

Laser-plasma physics of high energy in Lebedev

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Asian Forum for Accelerators and Detectors (AFAD-2021),
Budker Institute of Nuclear Physics, Novosibirsk, Russia
16-18, March 2021

High Energy Laser-Plasma Physics Lab

V. Yu. Bychenkov

High Field Science

S. G. Bochkarev, A. V. Brantov, D. A. Gozhev,
V. F. Kovalev, A. S. Kuratov, M. G. Lobok,
I. I. Metelskii, O. E. Vais



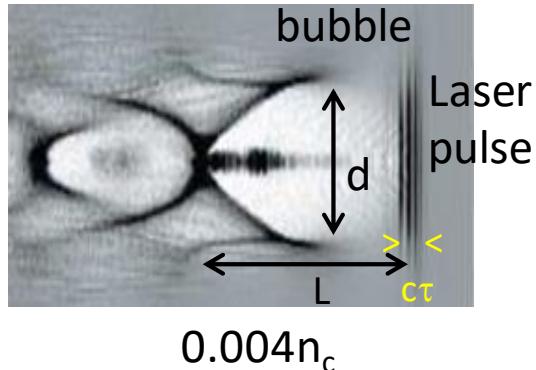
Relativistic optics
Acceleration
Relativistic nanoplasmatics
Secondary emission
Nuclear photonics

ICF Laser-Plasma Interaction

A.V. Brantov, S. I. Glazyrin,
S. A. Karpov, M. A. Rakitina

Instabilities in corona
Fast particles
Nonlocal transport
Nonlinear structures

New regime of wakefield acceleration → Relativistic self-trapping

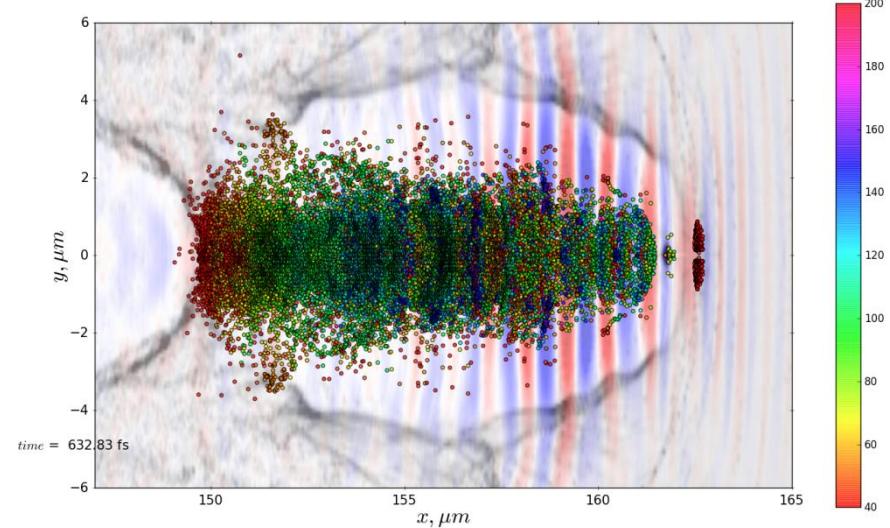
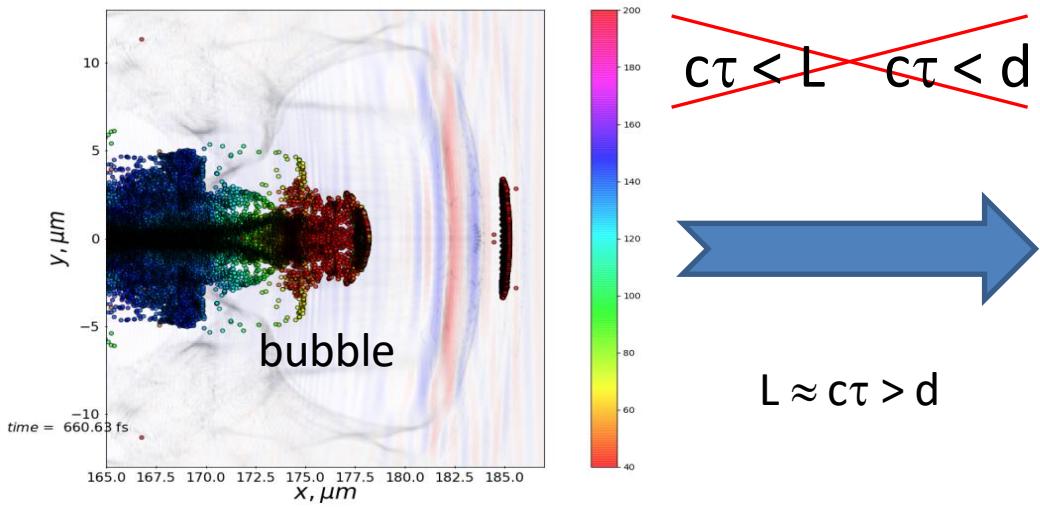


Tajima T. and Dawson J. M. 1979 Phys. Rev. Lett. **43** 267
3D plasma wave

Pukhov A. and Meyer-ter-Vehn J., 2002 Appl. Phys. B: Lasers Opt. **74** 355
"bubble" regime

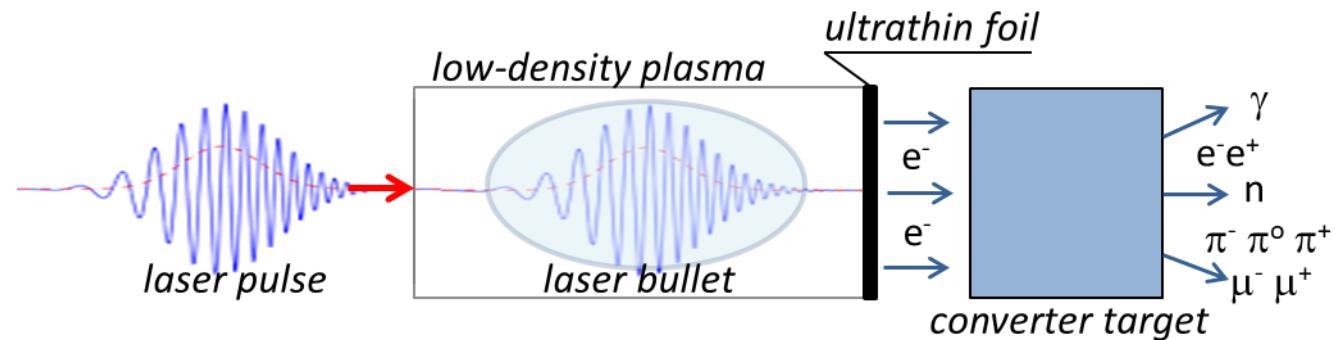
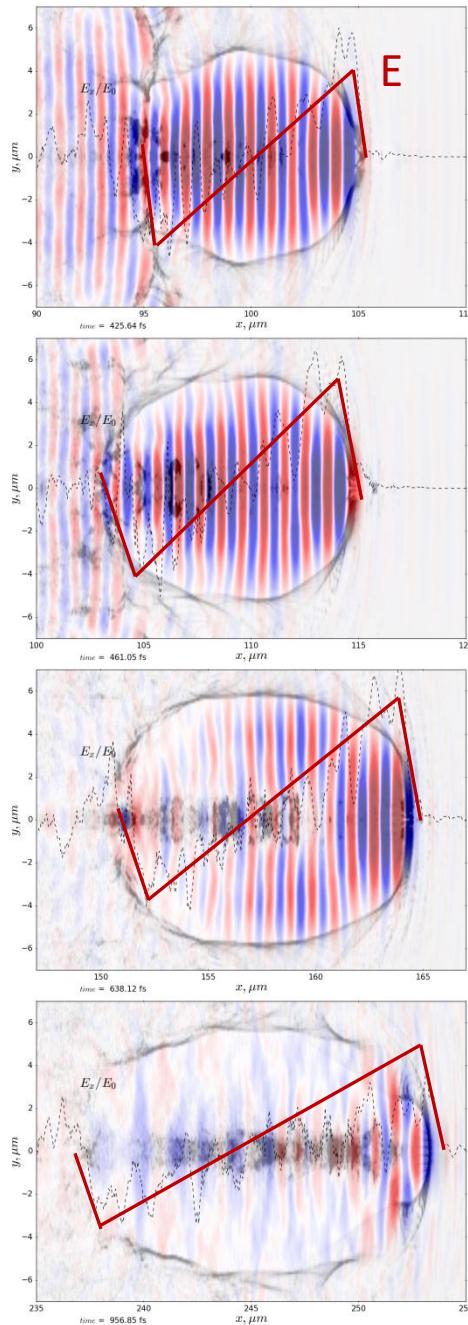
$$c\tau < L \quad c\tau < d$$

Quasi-monoenergetic electrons, pC charge
Up to 10 GeV (gas jet, gas capillary)



Laser bullet

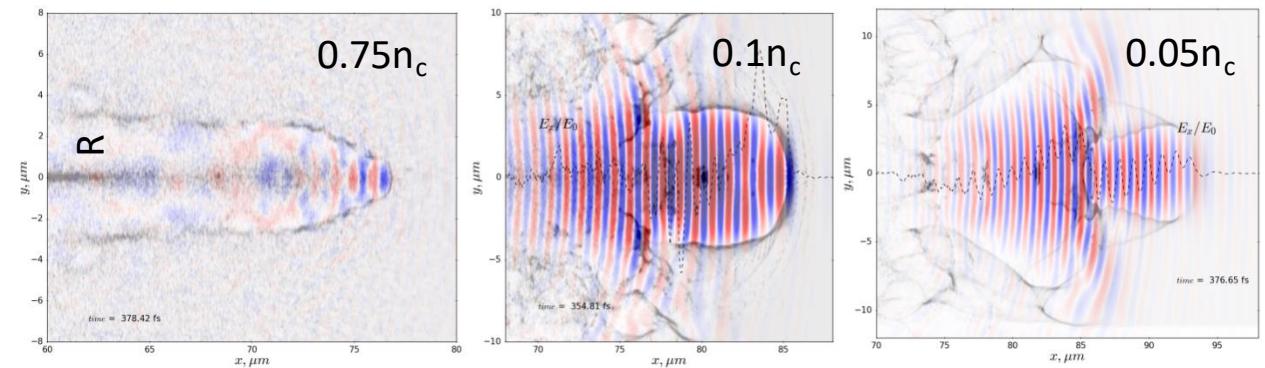
Laser bullet acceleration regime



$$a_0 = 0.85 \lambda [\mu\text{m}] \sqrt{I [\text{W/cm}^2]} \times 10^{-9}$$

$$R \cong \frac{c}{\omega_p} \alpha \sqrt{a_0}$$

α - numeric factor $\sim (1-2)$



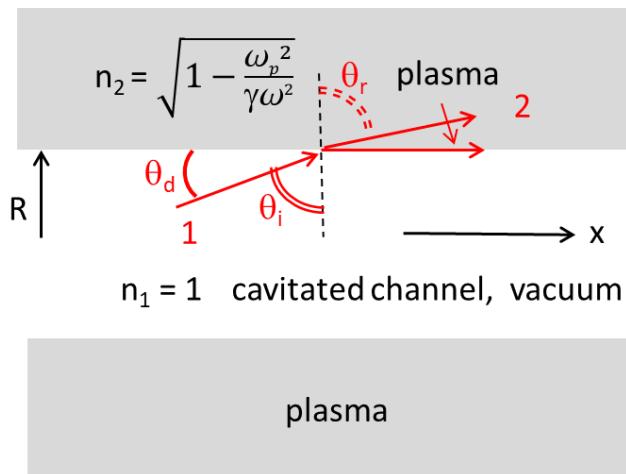
R (n, I) - matching

The matched cavern spot size condition

$$R \cong \frac{c}{\omega_p} \alpha \sqrt{a_0}$$

$$R = \frac{c}{\omega} \sqrt{\frac{n_c}{n_e}} \left(\frac{16\alpha^4 P}{P_c} \right)^{1/6} \quad P - \text{laser pulse power}$$

- 1 Gordienko S and Pukhov A 2005 *Phys. Plasmas* **12**, 043109 $\alpha \approx 1.12$
- 2 Lu M *et al* 2007 *Phys. Rev. STAB* **10**, 061301 $\alpha \approx 2$
- 3 Lobok M G Brantov A V Gozhev D A and Bychenkov V Yu 2018 *Plasma Phys. Control. Fusion* **60** 084010 $\alpha \approx 1.22$



$$\theta_d \simeq \lambda / \pi R \quad \theta_i = \pi/2 - \theta_d$$

$$n_1 \sin \theta_i = n_2 \sin \theta_r$$

M.G.Lobok, A.V.Brantov,
and V.Yu.Bychenkov,
Phys. Plasmas
26, 123107 (2019)

condition of the total internal reflection, $\theta_r = \pi/2$

$$\theta_d^2 \simeq \left(\frac{2c}{\omega R} \right)^2 \simeq \frac{\omega_p^2}{\gamma \omega^2} \simeq \frac{\sqrt{2}\omega_p^2}{a_0 \omega^2} \quad \gamma = \sqrt{1 + a_0^2/2} \simeq a_0/\sqrt{2}$$

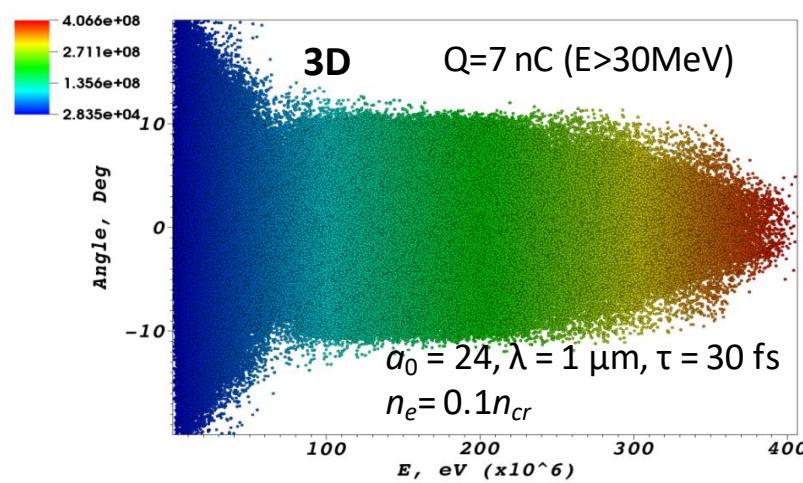
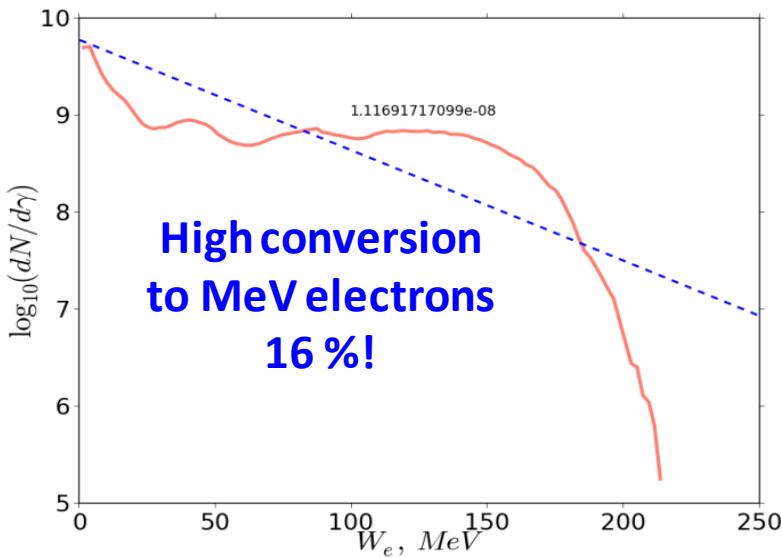


$$R \cong \frac{c}{\omega_p} 2^{3/4} \sqrt{a_0}$$

Theory:

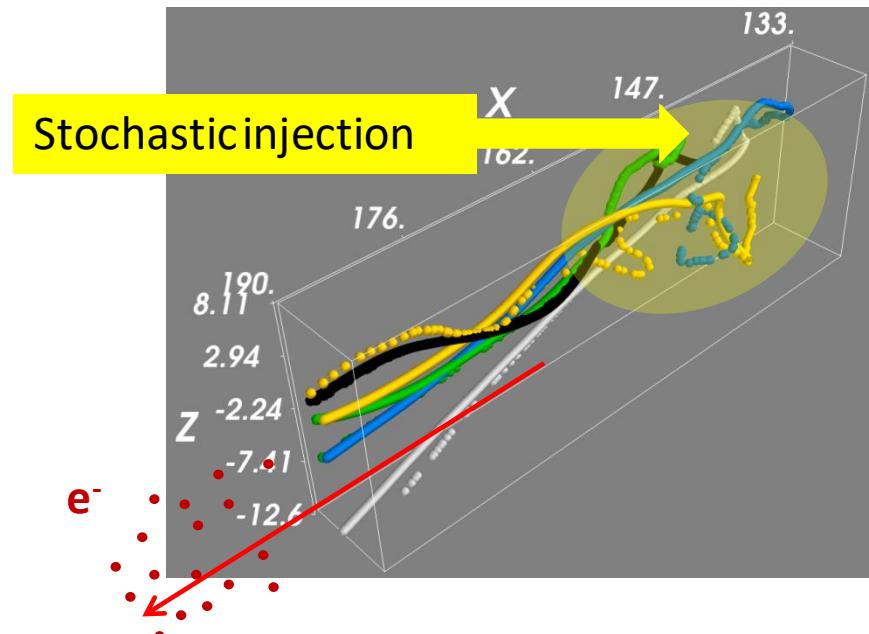
- V.F.Kovalev and V.Yu.Bychenkov, *Phys. Rev. E* **99**, 043201 (2019);
V.Yu.Bychenkov *et al.*, *Plasma Phys. Contr. Fus.* **61**, 124004 (2019);
V.Yu.Bychenkov and V.F.Kovalev, *Radiophysics & Quantum Electronics* **63**, N9 (2020)

Electron acceleration. Stochastic injection.



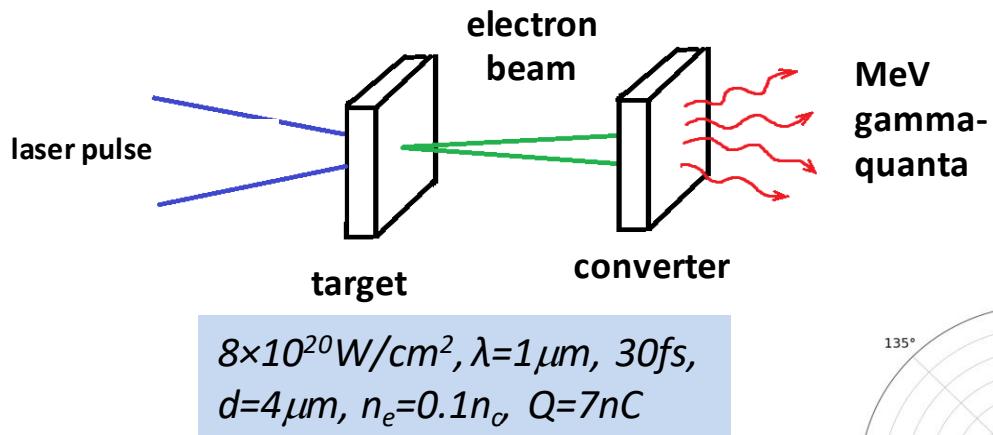
optimum target thickness:

$$l_0 = v_g c \tau_L / (c - v_g) \approx c^2 \tau_L / (c - v_g)$$
$$l_{opt} \simeq l_0 \simeq l_d$$



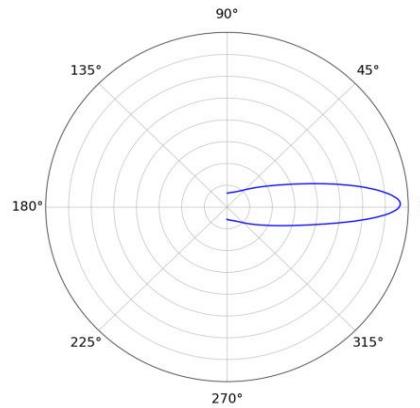
M.G.Lobok, A.V.Brantov, D.A.Gozhev,
and V.Yu.Bychenkov
Plasma Phys. Contr. Fus. 60, 084010 2018

Electron conversion to gamma emission



polar angular diagram
for high-energy γ

Pt of 6 mm thickness

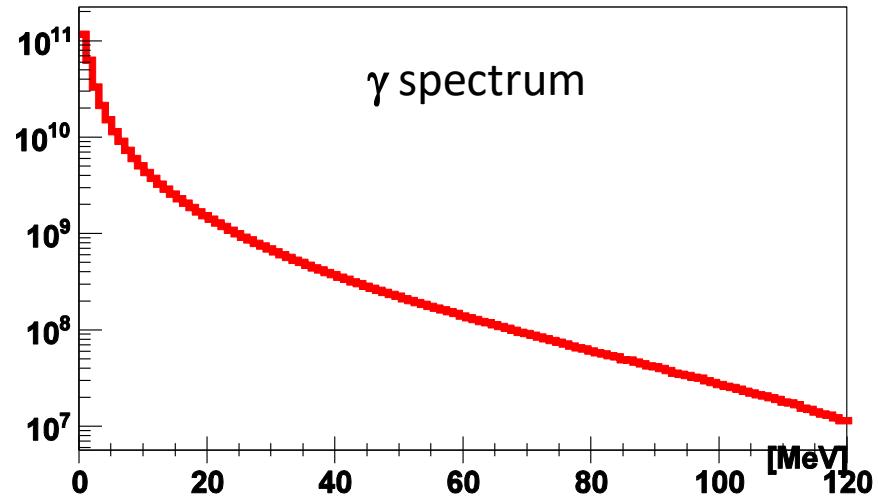
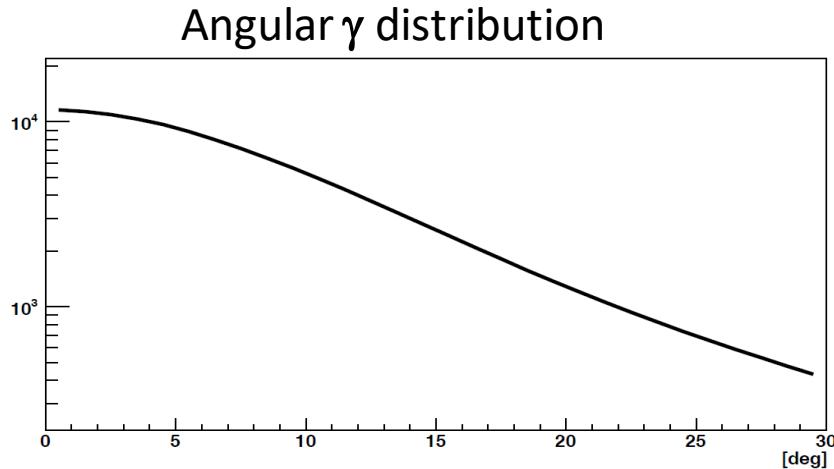


Deep γ -radiography
(10-20 cm dense matter)
Single shot radiography:
multi-tens of nQ

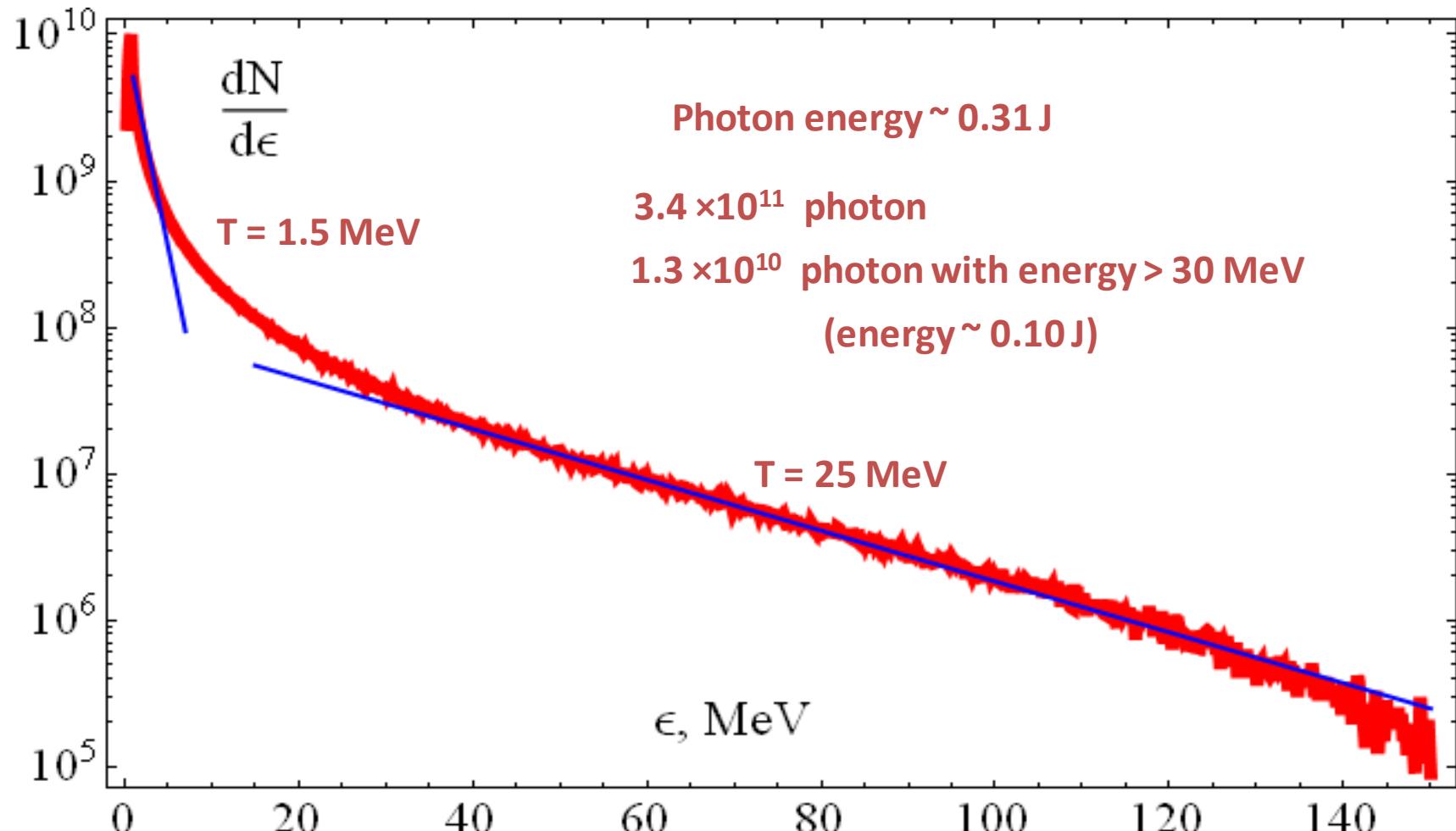


PW lasers

M.G.Lobok, A.V.Brantov,
and V.Yu.Bychenkov,
Phys. Plasmas
26, 123107 (2019)
gammas, pairs, neutrons, pions

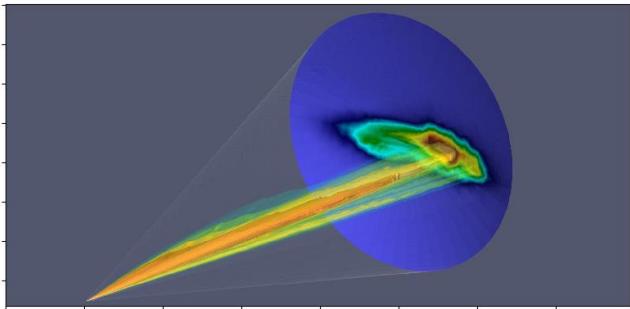


X-ray generation



Conversion from laser to gamma-rays of $\sim 7 \%$

Betatron emission

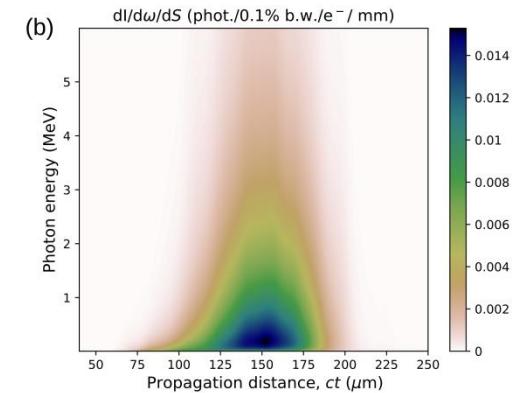
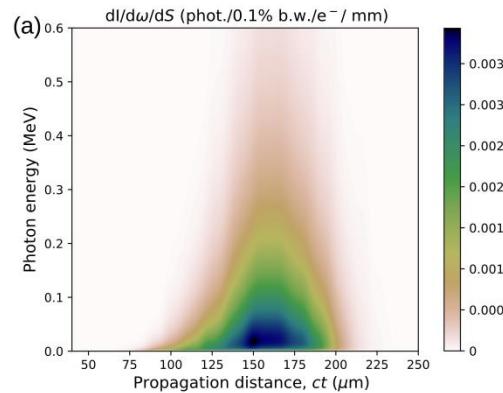


$135 \text{ TW} \rightarrow 3 \times 10^{10} \text{ ph} > 10 \text{ keV}$

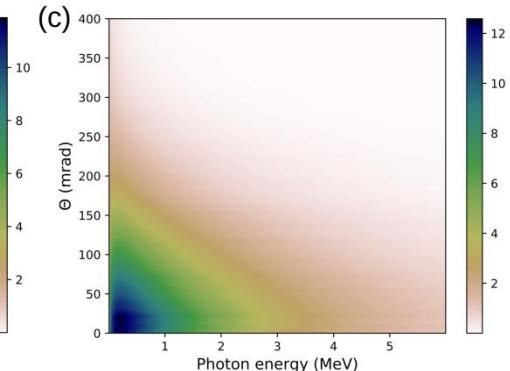
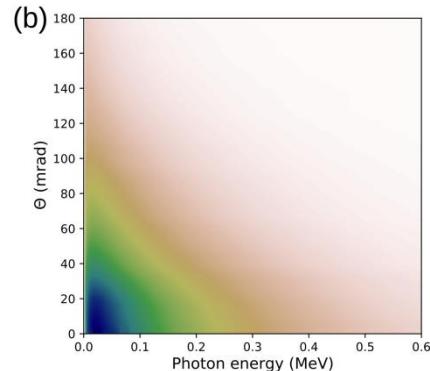
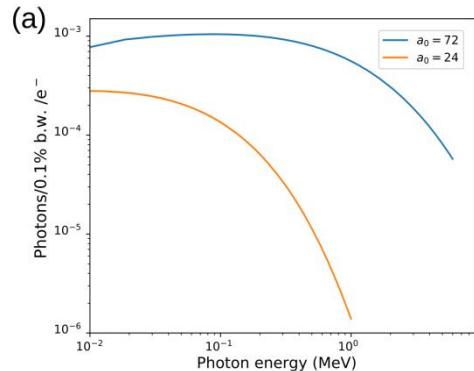
$1.2 \text{ PW} \rightarrow 10^{11} \text{ ph} > 10 \text{ keV}$

The laser-to-gammas energy conversion efficiency $\sim 10^{-4}$ for $\sim 100 \text{ keV}$ ph and few tenths % for entire keV-range.

Angular-integrated spectra of photons emitted per single electron along the laser propagation through the target: $a_0 = 24$ (a) and $a_0 = 72$ (b).



Angular-integrated photon spectra in 0.1% BW (a). The θ -distribution of the betatron radiation spectra for $a_0 = 24$ (b) and $a_0 = 72$ (c).



THz surface pulse generation by escaping electrons

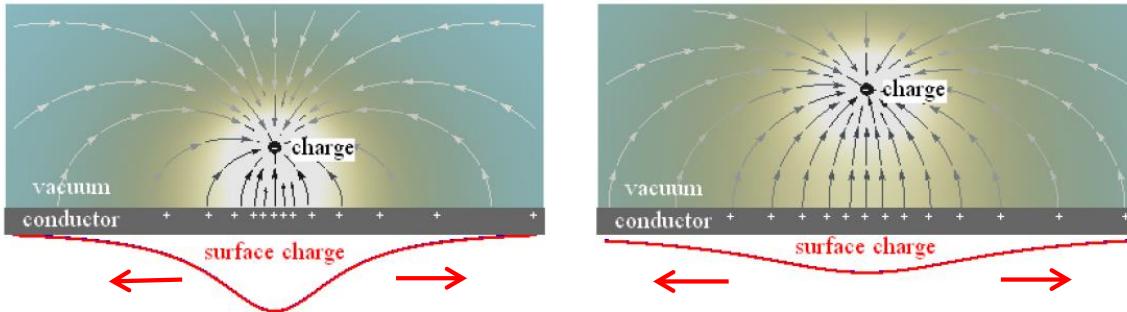


FIG. 1. Cartoon of electric field distribution and surface charge (shown by red line at the bottom) for electron line (position shown by black circle) moving with $v = 0.1c$ in vacuum from conductor boundary at two instants.

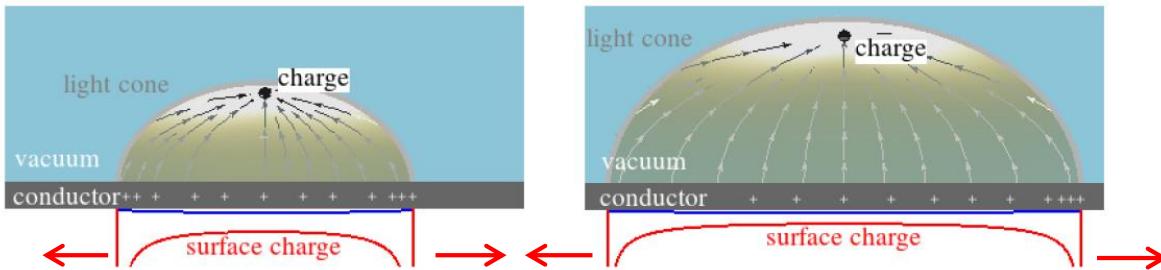
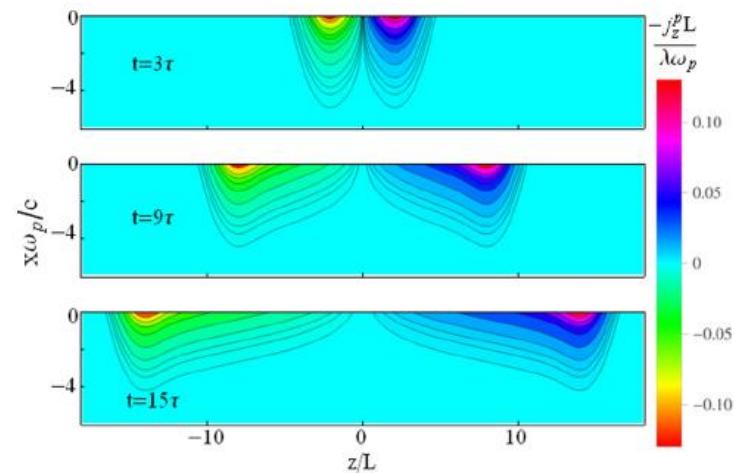


FIG. 2. Cartoon of electric field distribution and surface charge (shown by red line at the bottom) for electron line (position shown by black circle) moving with $v = 0.9c$ in vacuum from conductor boundary at two instants.

A. V. Brantov, A. S. Kuratov, Yu. M. Aliev,
and V. Yu. Bychenkov, Ultrafast target
charging due to polarization triggered by
laser-accelerated electrons,
PHYS. REV. E **102**, 021202(R) (2020)

theory, FDTP numerical
modeling, PIC simulation

Generation and propagation of a
transient electromagnetic pulse and a
lateral current in the wave form along
the target surface at the speed of light



Polarization triggered current

FDTD numerical simulations

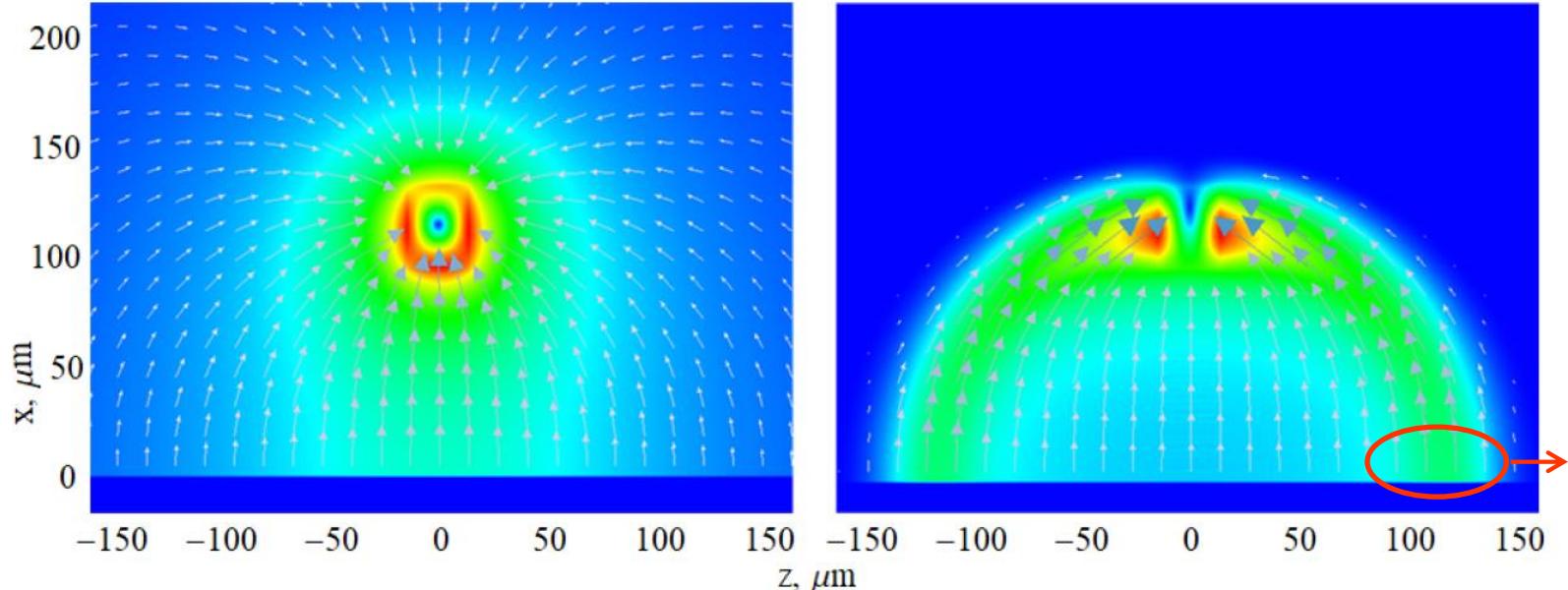
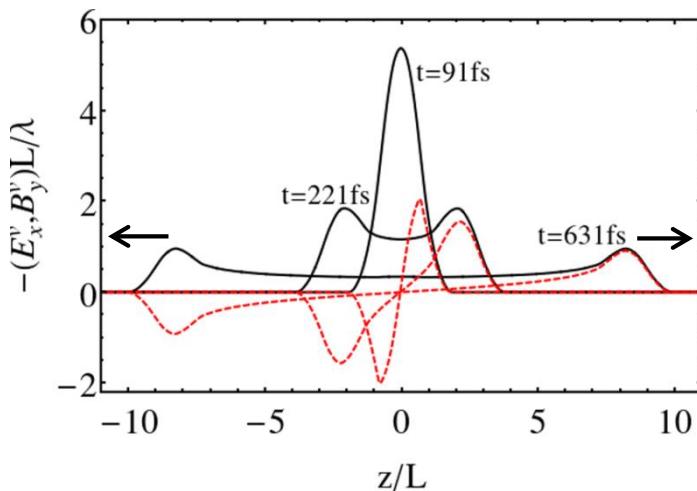


FIG. 3. Electric field distribution for electron bunch moving with $v = 0.2c$ at instant of 2.44 ps (left panel) and $v = 0.9c$ at instant of 551 fs (right panel) from FDTD modeling for $L = c\tau = 20 \mu\text{m}$.



Up to TV/m fields

FIG. 4. Evolution of the vacuum electric field $E_{x0}^v(t, z)$ (black curves) and magnetic field $B_{y0}^v(t, z)$ (red dashed curves) at the target-vacuum boundary from FDTD modeling for $L = c\tau = 20 \mu\text{m}$.

PIC simulations of surface wave

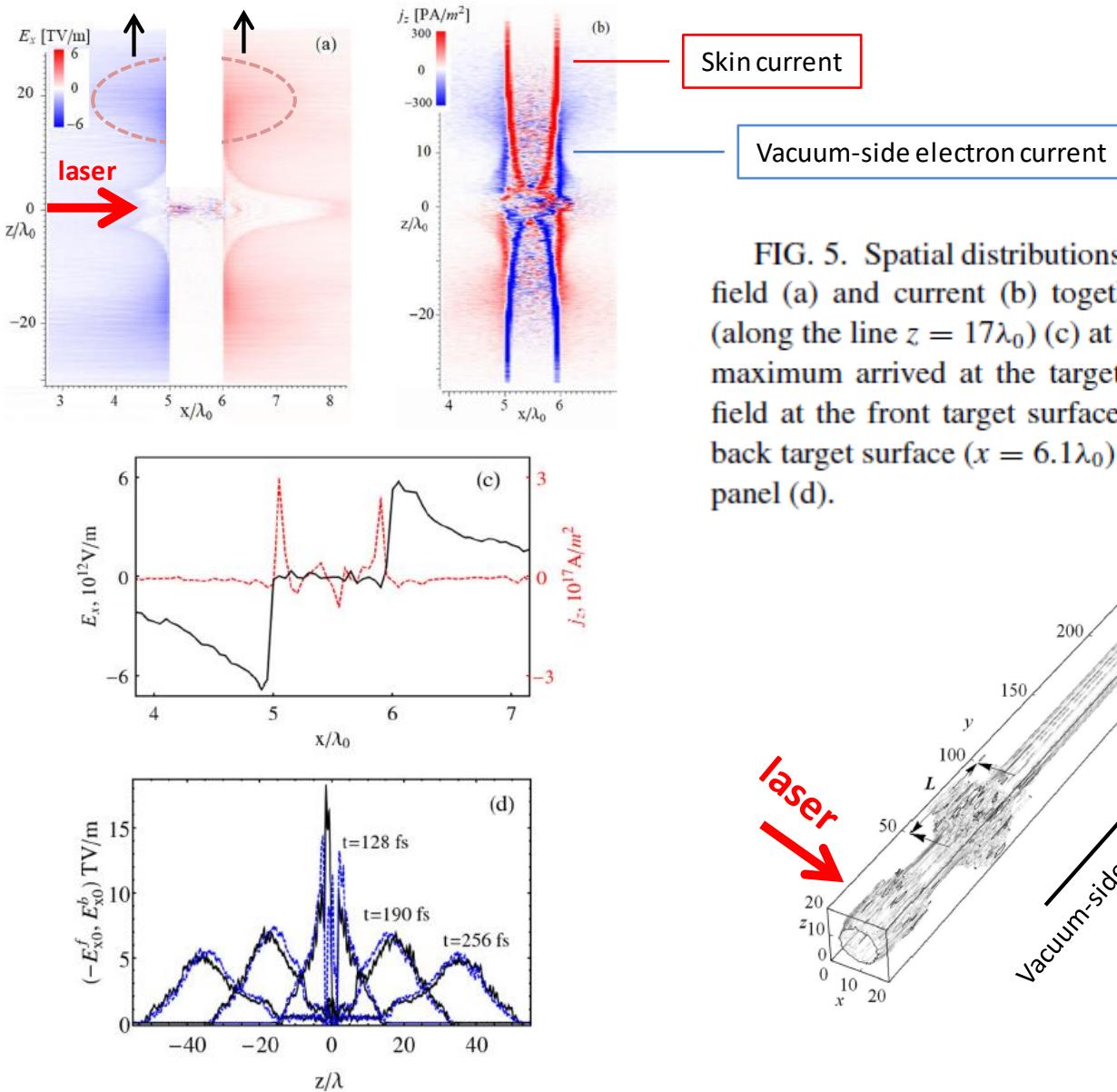
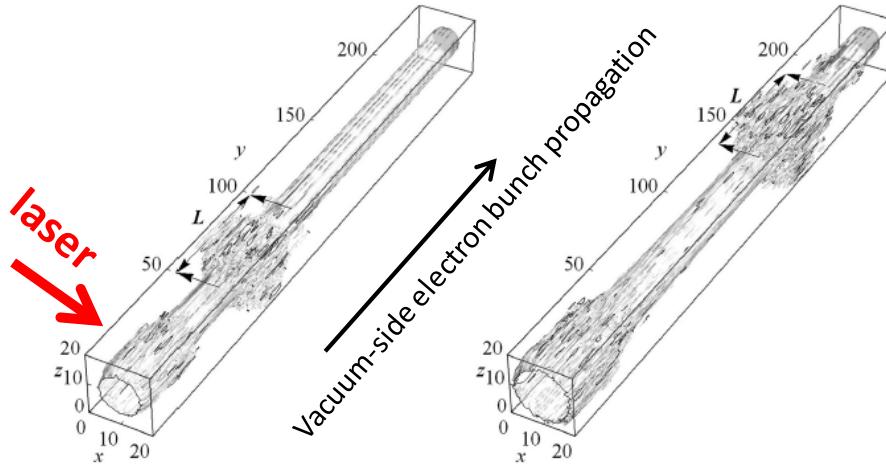
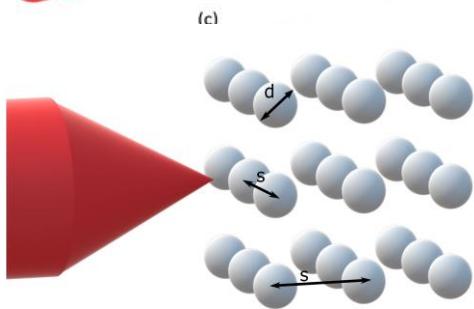
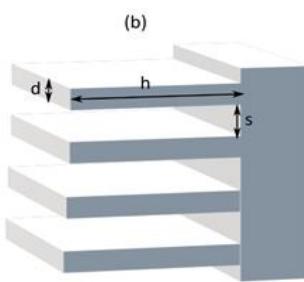
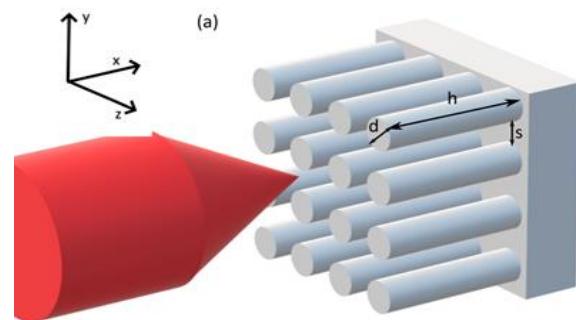


FIG. 5. Spatial distributions (from PIC simulation) of the electric field (a) and current (b) together with their cross-section profiles (along the line $z = 17\lambda_0$) (c) at $t = 190\text{ fs}$ (80 fs after the laser pulse maximum arrived at the target). Evolution of the vacuum electric field at the front target surface ($x = 4.9\lambda_0$) E_{x0}^f (black curves) and back target surface ($x = 6.1\lambda_0$) E_{y0}^b (blue dashed curves) is shown in panel (d).



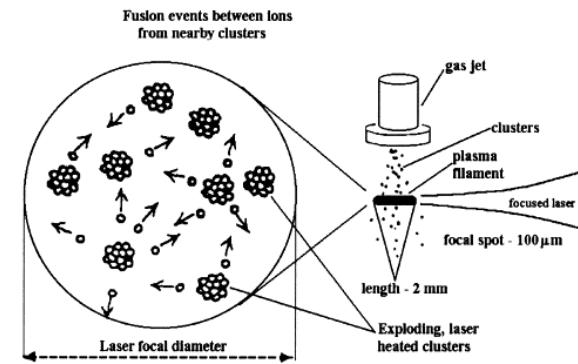
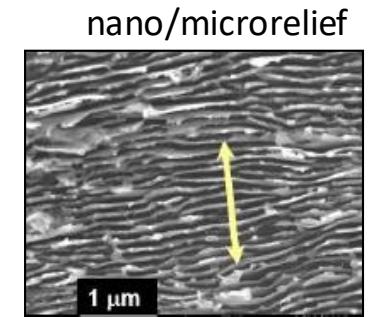
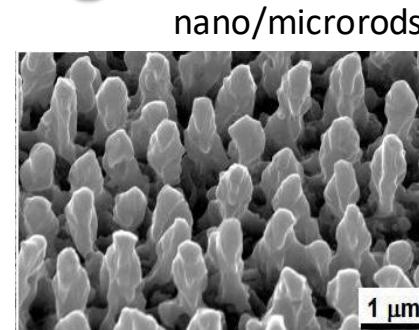
Laser triggered stochastic volumetric heating of micro-structured targets



super-ponderomotive
electrons

Semi-transparent
High average density
Volumetric heating

3D PIC simulation optimization
of laser-plasma interaction
through geometrical matching
(d , h , s)



- nuclear reactions on a table, neutron source
- secondary e.m. emission
- high energy density research

Stochastic electron dynamics in complex field

5 mJ, 60 fs laser of $2 \times 10^{18} \text{ W/cm}^2$

D. A. Gozhev *et al.*, HEDP **37**, 100856 (2020)

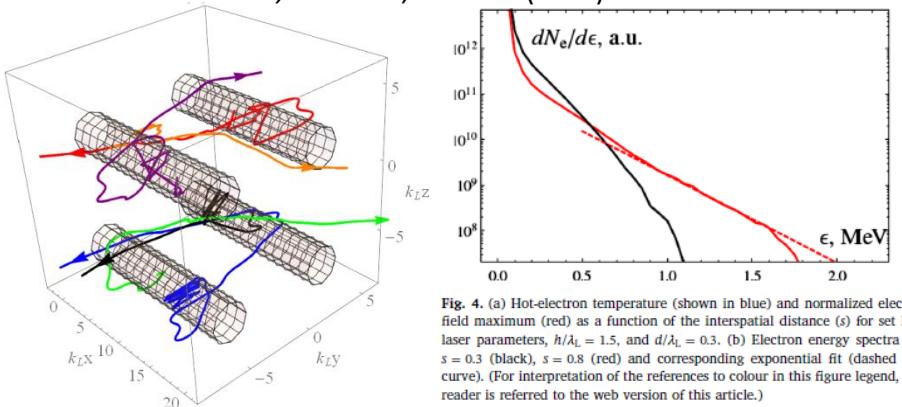


Fig. 4. (a) Hot-electron temperature (shown in blue) and normalized electric field maximum (red) as a function of the interspatial distance (s) for set I of laser parameters, $h/\lambda_L = 1.5$, and $d/\lambda_L = 0.3$. (b) Electron energy spectra for $s = 0.3$ (black), $s = 0.8$ (red) and corresponding exponential fit (dashed red curve). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

NEUTRON PRODUCTION FROM STRUCTURED TARGETS
IRRADIATED BY AN ULTRASHORT LASER PULSE,
S. G. Bochkarev *et al.*, JRLR **42**, N4 (2021)

25 mJ, 100 Hz laser, (1) Ti + 40% D or (2) 20% D and 20% T
1 : $\sim 5 \times 10^7 \text{ n}^\circ/\text{s}$ (DD); 2 : $\sim 10^9 \text{ n}^\circ/\text{s}$ (DT)

Electron acceleration from Au submicron clusters

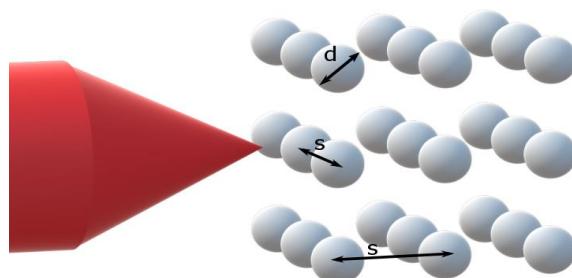
Laser pulse:

Linear polarized, Plane wave

$\lambda_L = 1 \mu\text{m}$

$\tau_{FWHM} = 10 \text{ fs}$

$$I_L = 2 \times 10^{18} - 3.4 \times 10^{19} \frac{\text{W}}{\text{cm}^2} \Rightarrow \\ a_L = 0.85 \sqrt{I \lambda_L^2 / I_{18} \lambda_1^2} \approx 1.2 - 5$$



Clusters electron density: $200 n_{cr}$
Diameter of clusters (d): $0.2\lambda_L$
Average distance between
cluster centers (s): $1.2\lambda_L$
Average electron densiry: $0.48 n_{cr}$

Simulation setup:

Simulation box: $4.5\lambda_L \times 4.5\lambda_L \times 4.5\lambda_L$
Time step: 0.0047 fs
Spatial step : $\frac{\lambda_L}{600} \times \frac{\lambda_L}{200} \times \frac{\lambda_L}{200}$
Particles in cells: 8
Duration of simulation: 60 fs

Boundary conditions:

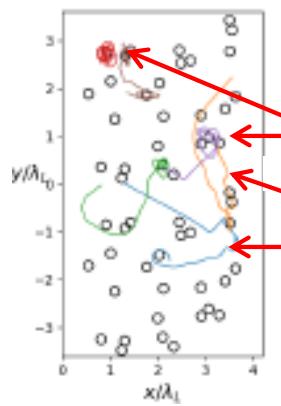
longitudinal - absorbing
transverse - periodic

Electron distribution in a cluster plasma

**PIC
simulations**

$I = 2 \times 10^{18} \text{ W/cm}^2$,

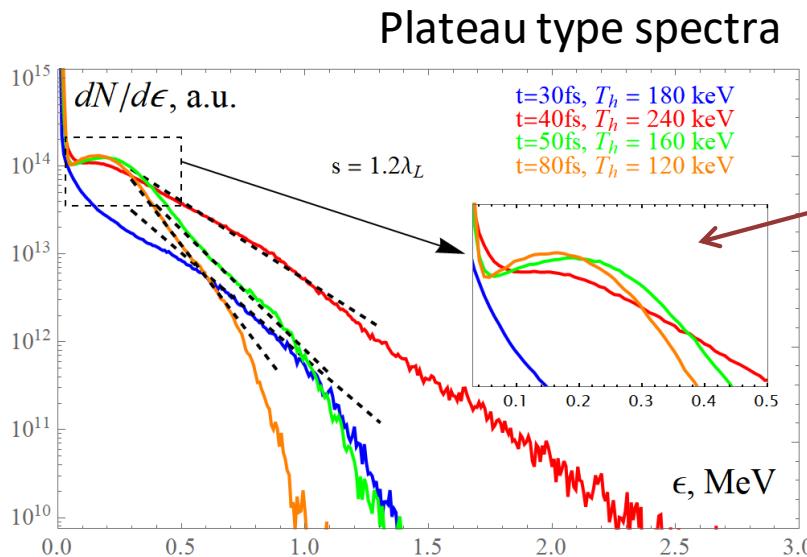
$\tau = 30 \text{ fs}$



Large angle scattering

Captured electrons, Kepler's orbits

Scattering with hopping

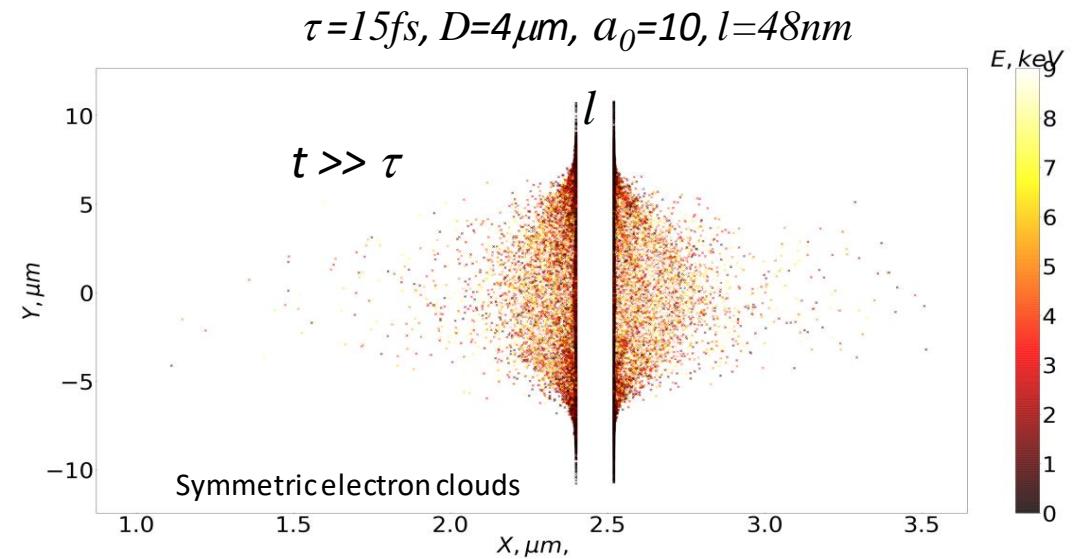
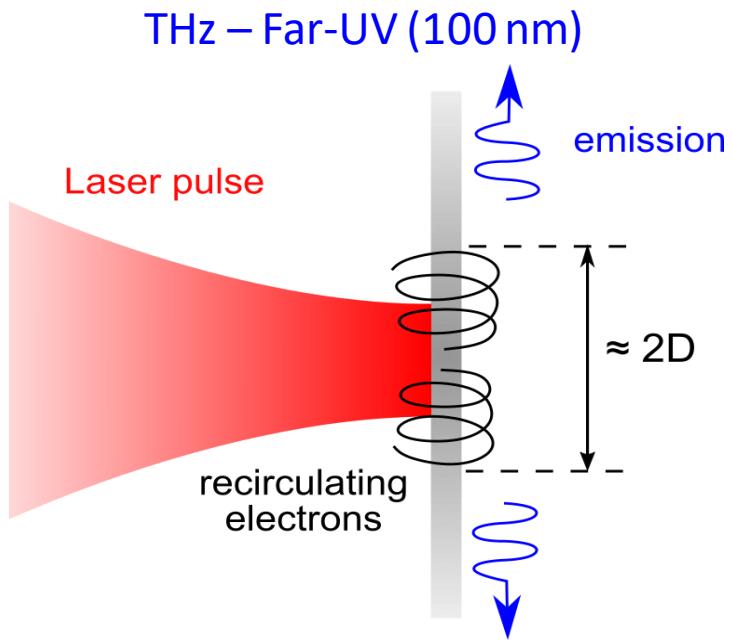


Plateau type spectra

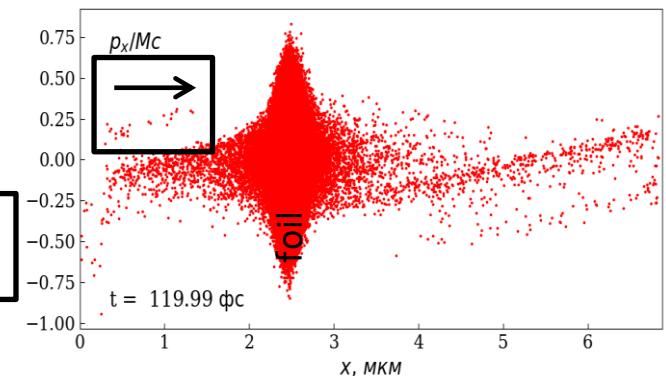
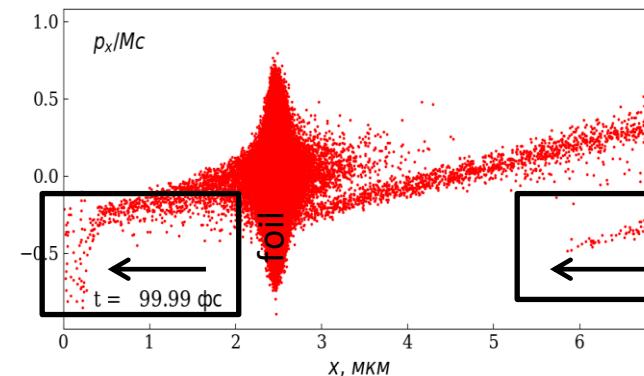
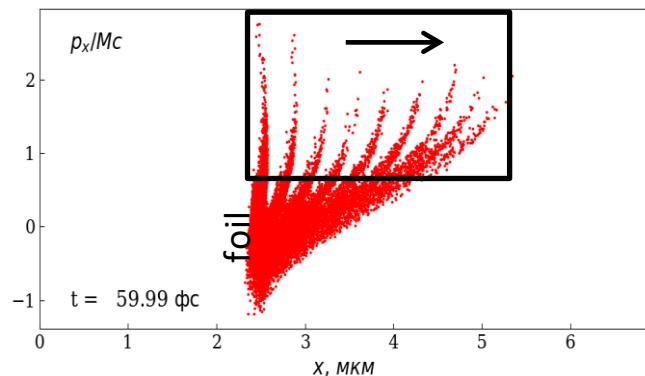
Increased number
of high-energy
electrons

X-ray cluster source

Recirculating electrons from ultrathin foil

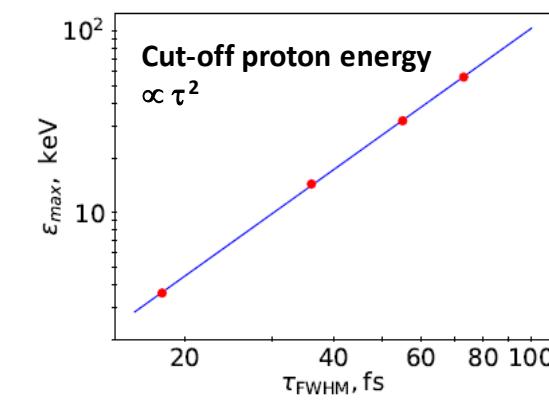
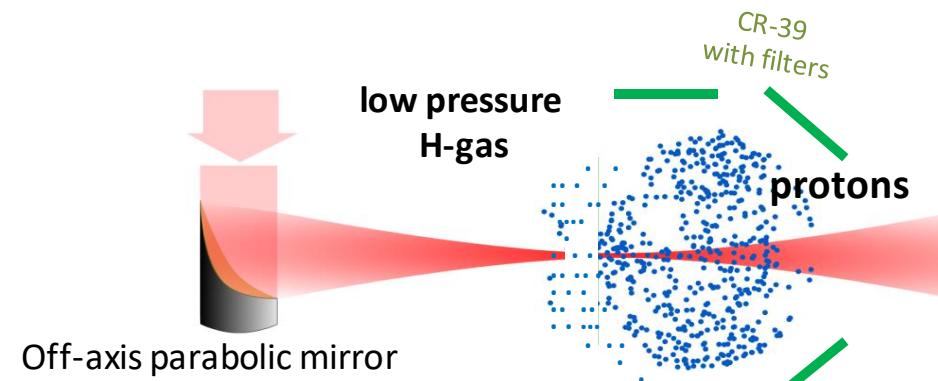


phase plane

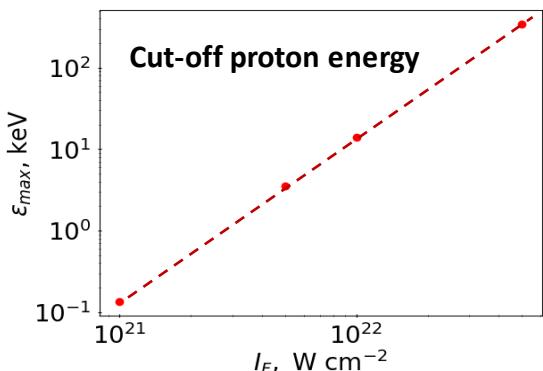


Laser proton acceleration for laser intensity measurement

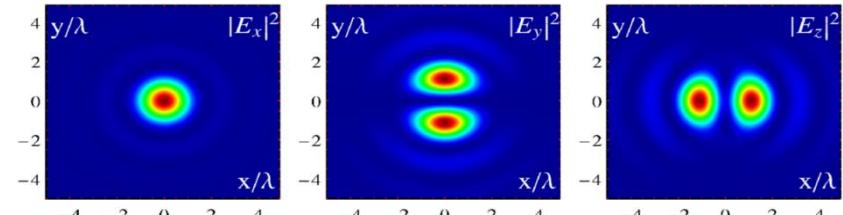
In collaboration with CUOS UoM, USA (A.M. Maksimchuk, A.G.R. Thomas, K. Krushelnick)



Cut-off energy vs pulse duration.
 $I = 10^{22} \text{ W/cm}^2$,
 $D_F = 1.32\lambda$.



- 1) O.E.Vais et al.,
New J. Phys.
22, 023003 (2020)
- 2) O.E.Vais et al.,
PPCF
63, 014002 (2021)



Focusing in diffraction limit

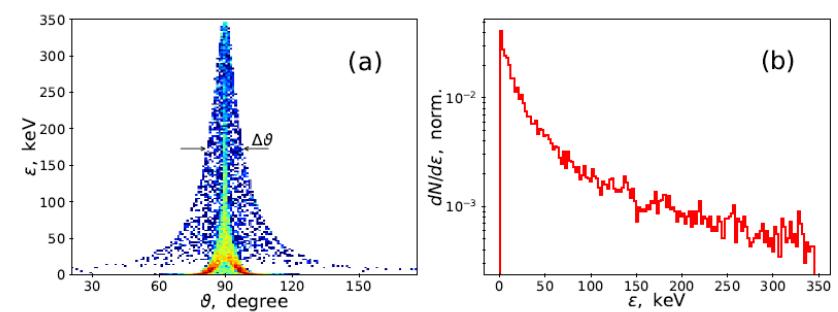
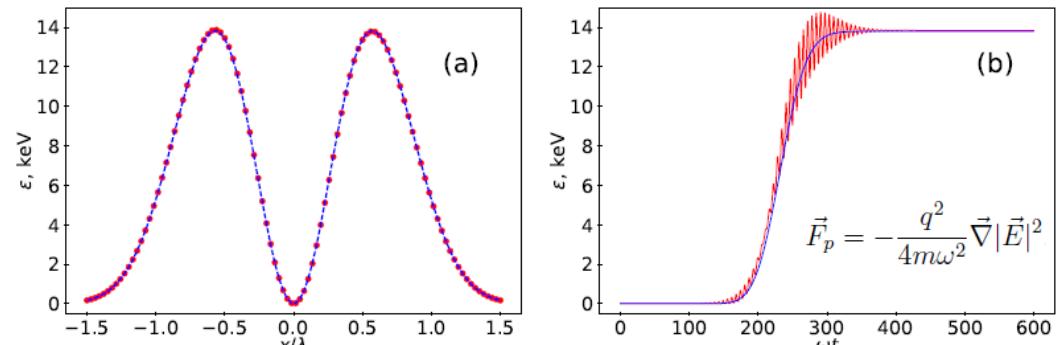


Fig. 3: φ -integrated (a) and angular-integrated (b) spectral distributions of protons accelerated by Gaussian laser pulse with duration $\tau_{FWHM} = 36\text{fs}$ and peak intensity $I_p = 5 \times 10^{22} \text{ W cm}^{-2}$, focused by off-axis parabolic mirror ($\rho = 5.08\text{cm}$, $\psi_{off} = 90^\circ$) into the focal spot $D_F = 1.3\lambda$.

*Thank
you*