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# High-brightness Radiation Enabled by Dielectric Laser Accelerator

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# OUTLINE

- 1. Dielectric laser accelerator (DLA)
- 2. Short-bunch radiation
- 3. Brilliance of DLA-driven coherent undulator radiation
- 4. Conclusions

# Accelerator on a Chip (ACHIP)

#### https://achip.stanford.edu/

- Direct-field acceleration
- High laser damage resistance on dielectric





#### Partner Institutions



# **Envisaged DLA-driven Coherent Undulator Radiation**

#### Dielectric laser accelerator (DLA)

**Dielectric undulator** 





electron bunch length  $\sim 1 \text{ nm}$ Bunch Charge =  $\sim 1 \text{ fC}$ Gamma =  $\sim 500 \text{ MeV}$  T. Plettner, R. L. Byer, Phys. Rev. ST Accel. Beams **11**, 030704 (2008).

← 100 cm

 $\lambda_{\rm u} = 1 \, {\rm mm} \, ({\rm N}_{\rm u} = 1000)$ 

B<sub>peak</sub> = 3 T (subject to laser damage)

 $a_u = \sim 0.22$  (undulator parameter) \*good scheme to keep  $a_u$  for small  $\lambda_u$ 

#### **Pulse structures from DLA vs. RF Accelerator**



### Short-bunch enhanced radiation - superradiance

Total radiation Spectral Energy in one driver laser pulse  $W_N = N_I [N_\mu + N_\mu (N_\mu - 1)b^2(\omega)]W_1$  $W_{1,N}$  radiation spectral energy of 1 & N electrons  $b(\omega)$ : bunching factor, Fourier transform of the micro-bunch profile  $N_{\mu}$ : # of electrons in a DLA micro-bunch  $N_{l}$ : # of micro-bunches in a driver-laser pulse  $N_L \times N_\mu$ : total number of electrons in a driver laser pulse For Gaussian bunch  $f(t) = \frac{\exp(-t^2/2\tau_b^2)}{\sqrt{2\pi\tau_b}}$  $b^2(\omega)$ 0.8  $\sim 1/\tau_{h}$  $\Rightarrow b(\omega) = \exp\left(-\frac{\omega^2 \tau_b^2}{2}\right)$ 0.4

0.2

0.2 0.4 0.6 0.8 1 1.2 1.4 1.6

 $\omega \tau_{h}$ 

1.8

## **Brilliance of undulator radiation**

photons/s/mm<sup>2</sup>/mrad<sup>2</sup>/0.1%BW

$$B(\omega) \propto [(N_{\mu} - 1)|b(\omega)|^2 + 1]) \times N_{\mu}I(A) \times [JJ]$$

Short-bunch enhancement

where 
$$[JJ] \equiv 4M \times [J_0(M) - J_1(M)]^2 \sim 1$$
 with  $M \equiv a_u^2 / 2(1 + a_u^2)$ 

 $N_{\mu}$ : # of  $\mu$ -bunch electrons  $N_{\nu}$ : # of undulator period I: current  $\infty$  bunch rate

Given a design wavelength, it is desirable to have a large  $N_{\mu}$ ,  $N_{\mu}$ , and I (high bunch rate).



a

System parameters for calculating peak and average brilliances of coherent undulator radiation (CUR) driven by DLA

System parameters		DLA CUR	remark
Bunch Charge	fC	0.5	~3000 electrons/bunch
bunch Rep Rate	MHz	100	30 optical cycles in a 100-fs pulse repeating at 3 MHz or 100 cycles in a 300 fs pulse repetition at 1 MHz
Max photon Energy	keV	1.5	For a shorter wavelength, CUR is not effective.
Undulator Length	m	1	1000 undulator periods with a 1-mm period (0.1% bandwidth)
rms Electron Bunch Length	as	1	Wavelength divided by 3500 times scaled from the demonstrated 100 fs bunch length for an RF-accelerator driven SASE FEL
Electron Energy	MeV	~500 MeV	~1-nm radiation wavelength
Undulator parameter	NA	0.217	Laser undulator field (3.3 T, peak) at laser damage of dielectric

# Peak brilliance higher or comparable to 3<sup>rd</sup>-generation light source (due to superradiance)



\*Curves other than DLA CUR are adapted from Zirong Huang, SLAC-PUB-15449.



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~GW circulating power



average brilliance/Watt 10 17 DLA CUR **DLA stands high when** 10<sup>16</sup> normalized to beam power <sup>2</sup> /0.1%BW/Watt) 10 15 **Diffraction-limited 2-9 GeV rings** 10 14 Upgraded 6-7 GeV rings



~GW circulating power (Taiwan photon source)



~25W power

## CONCLUSIONS

I. Dielectric laser accelerator (DLA) is potentially compact, stable, and high-gradient.

- 2. Nano-bunches from DLA permit high-brightness superradiance in the soft x-ray with small  $\gamma$ . DLA-driven CUR is predicted to have a peak brilliance comparable to a synchrotron.
- 3. DLA-driven coherent undulator radiation has a much higher brilliance/watt than a synchrotron.