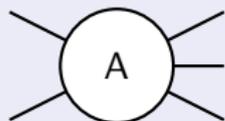


Miguel Albaladejo
JLab

JPAC effort

Collaboration of theoreticians and experimentalists

Amplitude analysis

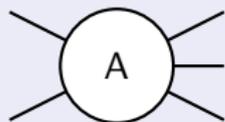


- Writing amplitudes using general QFT constraints
- Analysis of experimental data
- Analytic continuation, pole search

JPAC effort

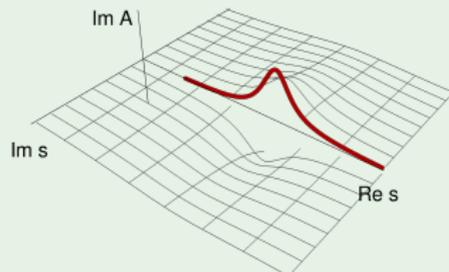
Collaboration of theoreticians and experimentalists

Amplitude analysis



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General properties of the scattering amplitude



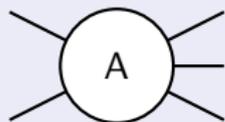
Analyticity + Unitary + Crossing symmetry

- Scattering amplitude is an analytic function in $s = E^2$ complex plane,
- The Real axis \rightarrow physical world,
- Resonances = poles of the unphysical sheet.

JPAC effort

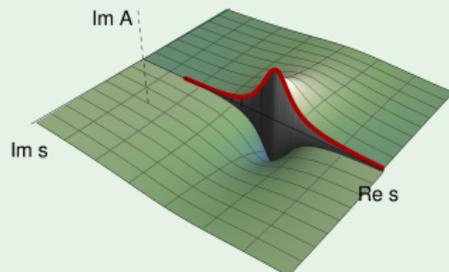
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Amplitude analysis



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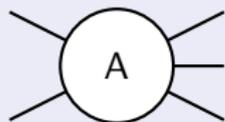
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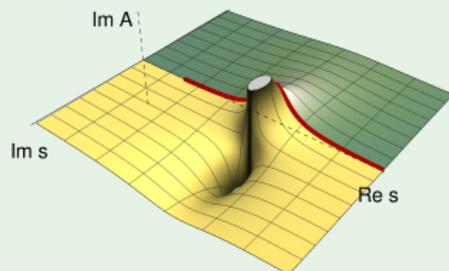
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Amplitude analysis



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General properties of the scattering amplitude



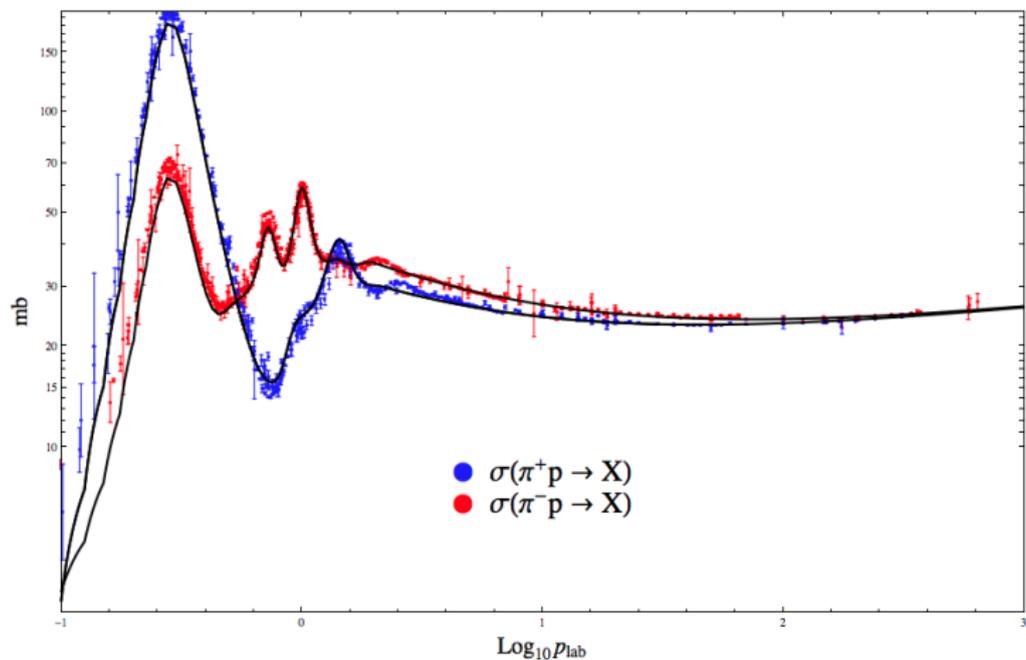
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Two regimes of scattering

Hadronic duality

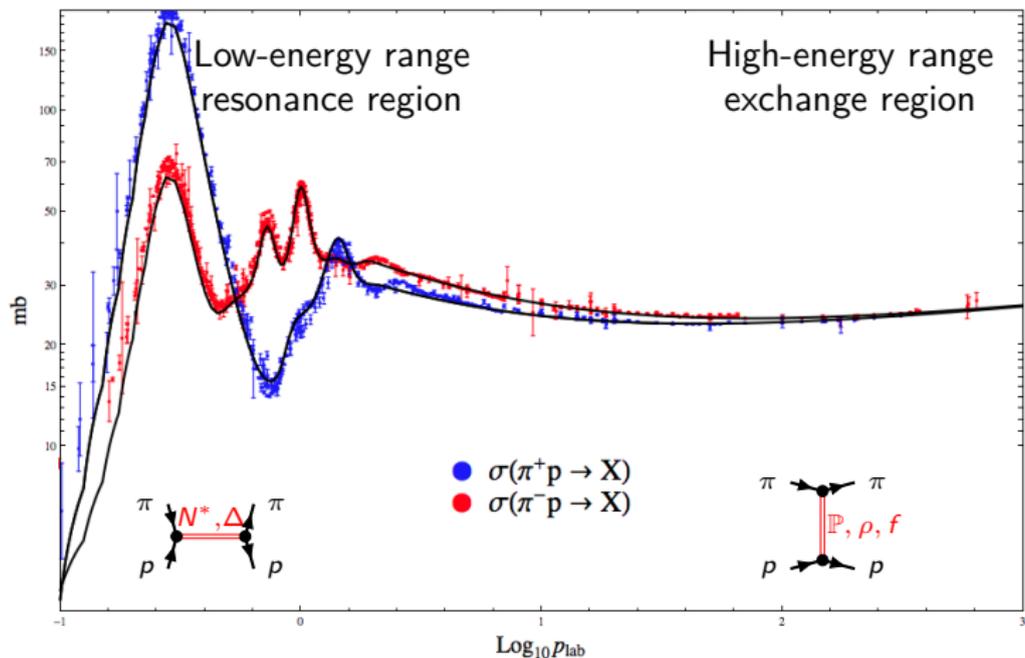
[V.Mathieu, et al., PRD92 (2015), 074004]



Two regimes of scattering

Hadronic duality

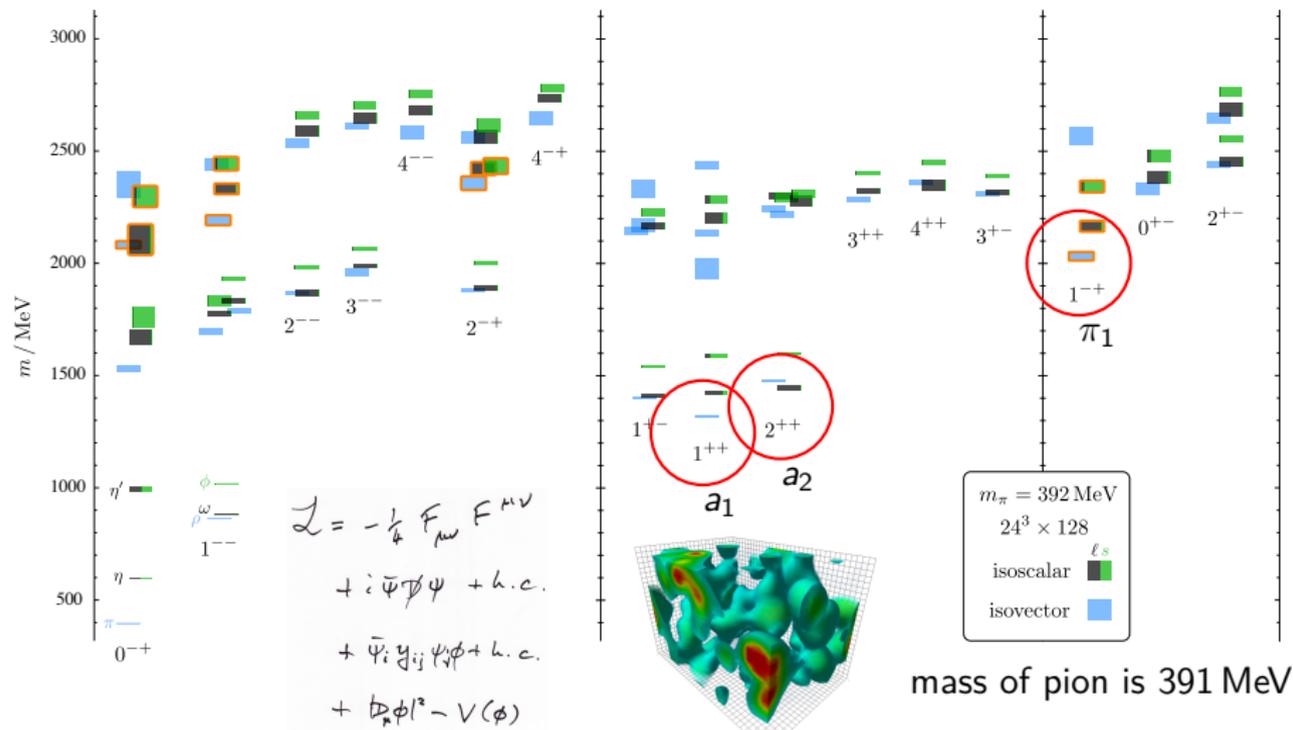
[V.Mathieu, et al., PRD92 (2015), 074004]



Hadronic excitations

Results of lattice QCD

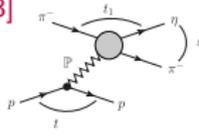
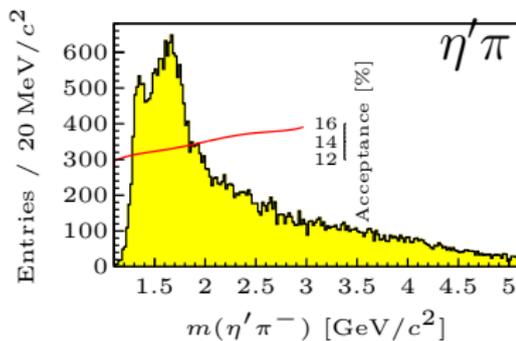
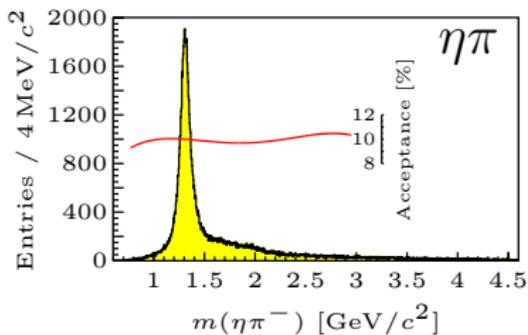
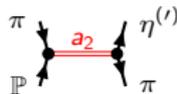
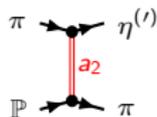
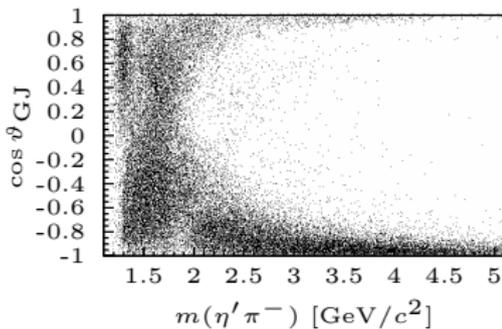
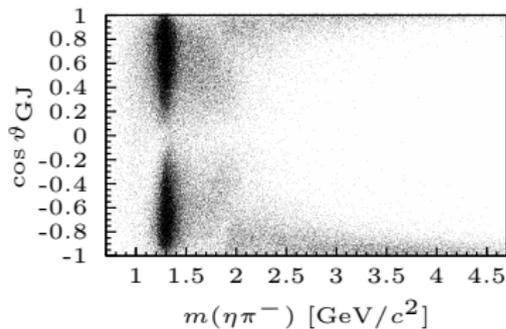
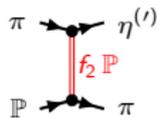
[Dudek et al., PRD 88, 094505 (2013)]



$\eta^{(\prime)}\pi$ analyses

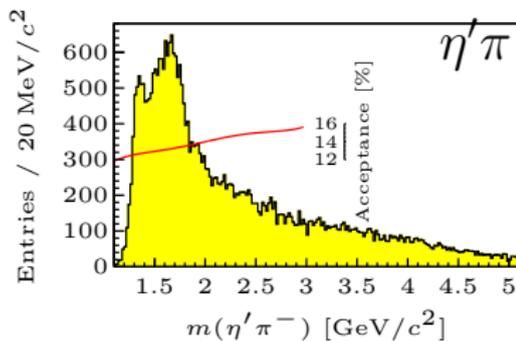
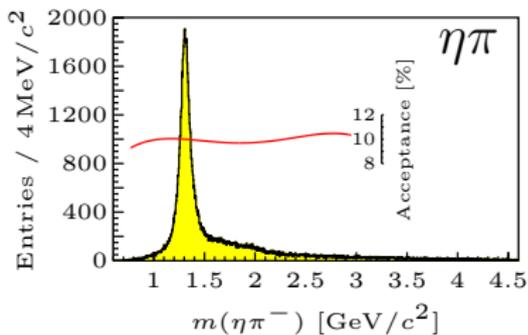
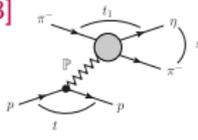
$\eta\pi$ vs $\eta'\pi$ at COMPASS

[[COMPASS) PLB 740 (2015) 303]

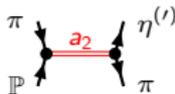
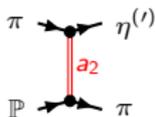
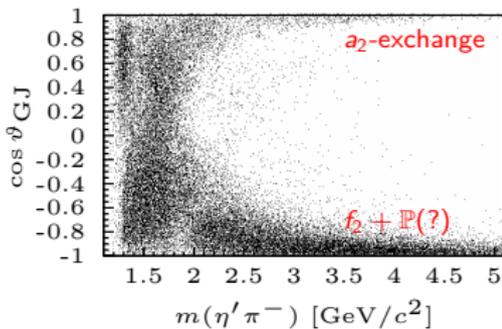
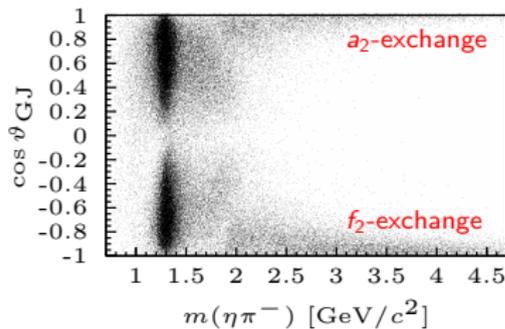
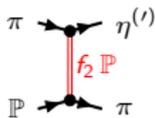
Resonance
production η -forward
production π -forward
production
 $\eta^{(\prime)}(0^-)\pi(0^-)$, $J^{PC} = L^{P+} \Rightarrow \cos\theta_{GJ}$ asymmetry \Rightarrow exotic waves!

$\eta\pi$ vs $\eta'\pi$ at COMPASS

[[COMPASS) PLB 740 (2015) 303]



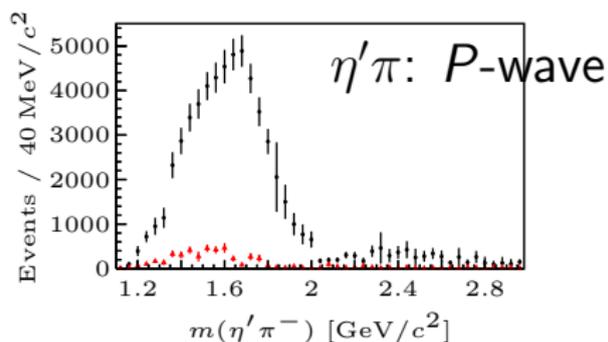
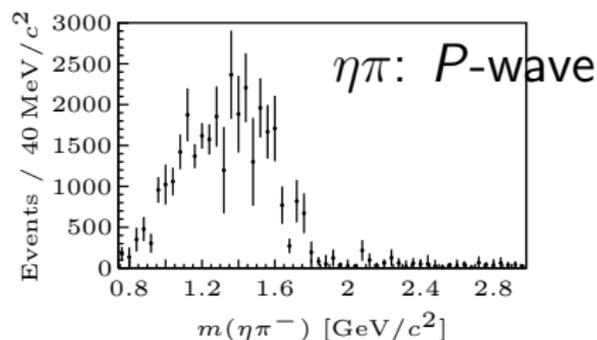
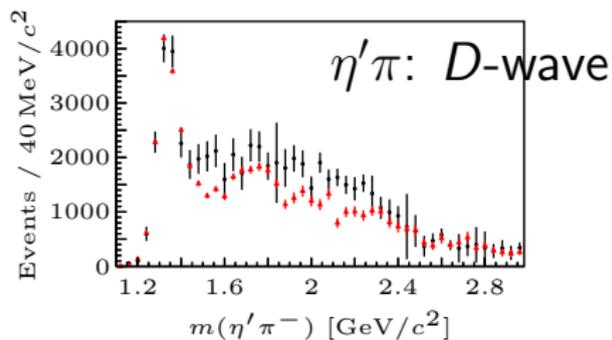
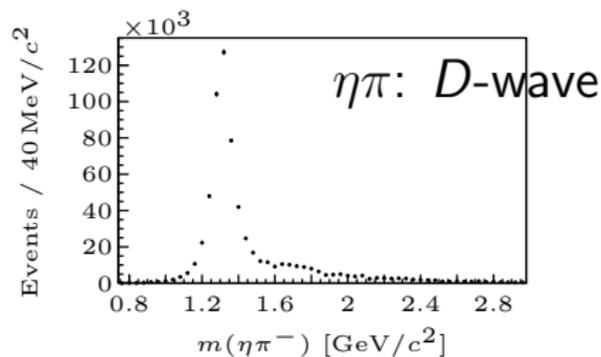
Resonance production

 η -forward production π -forward production
 $\eta^{(\prime)}(0^-)\pi(0^-)$, $J^{PC} = L^{P+} \Rightarrow \cos\theta_{GJ}$ asymmetry \Rightarrow exotic waves!

$\eta^{(\prime)}\pi$ partial wave analysis

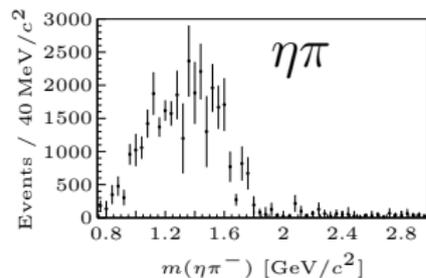
Two-main contribution: P - and D -waves

[(COMPASS) PLB 740 (2015) 303]



PDG status: exotic π_1 states

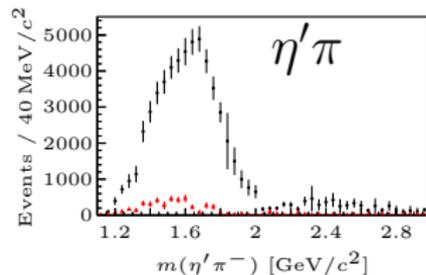
Two candidates



$$\pi_1(1400) \quad I^G(J^{PC}) = 1^-(1^{-+})$$

See also the mini-review under non- $q\bar{q}$ candidates in [PDG 2006](#), Journal of Physics G33 1 (2006).

$\pi_1(1400)$ MASS	1354 ± 25 MeV (S = 1.8)		
$\pi_1(1400)$ WIDTH	330 ± 35 MeV		
Decay Modes			
Mode	Fraction (Γ_i / Γ)	Scale Factor/ Conf. Level	P (MeV/c)
Γ_1	$\eta\pi^0$	seen	557
Γ_2	$\eta\pi^-$	seen	556
Γ_3	$\eta'\pi$		318



$$\pi_1(1600) \quad I^G(J^{PC}) = 1^-(1^{-+})$$

$\pi_1(1600)$ MASS	1662^{+8}_{-9} MeV		
$\pi_1(1600)$ WIDTH	241 ± 40 MeV (S = 1.4)		
Decay Modes			
Mode	Fraction (Γ_i / Γ)	Scale Factor/ Conf. Level	P (MeV/c)
Γ_1	$\pi\pi\pi$	seen	803
Γ_2	$\rho^0\pi^-$	seen	641
Γ_3	$f_2(1270)\pi^-$	not seen	318
Γ_4	$b_1(1235)\pi$	seen	357
Γ_5	$\eta'(958)\pi$	seen	543
Γ_6	$f_1(1285)\pi$	seen	314

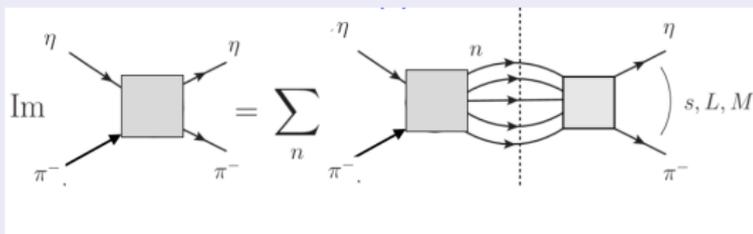
Amplitude for $\eta\pi$ production [A.Jackura,MM,A.Pilloni,et al. (JPAC-COMPASS),

PLB779, 464-472]

N -over- D method

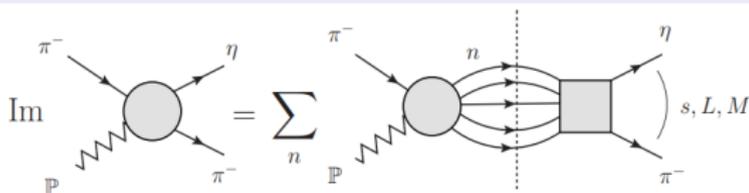
Scattering amplitude: $\eta\pi \rightarrow \eta\pi$, D -wave

$$T = \frac{N(s)}{D(s)}$$



Production amplitude: $\pi\mathbb{P} \rightarrow \eta\pi$, D -wave

$$a = \frac{n(s)}{D(s)}$$

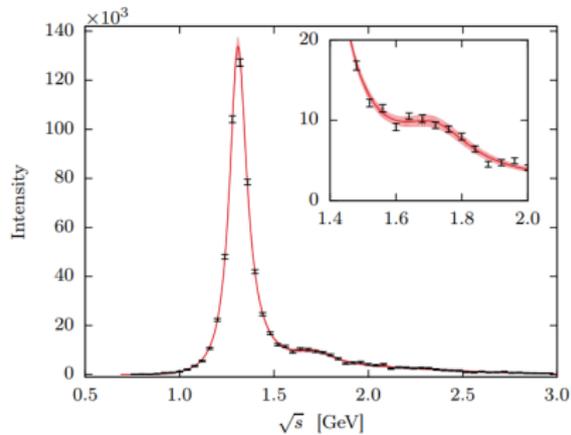


- $D(s)$ is universal, has only the right-hand cut.
- $N(s)$ and $n(s)$ have the left-hand cut only (exchanges)

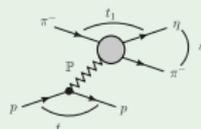
Tensor mesons ($J^{PC} = 2^{++}$)

Advanced $\eta\pi$ analysis

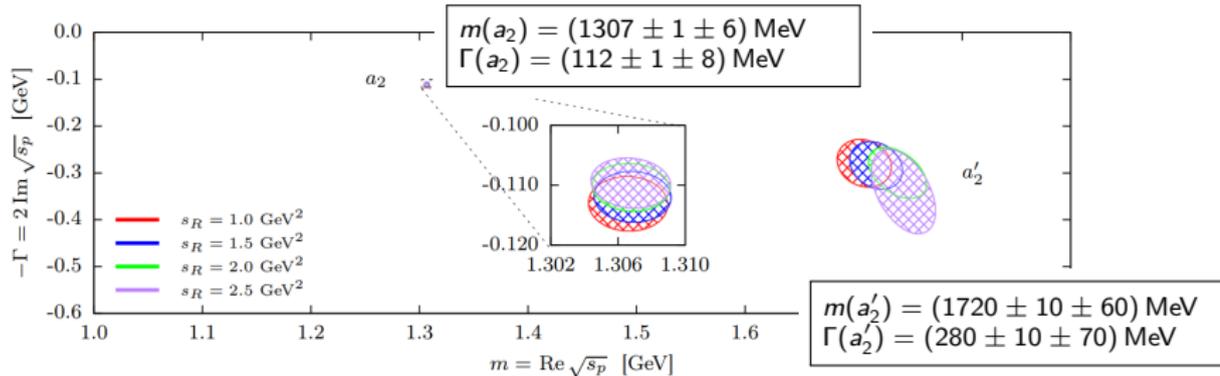
[A.Jackura,MM,A.Pilloni,et al. (JPAC-COMPASS), PLB779, 464-472]



Single channel: $\eta\pi$ D -wave



- Elastic unitarity
- Two CDD-poles



Coupled-channel amplitude [A.Rodas,A.Pilloni,MM,et al. (JPAC), PRL122 (2019)]

Scattering amplitude: $\eta^{(\prime)}\pi \rightarrow \eta^{(\prime)}\pi$, P/D -waves

$$T = \frac{N(s)}{D(s)}$$

$$\rho N_{ki}^J(s') = \delta_{ki} \frac{(p_{\eta^{(\prime)}\pi} \sqrt{s}/2)^{2J+1}}{(s' + s_L)^{2J+1+\alpha}},$$

- 2 K -matrix pole for D -wave
- 1 K -matrix pole for P -wave

$$D_{ki}^J(s) = [K^J(s)^{-1}]_{ki} - \frac{s}{\pi} \int_{s_k}^{\infty} ds' \frac{\rho N_{ki}^J(s')}{s'(s' - s - i\epsilon)}$$

Production amplitude: $\pi\mathbb{P} \rightarrow \eta^{(\prime)}\pi$, P/D -waves

$$a = \frac{n(s)}{D(s)}$$

$$a_i^J(s) = q^{J-1} p_i^J \sum_k n_k^J(s) [D^J(s)^{-1}]_{ki}$$

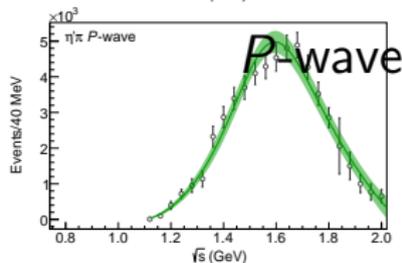
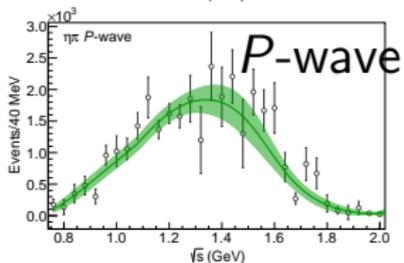
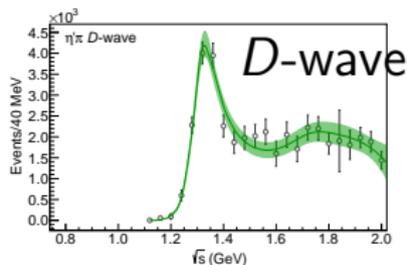
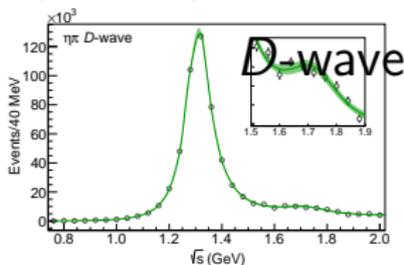
- left poles to model unknown production function $n(s)$

- $D(s)$ has only the right-hand cut.
- $N(s)$ and $n(s)$ have the left-hand cut only (exchanges)

Fit to the data

[A.Rodas,A.Pilloni,MM,et al. (JPAC), PRL122 (2019)]

$\chi^2/\text{ndf} = 162/122$, the band - 2σ bootstrap error

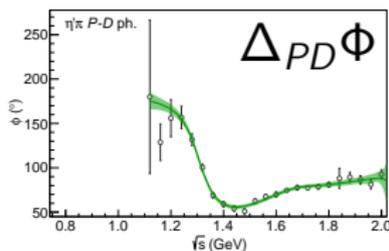
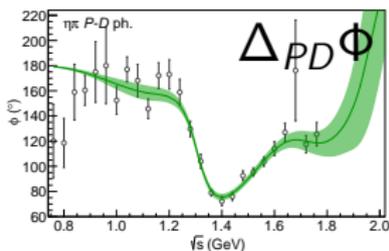


D-wave difference

- Kinematics
($m_{\eta'} > m_{\eta}$)
- ⇒ Same amplitude.

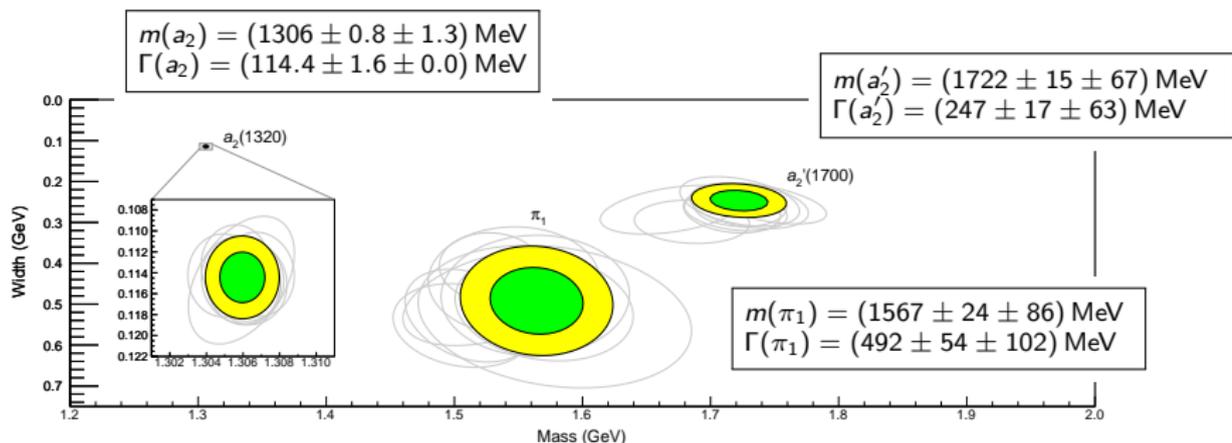
P-wave difference

- production mechanism
- + kinematics.



Results: pole positions

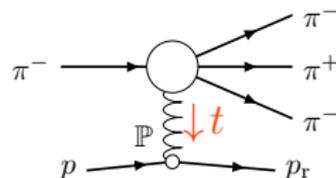
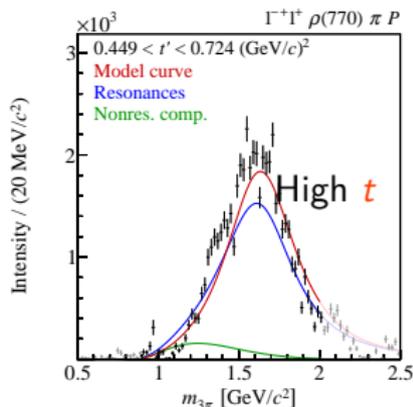
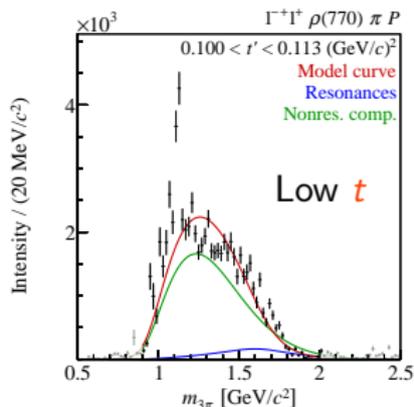
[A.Rodas,A.Pilloni,MM,et al. (JPAC), PRL122 (2019)]



- Change parametrization of the denominator $\rho N_{ki}^J(s') = \delta_{ki} \frac{(p_{\eta^{(\prime)}\pi} \sqrt{s}/2)^{2J+1}}{(s'+s_L)^{2J+1+\alpha}}$,
 - ▶ $s_R = 1 \text{ GeV} \rightarrow 0.8, 1.8 \text{ GeV}$.
 - ▶ $\alpha = 2 \rightarrow 1 \text{ GeV}$.
 - ▶ Different function, $\rho N_{ki}^J(s') = \delta_{ki} Q_J(z_{s'}) s'^{-\alpha} \lambda^{-1/2}(s', m_{\eta^{(\prime)}}^2, m_{\pi}^2)$
- Change of parameters in the numerator $n(s)$
 - ▶ Effective transferred momentum $t_{\text{eff}} = -0.1 \text{ GeV}^2 \rightarrow -0.5 \text{ GeV}^2$.
 - ▶ Order of the polynomial 3rd-order \rightarrow 4th-order.

Same π_1 as in 3π ?

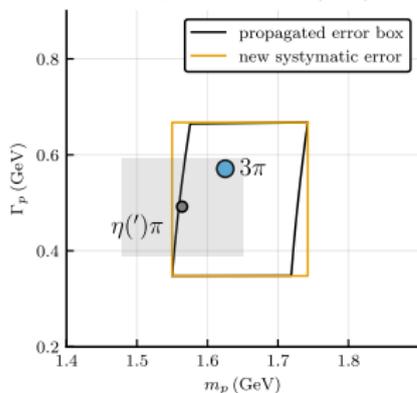
[See a talk of B.Ketzer, afternoon]



The COMPASS fit:

- Signal by BW amplitude
- Flexible background

Pole parameters of $\pi_1(1600)$



Consistent results on $\pi_1(1600)$ pole:

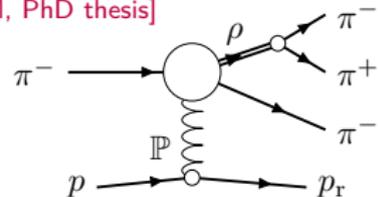
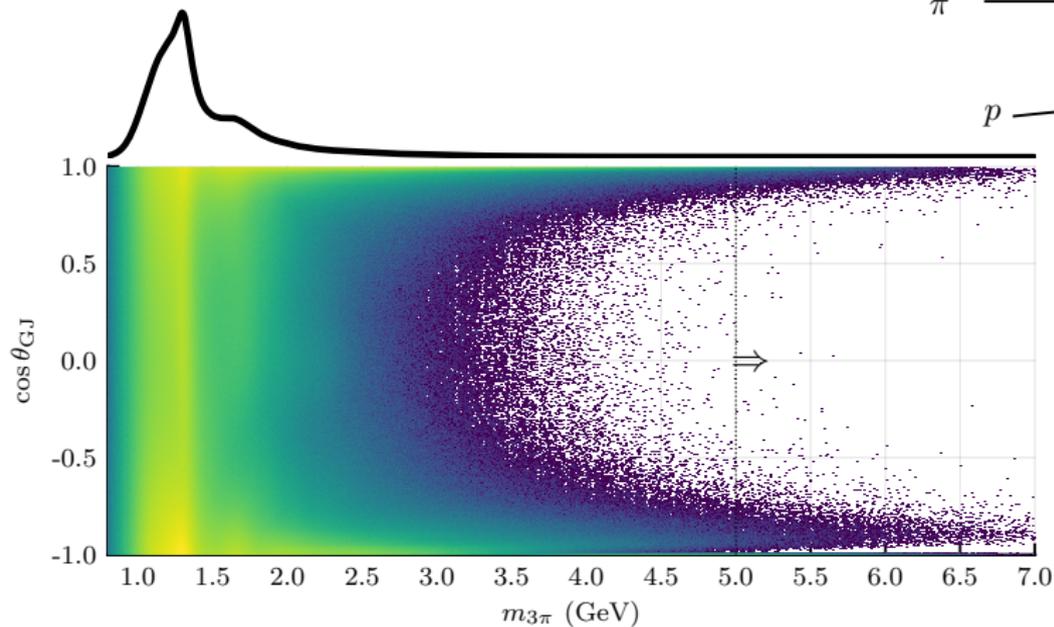
- $\rho\pi$ Breit-Wigner parameters
 \Rightarrow pole position
- $\eta^{(\prime)}\pi$ systematic margins

Three-pions physics

Diffractive production of 3π off proton target

Forward-background scattering

[COMPASS data, MM, PhD thesis]

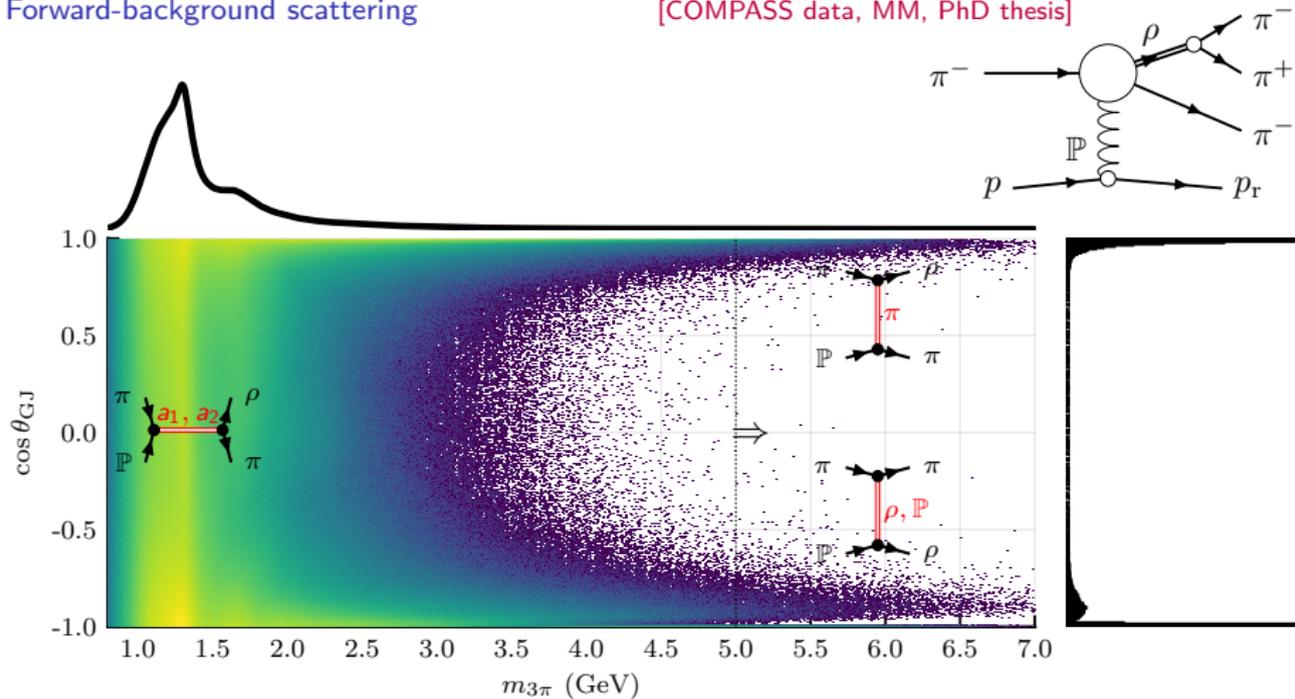


The high-energy exchange processes penetrate to the low energy and make resonance characterization difficult

Diffractive production of 3π off proton target

Forward-background scattering

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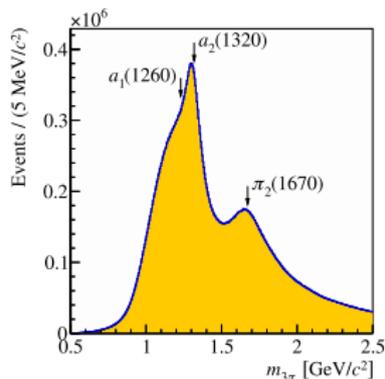


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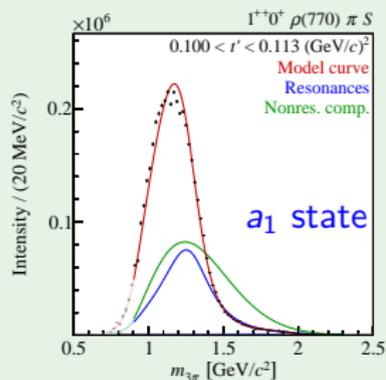
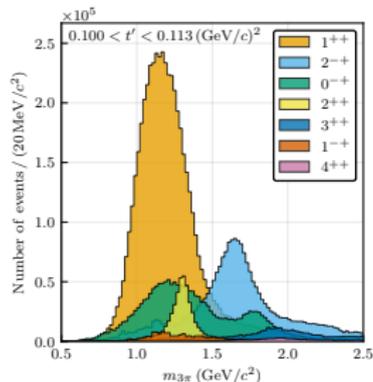
Three-pion resonances

[see talk by D.Ryabchikov, afternoon]

COMPASS PWA [PRD95 (2017) 032004]



Partial Wave Analysis →



COMPASS Fit:

- Signal by BW amplitude $\sim 50\%$
- Background $\sim 50\%$

$$m_{a_1}^{(BW)} = (1299_{-28}^{+12}) \text{ MeV,}$$

$$\Gamma_{a_1}^{(BW)} = (380 \pm 80) \text{ MeV.}$$

Large uncertainty due to unknown background.

Model for the forward scattering

[MM, A.Jackura (JPAC) in preparation]

Deck effect

$$= (T_{\pi_1 \rho})_{\lambda \lambda'} \frac{1}{m_\pi^2 - t_1} T_{\pi_2 \pi_3} \quad (1)$$

- Two diagrams (π^- symmetrization)

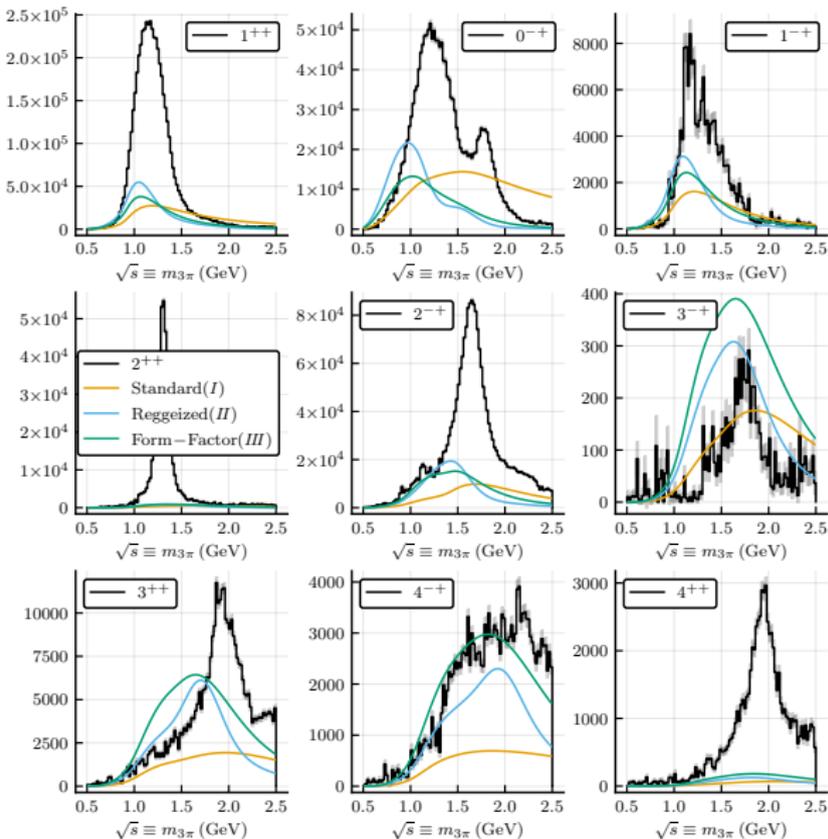
$$\mathcal{B}_{\lambda \lambda'} = \mathcal{B}_{\lambda \lambda'}^{(1)} + \mathcal{B}_{\lambda \lambda'}^{(3)}$$

- High energy $p\pi$ scattering
- $\pi\pi$ scattering dominated by resonances in lower partial waves
 - ▶ Relative strength of S , P , D -waves is controlled by unitarity

$$\mathcal{B}^{(1)} = s_{\pi\rho} F(t) \frac{\text{FF}(t_1)}{m_\pi^2 - t_1} \left[\frac{2}{3} t^{(\sigma_1, f_0)}(\sigma_1) + 3 t^{(\rho)}(\sigma_1, t_1) P_1(\cos \theta_{\pi\pi}) + \frac{10}{3} t^{(f_2)}(\sigma_1, t_1) P_2(\cos \theta_{\pi\pi}) \right].$$

Comparison with the COMPASS data

[MM PhD thesis]



Pion propagator:

- Standard Deck

$$\frac{1}{m_{\pi}^2 - t_1},$$

- Regge Deck

$$\frac{e^{-i\pi\alpha(t_1)/2}}{m_{\pi}^2 - t_1} \left(\frac{s' - u'}{2s_{\text{Sc}}} \right)^{\alpha(t_1)}$$

- Form-factored Deck

$$\frac{e^{bt_1}}{m_{\pi}^2 - t_1},$$

$a_1(1260)$ state – isospin partner of ρ $a_1(1260)$ WIDTH[INSPIRE search](#)

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
250 to 600	OUR ESTIMATE			
389 ± 29	OUR AVERAGE	Error includes scale factor of 1.3.		
430 ± 24 ± 31		DARGENT 2017	RVUE	$D^0 \rightarrow \pi^- \pi^+ \pi^- \pi^+$
367 ± 9 ⁺²⁸ ₋₂₅	420k	ALEKSEEV 2010	COMP	$190 \pi^- \rightarrow \pi^- \pi^- \pi^+ P b'$
••• We do not use the following data for averages, fits, limits, etc. •••				
410 ± 31 ± 30		1 AUBERT 2007AU	BABR	$10.6 e^+ e^- \rightarrow \rho^0 \rho^\pm \pi^\mp \gamma$
520 - 680	6360	2 LINK 2007A	FOCS	$D^0 \rightarrow \pi^- \pi^+ \pi^- \pi^+$
480 ± 20		3 GOMEZ-DUMM 2004	RVUE	$\tau^+ \rightarrow \pi^+ \pi^+ \pi^- \nu_\tau$
580 ± 41	90k	SALVINI 2004	OBLX	$\bar{p} p \rightarrow 2 \pi^+ 2 \pi^-$
460 ± 85	205	4 DRUTSKOY 2002	BELL	$B^{(*)} K^- K^0$
814 ± 36 ± 13	37k	5 ASNER 2000	CLE2	$10.6 e^+ e^- \rightarrow \tau^+ \tau^-$, $\tau^- \rightarrow \pi^- \pi^0 \pi^0 \nu_\tau$

$a_1(1260)$ state – isospin partner of ρ $a_1(1260)$ WIDTH[INSPIRE search](#)

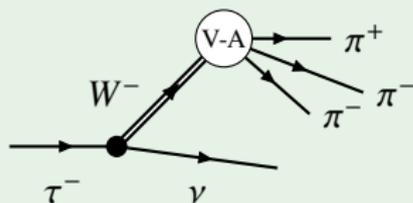
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$a_1(1260)$ state – isospin partner of ρ $a_1(1260)$ WIDTH

INSPIRE search

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$$\tau^- \rightarrow \pi^- \pi^+ \pi^- \nu$$

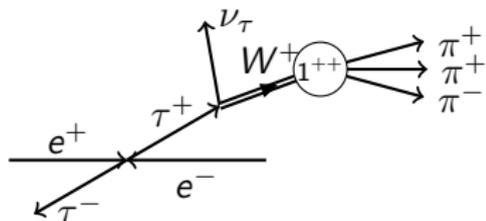
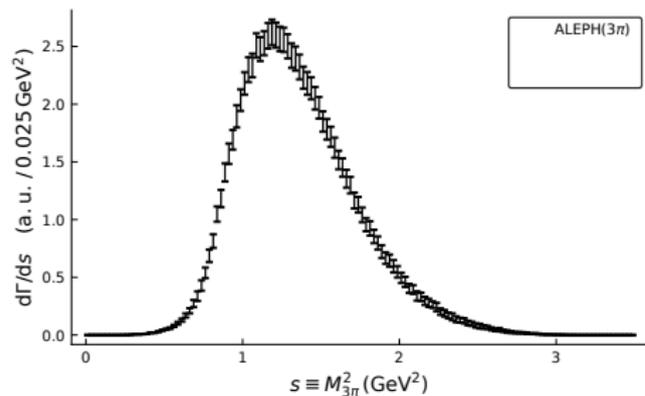


- V-A: Vector (1^{--}) or Axial (1^{++})
- Isospin 1 due to the charge
- Negative G-parity \Rightarrow positive C-parity

$$\Rightarrow J^{PC} = 1^{++}$$

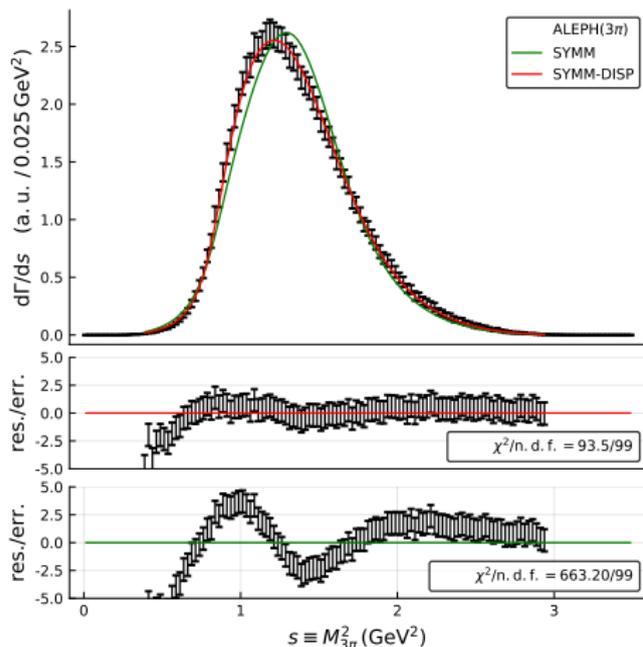
Fit to ALEPH data

[data from ALEPH, Phys.Rept.421 (2005)]



Fit to ALEPH data

[data from ALEPH, Phys.Rept.421 (2005)]



Dispersive model vs Non-dispersive model

- Difference: LH singularities
- The dispersive model fits significantly better

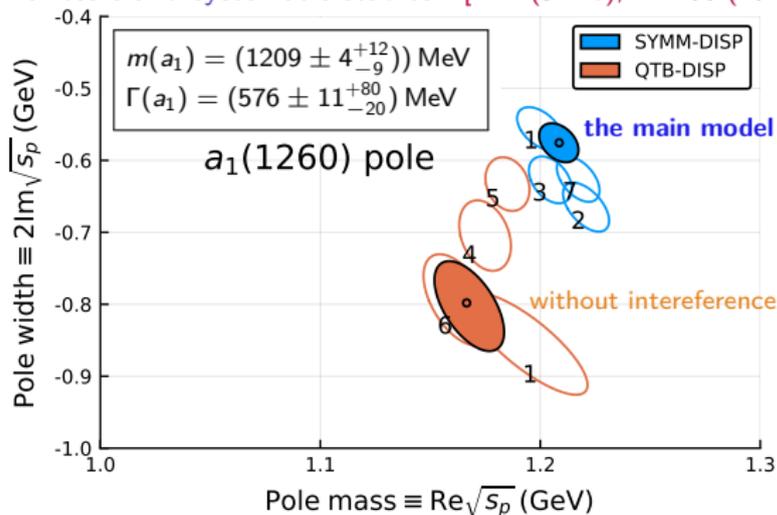
Fit function

$$\chi^2(c, m, g) = (\vec{D} - \vec{M}(c, m, g))^T C_{\text{stat}}^{-1} (\vec{D} - \vec{M}(c, m, g)),$$

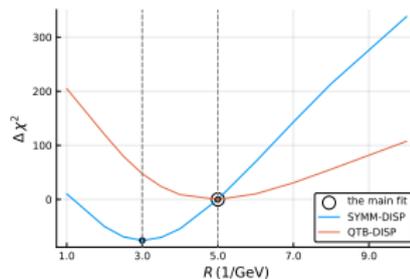
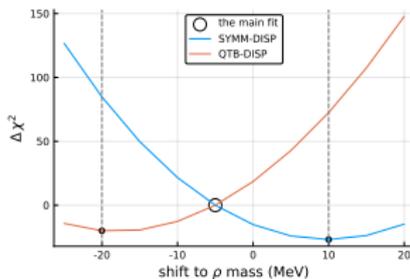
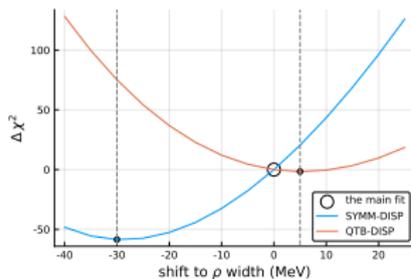
- Stat. cov. matrix is used in the fit
- Syst. cov. matrix – in the bootstrap

First measurement of the $a_1(1260)$ pole position

The result and systematic studies [MM (JPAC), PRD98 (2018), 096021]



#	Fit studies
1	$s < 2 \text{ GeV}^2$
2	$R' = 3 \text{ GeV}^{-1}$
3	$m'_\rho = m_\rho + 10 \text{ MeV}$
4	$m'_\rho = m_\rho - 10 \text{ MeV}$
5	$m'_\rho = m_\rho - 20 \text{ MeV}$
6	$\Gamma'_\rho = \Gamma_\rho + 5 \text{ MeV}$
7	$\Gamma'_\rho = \Gamma_\rho - 30 \text{ MeV}$



Summary

Meson spectroscopy

- Using hadronic scattering as QCD excitation laboratory.
- Mapping gluonic degrees of freedom to structures of excited states is an essential test of QCD.
- Non-perturbative methods are required

Recent impact of JPAC to light meson spectroscopy

Extensive analyses and extraction of resonance poles:

- Tensor states: $a_2(1320)$ and $a_2(1700)$
- Establishing single exotic $\pi_1(1600)$
- Ground axial state $a_1(1260)$

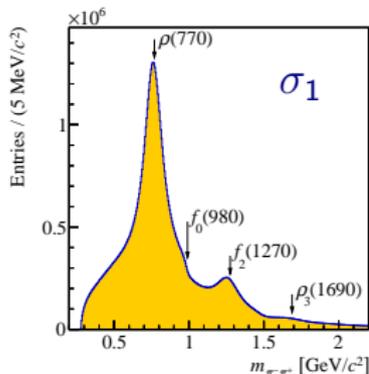
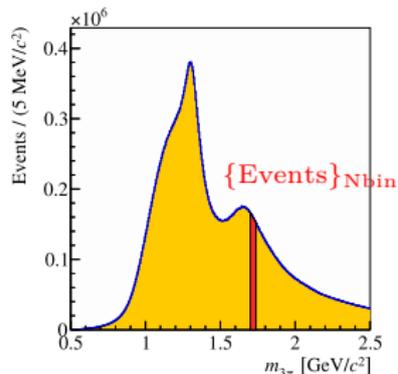
JPAC effort

- > 50 research papers in PRD, PLB, PRL, EJPC (> 10 in 2018)
- > 100 invited talks and seminars
- Collaboration with GlueX, CLAS12, COMPASS, MAMI, BaBar, LHCb,...
- Summer Schools on Reaction Theory (2015, 2017)

Thank you

Three-particles PW technique

COMPASS@PWA

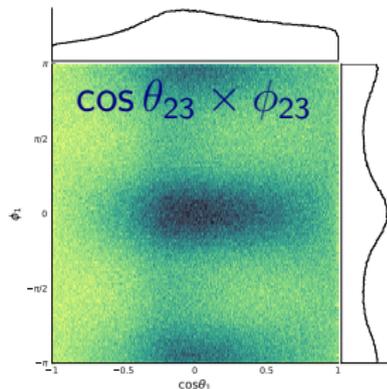
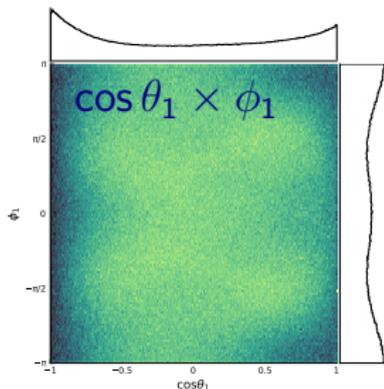


Model: a sum of 88* partial waves

$$A = c_1 \left(\begin{array}{c} \rho^{++} \\ \rho^{+-} \\ \rho^{00} \\ \rho^{-+} \\ \rho^{--} \end{array} \right) + c_2 \left(\begin{array}{c} f_2^{++} \\ f_2^{+-} \\ f_2^{00} \\ f_2^{-+} \\ f_2^{--} \end{array} \right) + c_3 \left(\begin{array}{c} f_0^{++} \\ f_0^{+-} \\ f_0^{00} \\ f_0^{-+} \\ f_0^{--} \end{array} \right) + \dots$$

Likelihood Fit: a product of probability per event

$$L = \prod_{e=1}^{N_{events}} \frac{A(\tau_e)}{\int d\tau_e A(\tau_e)}, \quad \tau = (\sigma_1, \cos \theta_1, \phi_1, \cos \theta_{23}, \phi_{23})$$



Tour to the complex plane

[MM (JPAC), PRD98 (2018), 096021]

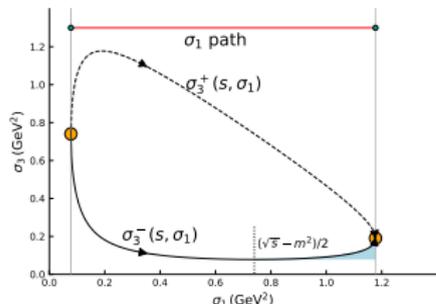
Analytical continuation

$$|t_I^{-1}(s)| = \left| \frac{m^2 - s}{g^2} - i \left(\frac{\tilde{\rho}(s)}{2} + \rho(s) \right) \right|.$$

- Analytical continuation of $\rho(s)$: integral over the Dalitz plot for the complex energy

$$\rho(s) = \sum_{\lambda} \int d\Phi_3 \left| f_{\rho}(\sigma_1) d_{\lambda 0}(\theta_{23}) - f_{\rho}(\sigma_3) d_{\lambda 0}(\hat{\theta}_3 + \theta_{12}) \right|^2$$

- Analytic continuation of ρ -meson decay amplitude f_{ρ}
 - Breit-Wigner amplitude with the dynamic width
 - P -wave Blatt-Weisskopf factors



Tour to the complex plane

[MM (JPAC), PRD98 (2018), 096021]

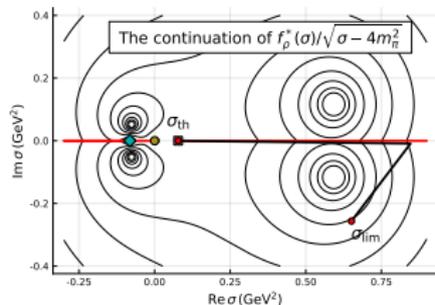
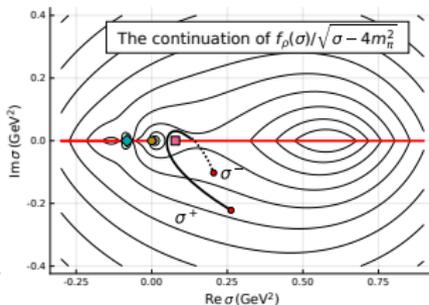
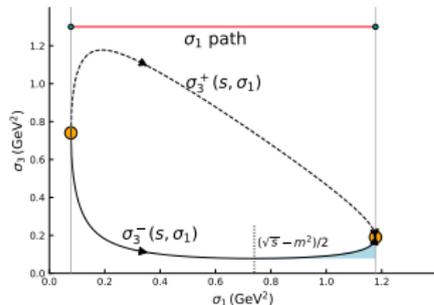
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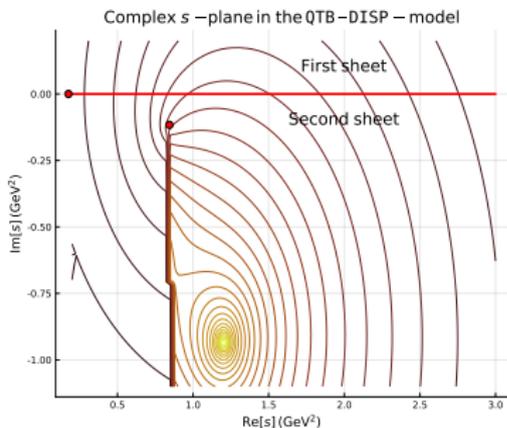
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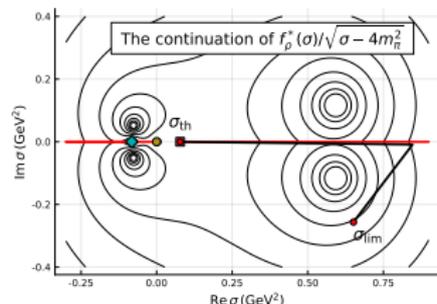
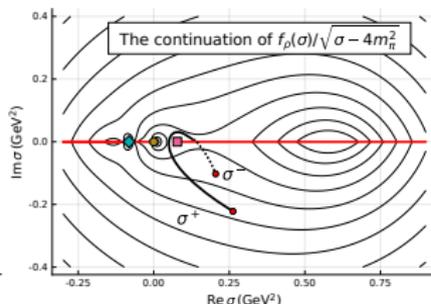
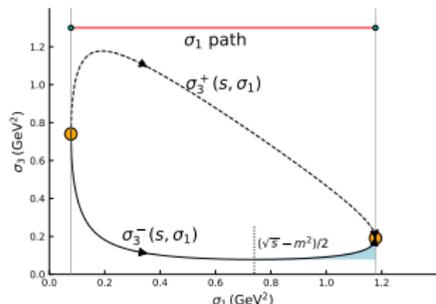
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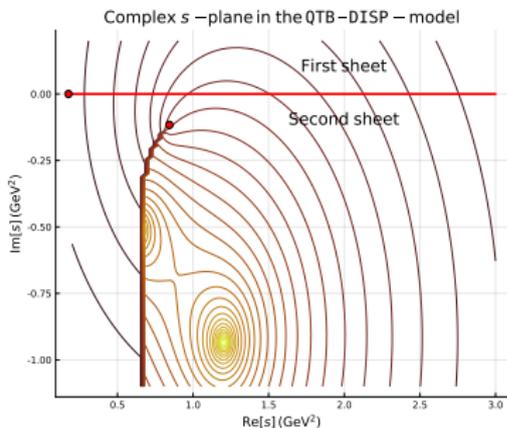
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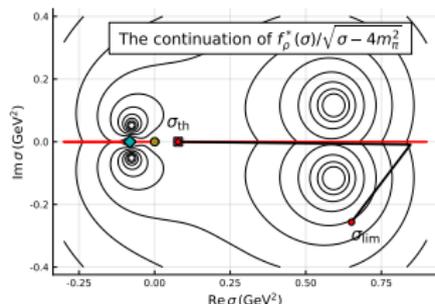
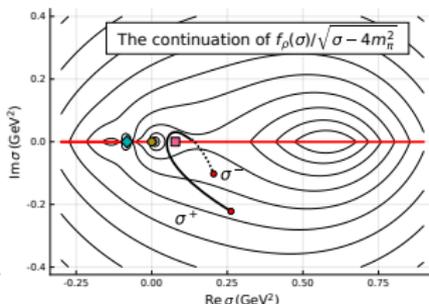
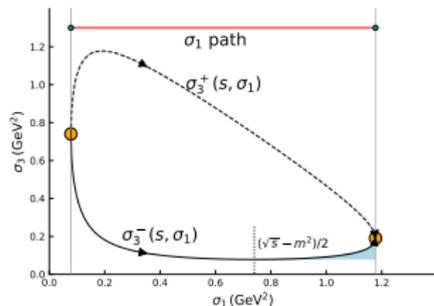
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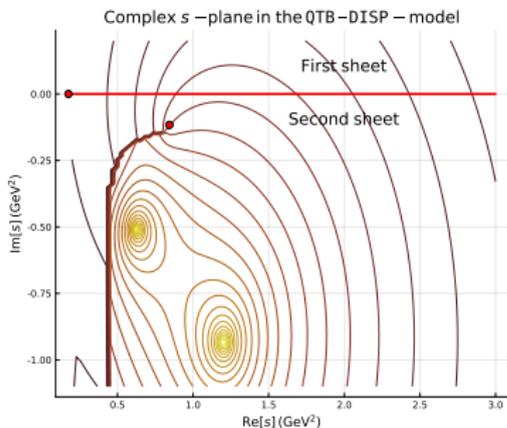
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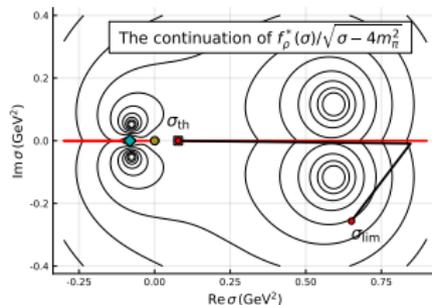
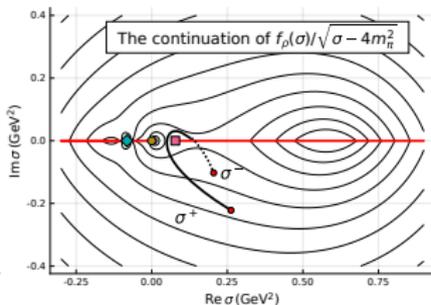
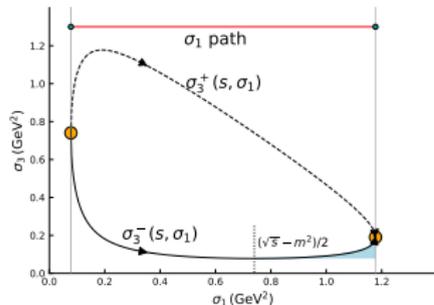
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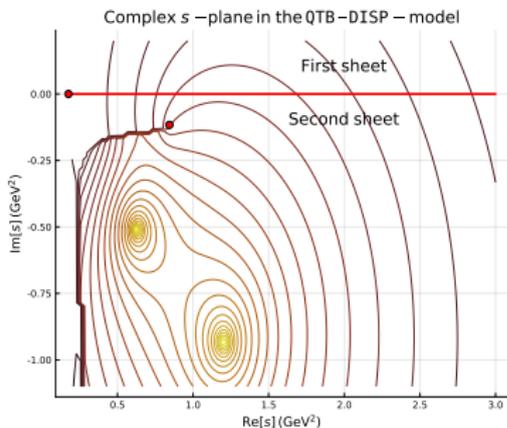
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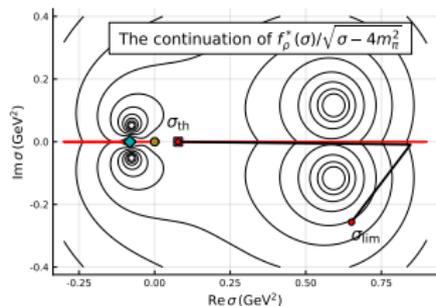
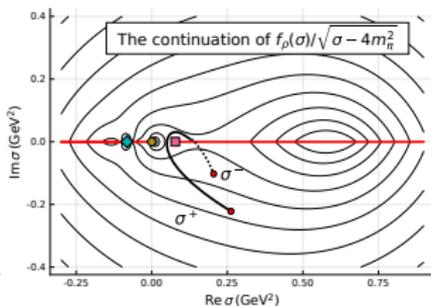
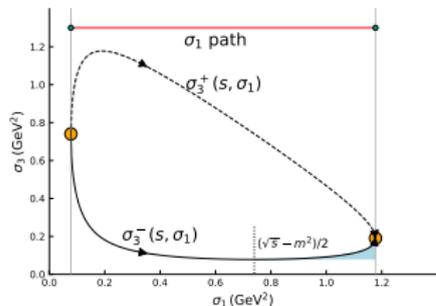
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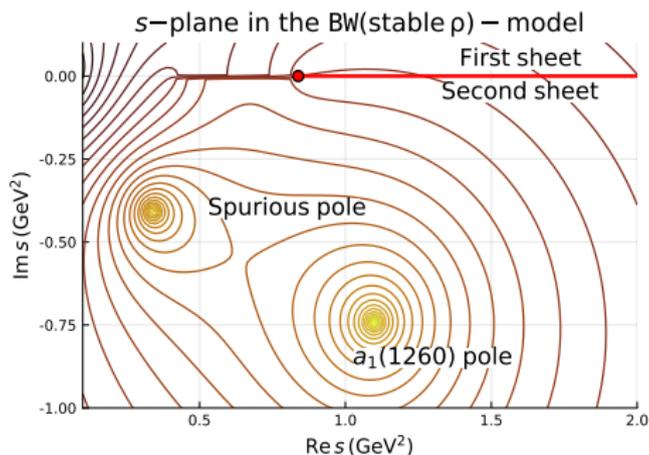


The spurious pole in the Breit-Wigner model

Energy dependent width, stable particles

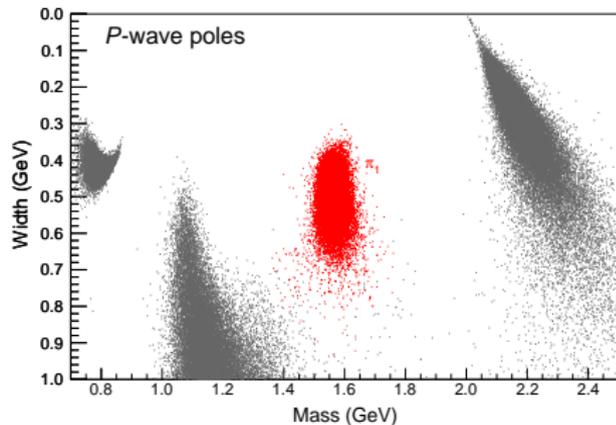
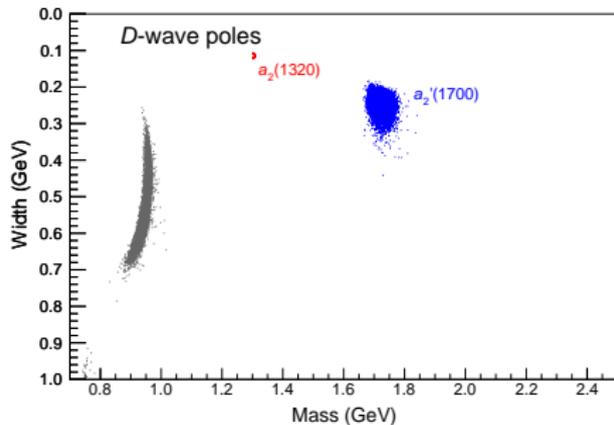
$$t(s) = \frac{1}{m^2 - s - im\Gamma(s)}, \quad \Gamma(s) = \Gamma_0 \frac{p(s)}{p(m^2)} \frac{m}{\sqrt{s}}, \quad p(s) = \frac{\sqrt{(s - (m_1 + m_2)^2)(s - (m_1 - m_2)^2)}}{2\sqrt{s}}.$$

Example: $m_1 = 140 \text{ MeV}$, $m_2 = 770 \text{ MeV}$, $m = 1.26 \text{ GeV}$, $\Gamma_0 = 0.5 \text{ GeV}$



Live demo

Bootstrap: stability of the poles



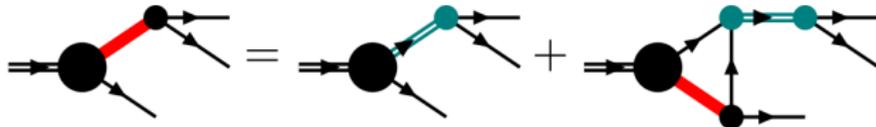
- Statistical bands are obtained by 50k bootstrap samples

Subchannel dynamics

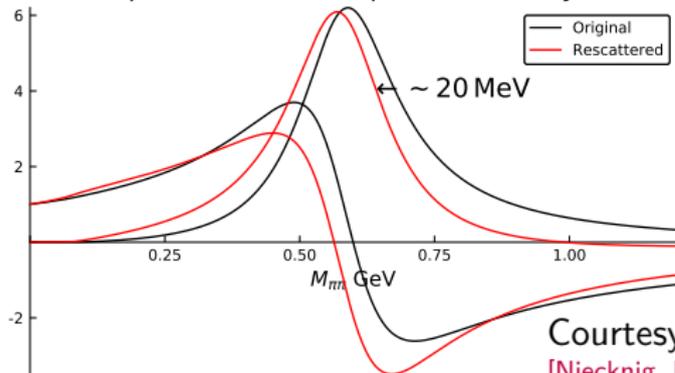
Khuri-Treiman equations

Consistency equations for the isobar lineshape

- Governed by two-body unitarity
- Model: only RHC for the isobar amplitude
- Uses Analyticity / Cauchy theorem / Omnès trick



ρ – meson lineshape in ω – decay



KT analysis of $\omega \rightarrow 3\pi$

- $\rho \rightarrow \rho$ rescattering
- Solved by iteration

Courtesy of Tobias Isken

[Niecknig, Kubis, JHEP 1510 (2015) 142]