# Systematic Uncertainties in measuring the Weinberg angle at SCT

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Workshop on future charm-tau factory

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#### The weak mixing angle

• Electroweak model  $SU(2)_L \times U(1)_Y$  (Glashow, 1961)

$$A_{\mu} = B_{\mu}^{0} \cos \theta_{W} + W_{\mu}^{0} \sin \theta_{W}$$

$$Z_{\mu} = W_{\mu}^{0} \cos \theta_{W} - B_{\mu}^{0} \sin \theta_{W}$$

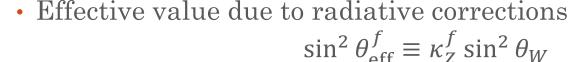
Two independent coupling constants g and g'

• On-shell definition of the weak mixing angle

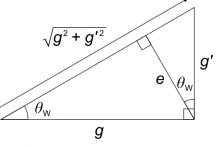
$$\sin^2 \theta_W \equiv \frac{{g'}^2}{g^2 + {g'}^2} = 1 - \frac{m_W^2}{m_Z^2}$$
  $z \sim \sqrt{y/z}$ 

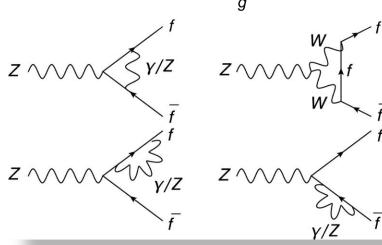


$$\frac{g}{\cos\theta_W} Z_\mu \bar{f} \gamma^\mu \left( I_3^f - 2Q_f \sin^2\theta_W - I_3^f \gamma_5 \right) f, \qquad I_3^f = 0, \pm 1/2$$



Full two-loop EW fermionic and bosonic corrections completed recently

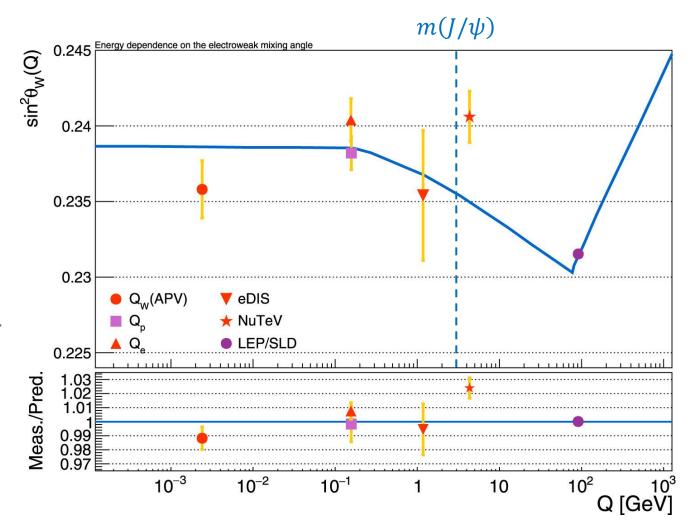




$$I_3^f = 0, \pm 1/2$$

### $\sin^2 \theta_{\rm eff}$ measurements

- $A_{FB}$  close to the Z pole
  - $\delta(\sin^2\theta_{\rm eff}) \approx 0.0006$
  - $Q = m_Z = 91 \text{ GeV}$
- Atomic parity violation
  - $\delta(\sin^2\theta_{\rm eff}) \approx 0.002$
  - $Q \sim 10^{-3} \text{ GeV}$
- ν and polarized e<sup>-</sup> scattering on fixed targets
  - $\delta(\sin^2\theta_{\rm eff}) \approx 0.0012 0.004$
  - $Q \sim 0.1 1 \, \text{GeV}$
- Planned experiments
  - P2 at MESA (Mainz)
  - Moller at JLab



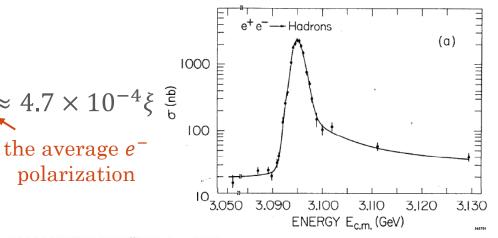
#### Left-right asymmetry at $J/\psi$

• Interference of  $\gamma^*$  and  $Z^*$  annihilation

$$A_{LR} \equiv \frac{\sigma_{+} - \sigma_{-}}{\sigma_{+} + \sigma_{-}} = \frac{3/8 - \sin^{2}\theta_{eff}^{c}}{2\sin^{2}\theta_{eff}^{c} \left(1 - \sin^{2}\theta_{eff}^{c}\right)} \left(\frac{m_{J/\psi}}{m_{Z}}\right)^{2} \xi \approx 4.7 \times 10^{-4} \xi^{\frac{2}{6}}$$

- Parameters of the SCT experiment
  - Luminosity  $L = 10^{35} \text{ cm}^{-2} \text{s}^{-1}$
  - Cross-section  $\sigma(e^+e^- \to J/\psi) \approx 3 \times 10^{-30} \text{ cm}^2$
  - One data-taking season  $T_{tot} = 10^7 \text{ s}$
  - Efficiency of  $J/\psi$  decays used in the analysis  $\varepsilon \approx 0.9$

$$dA_{LR} \approx \frac{1}{\sqrt{L\sigma T tot \varepsilon}} \sim 10^{-6}$$



#### СЛАБЫЕ НЕЙТРАЛЬНЫЕ ТОКИ НОВЫХ КВАРКОВ В $e^+e^-$ -АННИГИЛЯЦИИ

ю. и. сковпень, и. б. хриплович

ИНСТИТУТ ЯДЕРНОЙ ФИЗИКИ СО АН СССР (Поступило в редакцию 11 апреля 1979 г.)

и при полной продольной поляризации обеих или хотя бы одной из начальных частиц находим для qar q-резонанса

$$\eta(1,-1) = \eta(1,0) = \eta(0,-1) = \frac{\sqrt{2}Gm^2}{8\pi\alpha|Q|} (1-4|Q|\sin^2\theta). \tag{6}$$

При  $\sin^2\theta = ^1/_4$  <sup>1)</sup> эта величина составляет соответственно для  $\psi$ - и  $\Upsilon$ -пиков

$$\eta_{\psi} = \frac{\sqrt{2} G m^2}{16\pi \alpha} \approx 4.10^{-4}, \quad \eta_{\rm r} = \frac{\sqrt{2} G m^2}{4\pi \alpha} \approx 1.6.10^{-2}.$$
(7)

## $\sin^2(\theta_{\rm eff}^c)$ at $J/\psi$

$$\frac{d\sin^2\theta_{\rm eff}^c}{\sin^2\theta_{\rm eff}^c} \approx -0.44 \frac{dA_{LR}}{A_{LR}} \oplus 0.44 \frac{d\xi}{\xi} \approx 0.1\%$$

• Ultimate one-year precision

$$\delta(\sin^2\theta_{\rm eff}^c) \approx 2.5 \times 10^{-4}$$

- The average electron beam polarization  $\xi$  should be controlled with precision better than  $10^{-3}$ 
  - On-line laser diagnostics
  - Off-line data-driven approach:

$$e^+e^- \to J/\psi \to [\Lambda \to p\pi^-][\overline{\Lambda} \to \overline{p}\pi^+]$$

#### JHEP 2020, 76 (2020)

A. Bondar, A. Grabovsky, A. Reznichenko, A. Rudenko & V. Vorobyev

#### Subtleties and difficulties

- 1. Due to natural radiative polarization the averaged beam polarization  $\xi_+ \neq -\xi_-$
- 2. Beam life time depends on square of polarization degree
- 3. J/Ψ cross-section depends on beam energy spread. Due to intrabeam scattering the energy spread depends on square of polarization degree
- 4. Not equal average positive and negative beam polarization  $\xi_{+} \neq -\xi_{-}$  may effect on asymmetry measurement
- 5. Effect of natural polarization of positrons

#### How to overcome above difficulties?

#### Experiment at SCT

- 1. Set beam energy at  $\sqrt{s} \approx m(J/\psi)$ , about 300 bunches circulate simultaneously
- 2. Set random polarization,  $\xi_+$  or  $\xi_-$ ,  $\xi_+ \approx -\xi_-$ , for each  $e^-$  bunch
- 3. During the data taken flip spin of all bunches with high as possible frequency
- 4. Count numbers of the  $J/\psi \to \text{hadrons events } N_+$  and  $N_-$  for the polarizations  $\xi_+$  and  $\xi_-$

$$N_{\pm} \sim 10^{12}$$
, event rate  $\approx 100 \text{ kHz}$ 

5. Calculate the cross sections and left-right asymmetry

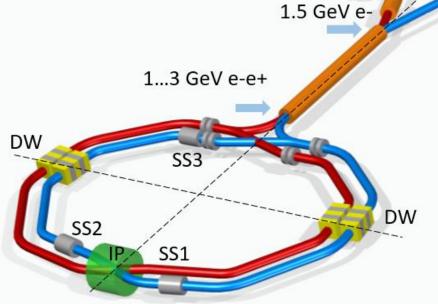
$$\sigma_{\pm} = \frac{N_{\pm}}{\mathcal{L}_{\pm}\varepsilon_{\mathrm{det}}}, \qquad A_{LR} = \frac{\sigma_{+} - \sigma_{-}}{\sigma_{+} + \sigma_{-}}$$

• Luminosity monitoring and backgrounds Statistical precision  $\sigma_{\mathcal{L}}/\mathcal{L} \sim 10^{-6}$  is needed

#### SCT collider

- Beam energy: from 1 to 3.5 GeV
- $\mathcal{L} = 10^{35} \text{ cm}^{-2} \text{s}^{-1} @ 2 \text{ GeV}$
- Longitudinal polarization of the electron beam
- Crab-waist collisions

Perimeter	632.94 m		
$2\theta$	60 mrad		
$eta_x^*/eta_y^*$	100 mm / 1 mm		
$F_{RF}$	350 MHz		
E <sub>beam</sub> (GeV)	1.5	2.5	3.5
<i>I</i> (A)	2	2	2
$N_{ m bunch}$	292	328	262
$L_{\rm peak} \times 10^{35}$ (cm <sup>-2</sup> s <sup>-1</sup> )	0.8	1.0	1.0



1.5 GeV er bypass e+ DR 50 m 100 m **Novosibirsk Super Charm Tau Factory** 

1.5 GeV e-

2.5 GeV e-

1.5 GeV e+

Pol e-

e+ DR - positron damping ring

DW - damping wiggler

SS - Siberian Snake

CV – electron-positron converter

Pol e-/e- - polarized/un-polarized electron

source

see Tomorrow Alexey Petrenko and Anton Bogomyagkov talks

#### Systematic errors suppression

A measurement of spin-dependent effects requires periodic reversal of the polarizations of the electron beams. As a rule this procedure does not effect other parameters of the beam.

$$\delta A_{LR}^{Syst} = \frac{\Delta_t}{\tau_L \sqrt{n_B T_{tot}/\hat{T}}} \sim 10^{-6}$$

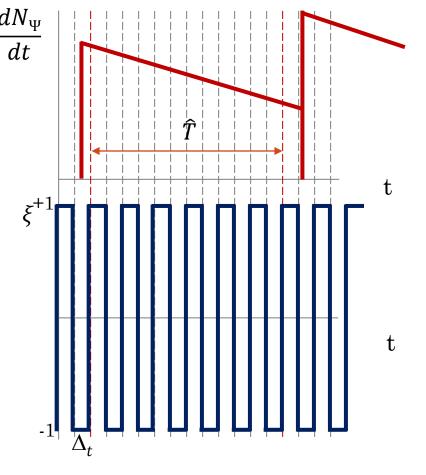
 $\Delta_t$ = 5 sec - time interval between spin-flips

 $\hat{T}$ = 300 sec duration of the single bunch data taking

 $\tau_L$ = 2000 sec - luminosity life time

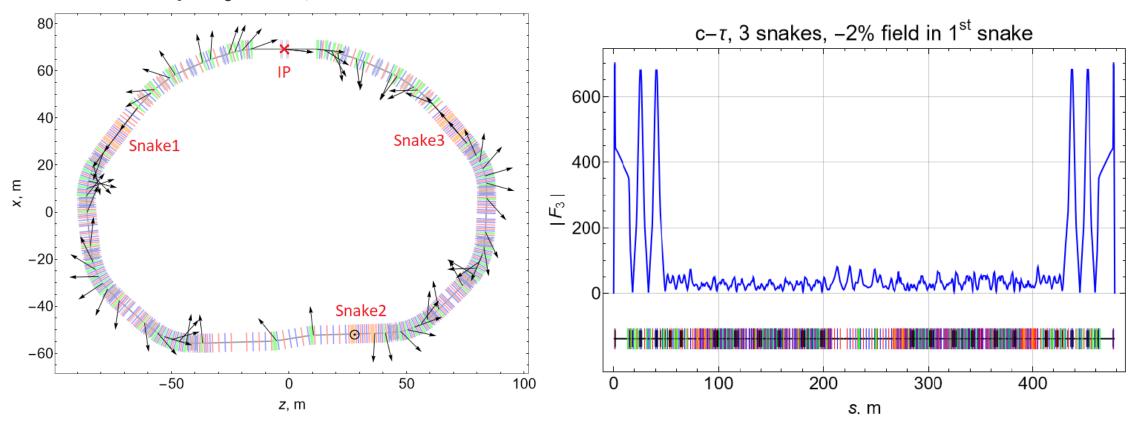
 $n_R = 300$  - number of bunches

 $T_{tot} = 10^7 \text{ sec}$ - time of measurement



#### Spin motion in the storage ring

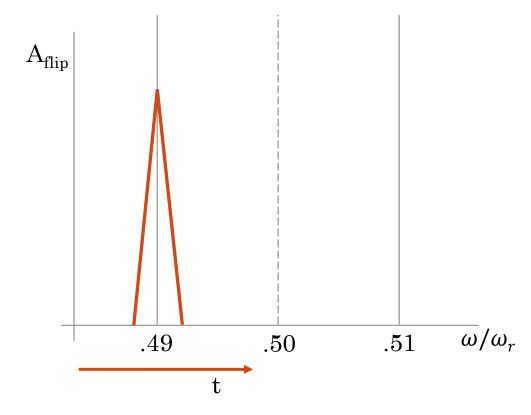
 $\mathbf{n}_0$  along machine,  $E = 1.55 \,\text{GeV}$ 



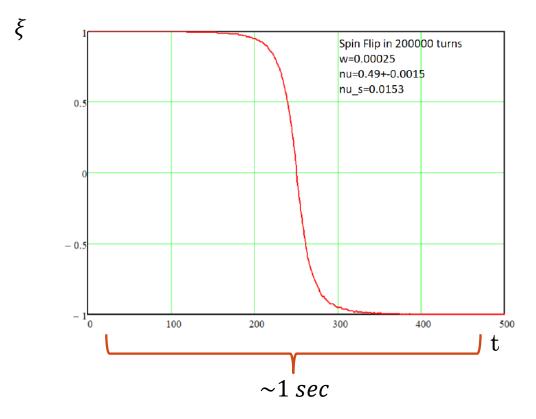
Equilibrium spin direction along the ring. The spin turn is equal to half integer number.

Spin response function to the flipper transverse magnetic field along the ring

#### Spin flip



Spin Flipper RF-field amplitude variation with field frequency for adiabatic resonant spin flip



Result of the adiabatic resonant spin flip simulation

#### Requirements to Detector

- 1. Event time resolution is better than 1 nsec
- 2. High rate (>100 kHz) event record capability
- 3. High rate luminosity measurement at small angles

#### Requirements to Collider

- 1. High luminosity
- 2. Electron beams with high degree of longitudinal polarization
- 3. Fast control bunch by bunch polarization
- 4. Adiabatic spin flip
- 5. Flexible beams injection
- 6. Guaranty depolarization of the positrons ( $\xi_{\parallel}$ <10<sup>-6</sup> at the IP)

#### Forward - Backward asymmetry

• Annihilation process  $e^+e^- \to f\bar{f}$ 

$$\frac{d\sigma}{d\cos\theta} \propto A(1+\cos^2\theta) + B\cos\theta$$

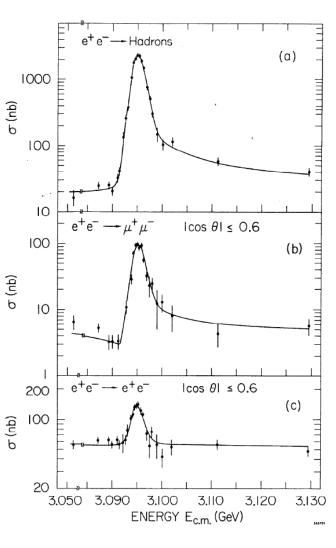
Forward-backward asymmetry

$$A_{FB}^{f} \equiv \frac{(\sigma_{LF} - \sigma_{LB}) - (\sigma_{RF} - \sigma_{RB})}{(\sigma_{LF} + \sigma_{LB}) + (\sigma_{RF} + \sigma_{RB})} = \frac{3}{4} |\xi| A_f,$$

$$A_{f} \equiv \frac{2g_{v}^{f}g_{a}^{f}}{\left(g_{a}^{f}\right)^{2} + \left(g_{v}^{f}\right)^{2}} = \frac{1 - 4|Q_{f}|\sin^{2}\theta_{\text{eff}}^{f}}{1 - 4|Q_{f}|\sin^{2}\theta_{\text{eff}}^{f} + 8|Q_{f}|\sin^{4}\theta_{\text{eff}}^{f}} \left(\frac{m_{J/\psi}}{m_{Z}}\right)^{2}$$

Counting experiment

#### Forward - Backward asymmetry



• Interference of  $\gamma^*$  and  $Z^*$  annihilation

• 
$$A_{FB}^f \equiv \frac{(\sigma_{LF} - \sigma_{LB}) - (\sigma_{RF} - \sigma_{RB})}{(\sigma_{LF} + \sigma_{LB}) + (\sigma_{RF} + \sigma_{RB})} = \frac{3/8 - \sin^2 \theta_{\text{eff}}^c}{2 \sin^2 \theta_{\text{eff}}^c (1 - \sin^2 \theta_{\text{eff}}^c)} \left(\frac{m_{J/\psi}}{m_Z}\right)^2 \xi$$
 (A.Milstein)

Parameters of the SCT experiment

- Luminosity  $L = 10^{35} \text{ cm}^{-2} \text{s}^{-1}$
- Cross-section  $\sigma(e^+e^- \to J/\psi \to \mu^+\mu^-) \approx 0.2 \times 10^{-30} \text{ cm}^2$
- One data-taking season  $T_{tot} = 10^7 \text{ s}$

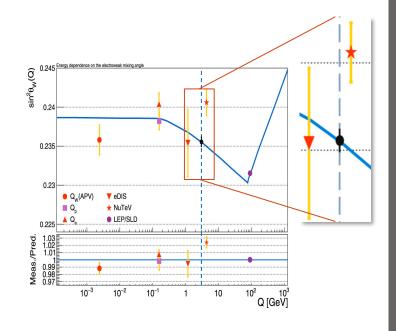
$$dA_{FB}^f \approx \frac{2}{\sqrt{L\sigma Ttot}} \sim 5 \times dA_{LR} = 5 \times 10^{-6}$$

$$\delta(\sin^2\theta_{\rm eff}^c) \approx 1.25 \times 10^{-3}$$

5 times less statistical sensitivity, but much less systematic error

#### Conclusions

- 1. SCT with polarized electron beam is a unique experiment to study neutral weak coupling of the charm quark and to measure  $\sin^2\theta_{\rm eff}^f$
- 2. The decay  $J/\psi \to \Lambda \overline{\Lambda}$  can be used as a precise monitor of the average polarization of electrons
- 3. Spin "gymnastics" radically suppresses systematic uncertainties in measuring asymmetry of  $A_{LR}$
- 4. Measurement of forward backward asymmetry in the  $J/\psi \rightarrow \mu \bar{\mu}$  even more free from possible systematic errors



Thus, it can be conclude that precision measurement of  $sin^2\theta_w$  at Super C-Tau factory could be feasible

# Backup

#### Luminosity monitoring

$$\sigma_{\pm} = \frac{N_{\pm}}{\mathcal{L}_{+}\varepsilon_{\rm eff}}$$

- Statistical precision  $\sigma_{\mathcal{L}}/\mathcal{L} \sim 10^{-6}$  is needed
  - Multiplicative biases vanish in asymmetry
- £ monitoring with Bhabha events

$$\sigma(e^+e^- \to e^+e^-)_{\theta > 10^\circ} \approx 1 \times 10^{-30} \text{ cm}^2 \approx \sigma(e^+e^- \to J/\psi)$$

- Bhabha events statistics will limit precision
- L monitoring with dedicated device at low angle
  - Would provide good support for the  $\sin^2 \theta_{\rm eff}$  measurement
  - The device should be able to measure bunch-by-bunch luminosity

#### $A_{FB}$ at LEP

• Annihilation process  $e^+e^- \to Z \to f\bar{f}$ , unpolarized cross-section

$$\frac{d\sigma}{d\cos\theta} \propto A(1+\cos^2\theta) + B\cos\theta$$

Forward-backward asymmetry

$$A_{FB}^{f} \equiv \frac{\sigma_{F} - \sigma_{B}}{\sigma_{F} + \sigma_{B}} = \frac{3}{4} A_{e} A_{f},$$

$$A_{f} \equiv \frac{2g_{v}^{f} g_{a}^{f}}{\left(g_{a}^{f}\right)^{2} + \left(g_{v}^{f}\right)^{2}} = \frac{1 - 4|Q_{f}|\sin^{2}\theta_{\text{eff}}^{f}}{1 - 4|Q_{f}|\sin^{2}\theta_{\text{eff}}^{f} + 8|Q_{f}|\sin^{4}\theta_{\text{eff}}^{f}}$$

Counting experiment

### $\sin^2 \theta_{\rm eff}$ at colliders

#### 1. LEP

- Unpolarized  $e^+e^-$  beams near the Z pole,  $17 \times 10^6$  Zs
- Forward-backward asymmetry
- 2. SLAC Large Detector (SLD)
  - Polarized  $e^+e^-$  beams near the Z pole,  $50 \times 10^3 Zs$
  - Average beam polarization of 60%
  - Combinations of the forward-backward and left-right asymmetries
- 3. LHC: ATLAS, CMS, LHCb
  - Unpolarized proton beams
  - Tests of the  $Z \to l\bar{l}$  couplings and measurement of the  $\sin^2 \theta_{\rm eff}^l$
  - Model-dependent

#### **SLC Experiment**

Polarized beam gives access to the left-right asymmetry

$$A_{LR} \equiv \frac{\sigma_{+} - \sigma_{-}}{\sigma_{+} + \sigma_{-}} = A_{e}\xi$$

 $\xi$  is the average polarization of the electron beam

Forward-backward asymmetry with polarized beam

$$A_{FB}^{f} = \frac{3}{4} A_f \frac{A_e + \xi}{1 + A_e \xi}$$

• Left-right forward-backward cross-section ratio

$$A_f = \frac{4}{3} \frac{\sigma_{LF}^f + \sigma_{RB}^f - \sigma_{LB}^f - \sigma_{RF}^f}{\sigma_{LF}^f + \sigma_{RB}^f + \sigma_{LB}^f + \sigma_{RF}^f}$$

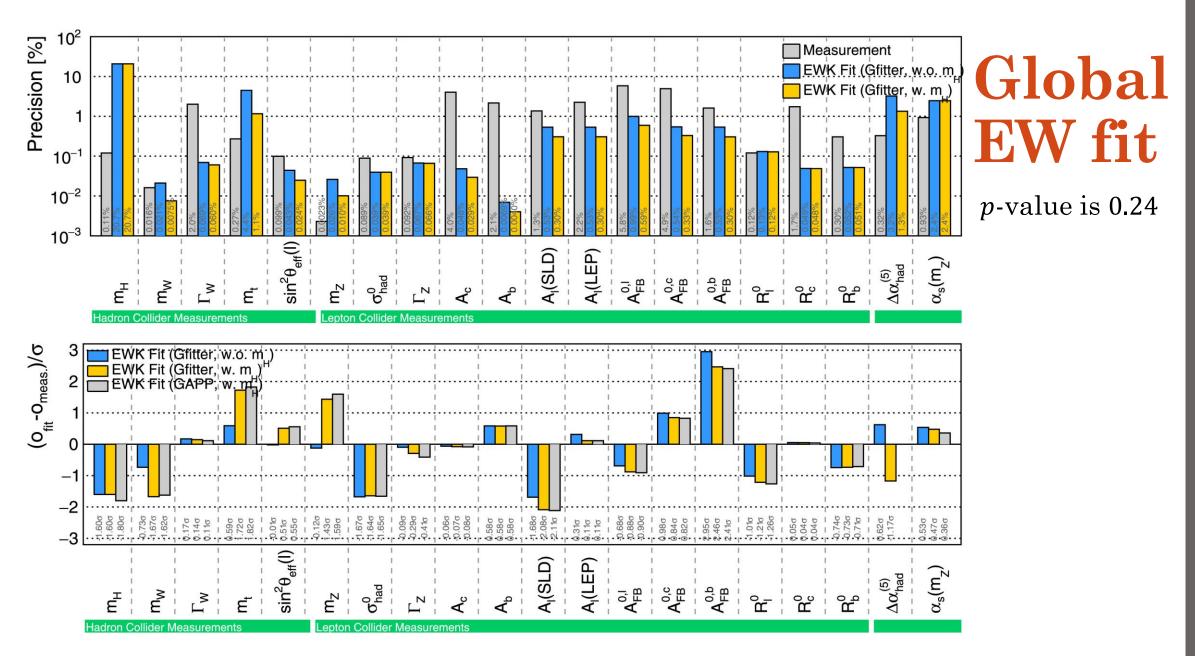
Counting experiment with direct measurement of  $A_f$ 

Table 3.8 Overview of the measured asymmetries at the Z pole from the LEP and SLD experiments [19]. The values are compared to the SM prediction and a pull value for each observable,  $(\mathcal{O}_{\text{measured}} - \mathcal{O}_{\text{predicted}})/\Delta\mathcal{O}$ , is calculated. In addition, the corresponding effective weak mixing angle  $\sin^2\theta_{\text{eff}}^l$  is given. The values indicated with an asterisk have been derived within this work.

Observable	Collider	Value	Total unc.	SM expectation	Pull	Corresponding $\sin^2 \theta_{\text{eff}}^l$
$A_e$	LEP	0.1498	0.0049	$0.1473 \pm 0.0012$	0.5	$0.23117 \pm 0.00062^*$
$A_{ au}$	LEP	0.1439	0.0043	$0.1473 \pm 0.0012$	-0.8	$0.23192 \pm 0.00055$
$A_{\rm FB}^{0,e}$	LEP	0.0145	0.0025	$0.01627 \pm 0.00027$	-0.7	$0.23254 \pm 0.0015^*$
$A_{\mathrm{FB}}^{0,\mu}$	LEP	0.0169	0.0013	$0.01627 \pm 0.00027$	0.5	$0.23113 \pm 0.0007^*$
$A_{\mathrm{FB}}^{0, au}$	LEP	0.0188	0.0017	$0.01627 \pm 0.00027$	1.5	$0.23000 \pm 0.0009^*$
$A_{\rm FB}^{0,l}$	LEP	0.0171	0.001	$0.01627 \pm 0.00027$	0.8	$0.23099 \pm 0.00053$
$A_{\rm FB}^{0,c}$	LEP	0.0699	0.0036	$0.07378 \pm 0.00068$	-1.1	$0.23220 \pm 0.00081$
$A^{0,e}_{{ m FB}} \ A^{0,\mu}_{{ m FB}} \ A^{0, au}_{{ m FB}} \ A^{0, au}_{{ m FB}} \ A^{0,l}_{{ m FB}} \ A^{0,c}_{{ m FB}} \ A^{0,b}_{{ m FB}} \ A^{0,b}_{{ m FB}}$	LEP	0.0992	0.0017	$0.10324 \pm 0.00088$	-2.4	$0.23221 \pm 0.00029$
$A_e$	SLD	0.1516	0.0021	$0.1473 \pm 0.0012$	2.0	$0.23094 \pm 0.00027^*$
$A_{\mu}$	SLD	0.142	0.015	$0.1473 \pm 0.0012$	-0.4	$0.23216 \pm 0.002^*$
$A_{ au}^{'}$	SLD	0.136	0.015	$0.1473 \pm 0.0012$	-0.8	$0.23259 \pm 0.002^*$
$A_l$	SLD	0.1513	0.0021	$0.1473 \pm 0.0012$	1.9	$0.23098 \pm 0.00026$
$A_c$	SLD	0.67	0.027	$0.66798 \pm 0.00055$	0.1	$0.231 \pm 0.008$ *
$A_b$	SLD	0.923	0.02	$0.93462 \pm 0.00018$	-0.6	$0.25 \pm 0.03^*$

**Table 3.9** Overview of selected measurements at LEP, SLD, Tevatron and the LHC of the effective leptonic electroweak mixing angle  $\sin^2 \theta_{\rm eff}^l$  using different observables including a breakdown of different sources of uncertainties. Values which are indicated with an asterisk have not been published and hence only estimated within this work.

$\sin^2  heta_{ ext{eff}}^l$	Value	Stat. unc.	Syst. unc.	PDF unc.	Model unc.	Total unc.	Reference
DØ	0.23095	0.00035	0.00007	0.00019	0.00008	0.00047	[223]
CDF	0.23221	0.00043	0.00003	0.00016	0.00006	0.00046	[224]
Tevatron (combined)	0.23148	0.00027	0.00005	0.00018	0.00006	0.00033	[225]
CMS	0.23101	0.00036	0.00018	0.00030	0.00016	0.00053	[226]
ATLAS (central)	0.23119	0.00031	0.00018	0.00033	0.00006	0.00049	[227]
ATLAS (forward)	0.23166	0.00029	0.00021	0.00022	0.00010	0.00043	[227]
ATLAS (combined)	0.23140	0.00021	0.00014	0.00024	0.00007	0.00036	[227]
LHCb	0.23142	0.00073	0.00052	0.00043*	0.00036*	0.00106	[228]
A <sup>had</sup> (LEP)	0.23240	0.00070	0.00100	_	_	0.00120	[19]
$A_l$ (LEP)	0.23099	0.00042*	0.00032*	_	_	0.00053	[19]
$A_{\tau} + A_{e}$ (LEP)	0.23159	0.00037*	0.00018*	_	_	0.00041	[19]
$A_{\rm FB}^b$ (LEP)	0.23221	0.00023*	0.00017*	_	_	0.00029	[19]
$A_l$ (SLD)	0.23098	0.00024	0.00013	_	_	0.00026	[19]



https://doi.org/10.1016/j.ppnp.2019.02.007

#### Polarized electron source

190 - 560 hours

•	Beam polarization	80 %
•	Cathode voltage (pulsed)	-100 kV
•	Photocathode type	Strained InGaAsP
•	Laser type	Ti-Sapphire
•	Light wavelength	700 - 850  nm
•	Laser power in a pulse	200 W
•	Pulse duration	2.1 μs
•	Repetition rate	1 Hz
•	Maximum current from a gun	150 mA
•	Operational current	$15-20~\mathrm{mA}$
•	Photocathode recesiation time	

(depends on laser power)

