

Rare charm decays to invisibles



based on works with Rigo Bause, Marcel Golz and Hector Gisbert,
2007.05001 [hep-ph], 2010.02225 [hep-ph]

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... because, we can:

- sizable SM branching ratios within $10^{-7} - 10^{-6}$ (semileptonic $c \rightarrow ull$), and $10^{-6} - 10^{-4}$ (radiative $c \rightarrow u\gamma$)
- plenty of BSM opportunities (this talk and talk by Marcel Golz)
- in fact, its already happening LHCb'17,18,21 $D \rightarrow \pi\pi\mu\mu$, Belle'16 $D \rightarrow \rho\gamma$, BES III '18 $D \rightarrow \pi\pi ee$

AND its a **unique probe of up-sector physics**:

- 1) leaving no stone unturned (BSM searches)
- 2) complementarity (w.r.t. K,B) (flavor origins)

pursue general flavor physics , $t, b, c, s, \dots$, exploit correlations

BSM opportunities with $|\Delta c| = |\Delta u| = 1$ studies

In view of the hadronic backgrounds in rare charm decays, the name of the game in flavor/BSM probes is "null tests", based on (approximate) symmetries of the SM, or optimized observables with reduced SM uncertainties:

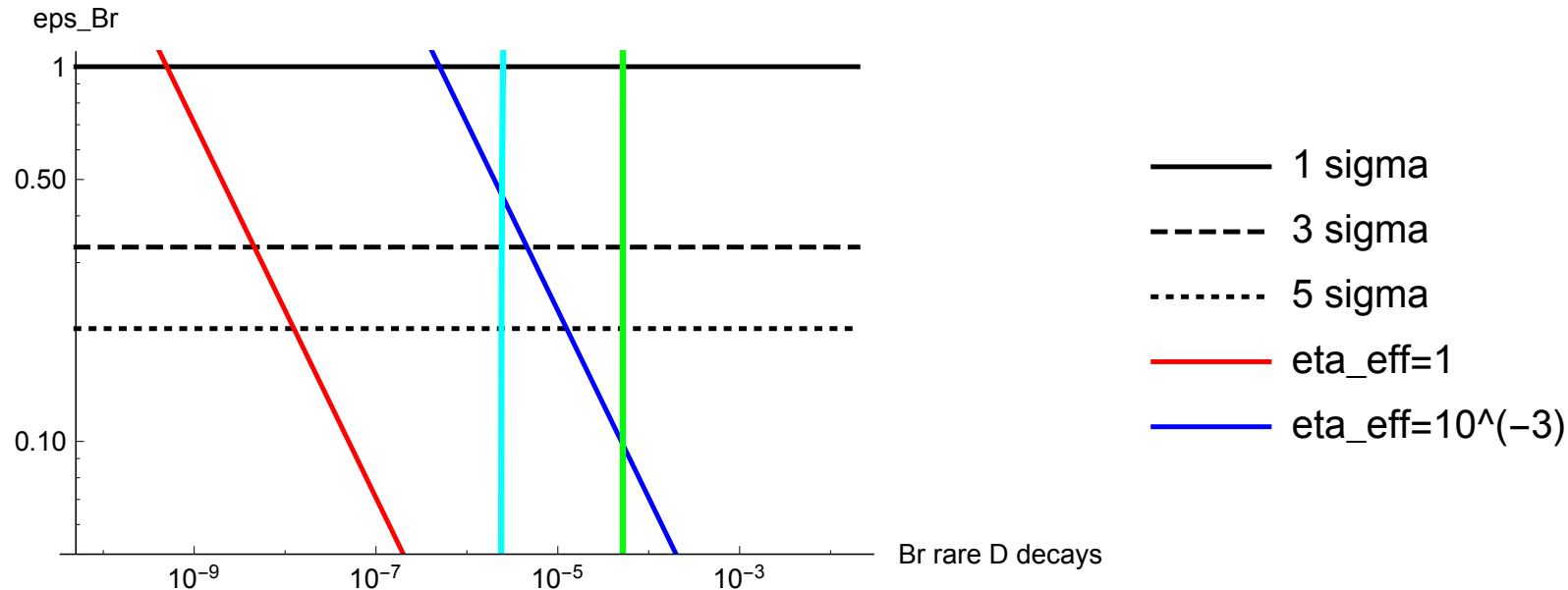
lepton-universality, lepton flavor conservation, CP, polarization studies, data-driven SM estimations, angular distributions

Procedure very well-known in state-of-the-art b -physics studies; essential in pursuit of beauty-anomalies, e.g., R_K , P'_5 .

Charm FCNCs are in addition subject to a superb **GIM-suppression**, implying that $c \rightarrow u\nu\bar{\nu}$ modes *this talk* and contributions to $c \rightarrow u\ell^+\ell^-$ with axial-vector lepton-currents " $C_{10}^{(\prime)}$ " *see talk by Marcel Golz* vanish in SM.

Testing the SM AND lepton universality with $c \rightarrow u\nu\bar{\nu}$

What can you do with $2 \cdot 10^9$ D -mesons:



Sensitivity $\epsilon(Br) \simeq 1/\sqrt{N(h_c)Br(h_c \rightarrow F\nu\bar{\nu})\eta_{eff}}$ for reconstruction efficiency η_{eff}

limits from data + theory/SMEFT [2010.02225](#); $Br(D^+ \rightarrow \pi^+ \nu\bar{\nu})^{\text{SM}} \simeq 0$

Green: $Br(D^+ \rightarrow \pi^+ \nu\bar{\nu})^{\text{EFTmax}} \simeq 5.2 \cdot 10^{-5}$

Cyan: $Br(D^+ \rightarrow \pi^+ \nu\bar{\nu})^{\text{LUmax}} \simeq 2.5 \cdot 10^{-6}$ assuming lepton universality

radiative decays, example $D^0 \rightarrow \rho^0 \gamma$, $Br = (1.77 \pm 0.31) \cdot 10^{-5}$. Belle '16: $943 fb^{-1}$: 500 ± 85 events measures the direct CP-asymmetry with $\delta A_{CP} = 16\%$;

semileptonic decays, example $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$, $Br \sim 9.6 \cdot 10^{-7}$. LHCb '18: $5 fb^{-1}$: 1083 ± 41 events measures angular observables with $\delta A_{CP} = 3.7\%$;

missing energy – only single search in PDG: $D^0 \rightarrow$ nothing, $Br < 9.4 \cdot 10^{-5}$ @ 90% CL UL. Belle'16: $924 fb^{-1}$:

LHCb excellent at $c \rightarrow u \mu^+ \mu^-$ precision studies.

missing energy, π^0 , photon, τ and $e(?)$ final states advantageous w.r.t hadronic collider territory

Upper limits on $c \rightarrow u\nu\bar{\nu}$ from data+theory

$h_c \rightarrow F\nu\bar{\nu}$	$\mathcal{B}_{\text{LU}}^{\text{max}}$ [10 ⁻⁷]	$\mathcal{B}_{\text{cLFC}}^{\text{max}}$ [10 ⁻⁶]	$\mathcal{B}^{\text{max}}$ [10 ⁻⁶]	$N_{\text{LU}}^{\text{max}}/\eta_{\text{eff}}$	$N_{\text{cLFC}}^{\text{max}}/\eta_{\text{eff}}$	$N^{\text{max}}/\eta_{\text{eff}}$
$D^0 \rightarrow \pi^0$	6.1	3.5	13	47 k (395 k)	270 k (2.3 M)	980 k (8.3 M)
$D^+ \rightarrow \pi^+$	25	14	52	77 k (650 k)	440 k (3.7 M)	1.6 M (14 M)
$D_s^+ \rightarrow K^+$	4.6	2.6	9.6	6 k (50 k)	34 k (290 k)	120 k (1.1 M)
$D^0 \rightarrow \pi^0\pi^0$	1.5	0.8	3.1	11 k (95 k)	64 k (540 k)	230 k (2.0 M)
$D^0 \rightarrow \pi^+\pi^-$	2.8	1.6	5.9	22 k (180 k)	120 k (1.0 M)	450 k (3.8 M)
$D^0 \rightarrow K^+K^-$	0.03	0.02	0.06	0.2 k (1.9 k)	1.3 k (11 k)	4.8 k (40 k)
$\Lambda_c^+ \rightarrow p^+$	18	11	39	14 k (120 k)	82 k (700 k)	300 k (2.6 M)
$\Xi_c^+ \rightarrow \Sigma^+$	36	21	76	28 k (240 k)	160 k (1.4 M)	590 k (5.0 M)
$D^0 \rightarrow X$	15	8.7	32	120 k (980 k)	660 k (5.6 M)	2.4 M (21 M)
$D^+ \rightarrow X$	38	22	80	120 k (1.0 M)	680 k (5.8 M)	2.5 M (21 M)
$D_s^+ \rightarrow X$	18	10	38	24 k (200 k)	140 k (1.1 M)	500 k (4.2 M)

Table 1: Upper limits on branching ratios and expected number of events $N_F^{\text{exp}} = \eta_{\text{eff}} N(h_c) \mathcal{B}(h_c \rightarrow F\nu\bar{\nu})$ per reconstruction efficiency η_{eff} for Belle II with 50 ab⁻¹ (FCC-ee in parentheses) corresponding to LU, cLFC, and general. [from 2010.02225](#)

Testing lepton universality and charged lepton flavor conservation with dineutrino* modes

method complementary to lepton-flavor specific tests a la R_K

concrete model-independent proposal 2007.05001 [hep-ph]

application to beauty 2109.01675 [hep-ph]

application to charm 2010.02225 [hep-ph]

Idea: Use $SU(2)_L$ -link for left-handed SM leptons $L = (\nu, \ell)$; assume only SM-like light neutrinos (crosscheck possible)

*neutrino flavors untagged!

SMEFT: SM plus higher dimensional operators at scale $\Lambda \gg \Lambda_{EWK}$, consistent with Lorentz and SM-gauge symmetries

$SU(3)_C \times SU(2)_L \times U(1)_Y$ made out of SM fields.

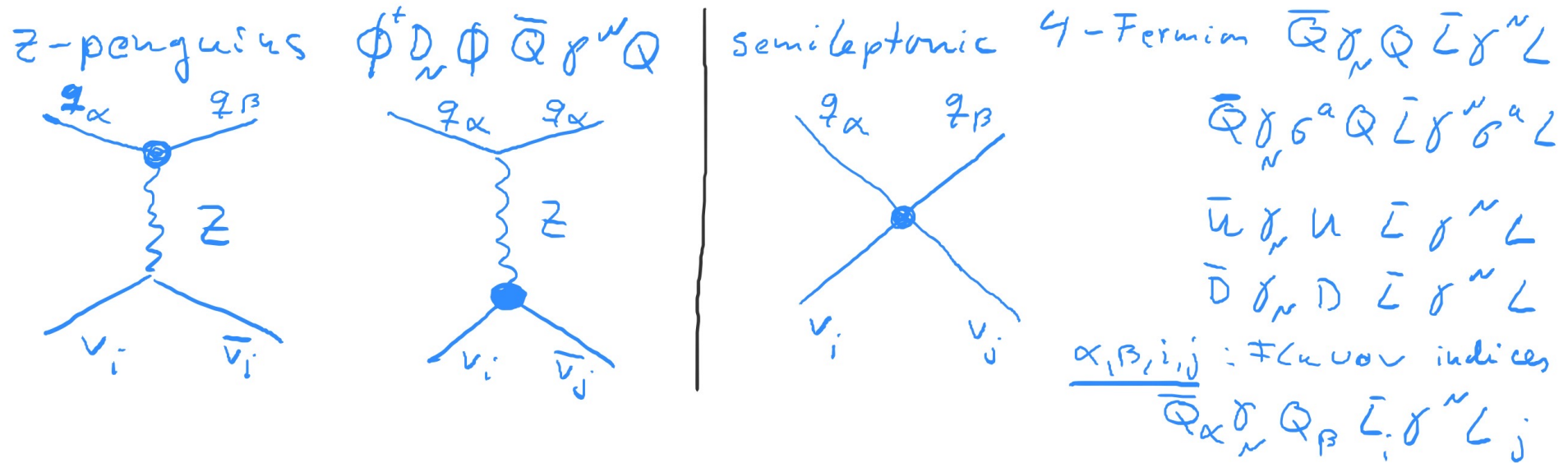
Leading terms: dim 6, Warsaw basis Misiak et al

For dineutrino modes $q \rightarrow q' \nu \bar{\nu}$ two types of contributions at tree-level:

leptonic and quark Z -penguins $\Phi^\dagger D_\mu \Phi \bar{F} \gamma^\mu F$, $F = L, Q, U, D$

semileptonic 4-fermion operators of type $\bar{Q} \gamma_\mu Q \bar{L} \gamma^\mu L$

moving ahead model-independently



For $SU(2)_L$ -doublet fermions, singlet and triplet $\sim \tau^3$ contributions

4-fermion operators: 2 lepton flavor indices i, j and 2 quark flavor indices α, β

Z -penguins strongly constrained Falkowski et al

Leading semileptonic 4-fermion operators at scale above m_W (SMEFT) contributing to dineutrino modes $q \rightarrow q' \nu \bar{\nu}$

$$\mathcal{L}_{\text{eff}} \supset \frac{C_{\ell q}^{(1)}}{v^2} \bar{Q} \gamma_\mu \textcolor{red}{Q} \bar{L} \gamma^\mu \textcolor{blue}{L} + \frac{C_{\ell q}^{(3)}}{v^2} \bar{Q} \gamma_\mu \tau^a Q \bar{L} \gamma^\mu \tau^a L + \frac{C_{\ell u}}{v^2} \bar{U} \gamma_\mu U \bar{L} \gamma^\mu L + \frac{C_{\ell d}}{v^2} \bar{D} \gamma_\mu D \bar{L} \gamma^\mu L. \quad (1)$$

$SU(2)_L \times U(1)_Y$ gauge invariance links up and down quarks, $\textcolor{red}{Q} = (u, d)$ and left-handed neutrinos and charged leptons $\textcolor{blue}{L} = (\nu, \ell)$.

$$\textcolor{blue}{C}_L^{\textcolor{red}{U}} = \textcolor{blue}{K}_L^{\textcolor{red}{D}} = C_{\ell q}^{(1)} + C_{\ell q}^{(3)}, \quad C_R^U = K_R^U = C_{\ell u}, \quad C_L^D = K_L^U = C_{\ell q}^{(1)} - C_{\ell q}^{(3)}, \quad C_R^D = K_R^D = C_{\ell d}.$$

contribution to $c \rightarrow u \nu \bar{\nu}$ (C_L^U) identical to $s \rightarrow d \ell \bar{\ell}$ (K_L^D) etc

L,R denotes left or right handed quark currents; only SM-like light neutrinos.

Since the neutrino flavors are not tagged, the branching ratio, say for charm, is obtained by an incoherent sum

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) = \sum_{i,j} \mathcal{B}(c \rightarrow u \nu_i \bar{\nu}_j)$$

in terms of Wilson coefficients

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) = \sum_{i,j} \mathcal{B}(c \rightarrow u \nu_i \bar{\nu}_j) \propto \sum_{i,j} |\mathcal{C}_L^{Uij}|^2 + |\mathcal{C}_R^{Uij}|^2$$

which can be written as a trace

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) \propto \sum_{i,j} |\mathcal{C}_L^{Uij}|^2 + |\mathcal{C}_R^{Uij}|^2 = \text{tr} \left[\mathcal{C}_L^U \mathcal{C}_L^{U\dagger} + \mathcal{C}_R^U \mathcal{C}_R^{U\dagger} \right]$$

and apparently PMNS rotations W between gauge and mass eigenstates (calligraphic) drop out due to unitarity.

Using SMEFT $SU(2)_L \times U(1)_Y$ gauge invariance follows

$$\mathcal{C}_L^U = W^\dagger \mathcal{K}_L^D W + \mathcal{O}(\lambda), \quad \mathcal{C}_R^U = W^\dagger \mathcal{K}_R^U W$$

$$\begin{aligned}
 \underbrace{\mathcal{B}(c \rightarrow u \nu \bar{\nu})}_{\text{dineutrino process}} &\propto \sum_{\nu=i,j} \left(|\mathcal{C}_L^{Uij}|^2 + |\mathcal{C}_R^{Uij}|^2 \right) = \text{tr} \left[\mathcal{C}_L^U \mathcal{C}_L^{U\dagger} + \mathcal{C}_R^U \mathcal{C}_R^{U\dagger} \right] \\
 &= \text{tr} \left[\mathcal{K}_L^D \mathcal{K}_L^{D\dagger} + \mathcal{K}_R^U \mathcal{K}_R^{U\dagger} \right] + \mathcal{O}(\lambda) = \underbrace{\sum_{\ell=i,j} \left(|\mathcal{K}_L^{Dij}|^2 + |\mathcal{K}_R^{Uij}|^2 \right)}_{\text{charged dilepton couplings}} + \mathcal{O}(\lambda) ,
 \end{aligned}$$

$\mathcal{K}_L^{Dij}$: coeffs for mass eigenstate charged leptons in $s \rightarrow d \ell^{i+} \ell^{j-}$

$\mathcal{K}_R^{Uij}$: coeffs mass eigenstates charged leptons in $c \rightarrow u \ell^{i+} \ell^{j-}$

lhs: observable dineutrino branching ratio

rhs: couplings to charged leptons/Wilson coefficients

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) \propto \sum_{\ell=i,j} \left(|\mathcal{K}_L^{Dij}|^2 + |\mathcal{K}_R^{Uij}|^2 \right) + \mathcal{O}(\lambda) ,$$

We obtain upper limits on dineutrino branching ratios from upper limits of charged dilepton modes [2007.05001](#) depending on scenarios

- i) $\mathcal{K}_{L,R}^{ij} \propto \delta_{ij}$, that is, *lepton-universality* (LU).
- ii) $\mathcal{K}_{L,R}^{ij}$ are diagonal, that is, *charged lepton flavor conservation* (cLFC).
- iii) $\mathcal{K}_{L,R}^{ij}$ arbitrary (cLFV).

universality tests with $q \rightarrow q' \nu \bar{\nu}$, quantitatively

		ee	$\mu\mu$	$\tau\tau$	$e\mu$	$e\tau$	$\mu\tau$
sd	$ \bar{\mathcal{K}}_{L,R}^{D\ell\ell'} $	3.5	1.9	6.7	2.0	6.1	6.6
cu	$ \bar{\mathcal{K}}_{L,R}^{U\ell\ell'} $	2.9	1.6	5.6	1.6	4.7	5.1
bd	$ \bar{\mathcal{K}}_{L,R}^{D_{13}\ell\ell'} $	4.9	2.7	9.5	3.1	9.7	10
bs	$ \bar{\mathcal{K}}_{L,R}^{D_{23}\ell\ell'} $	13	7.0	25	8.0	27	30

Upper limits on leptonic couplings $\bar{\mathcal{K}}_{L,R}$ from high- p_T [Fuentes-Martin et al 2020](#), [Angelescu et al 2020](#). The last two rows show limits on down sector couplings involving b -quarks. LFV-bounds are quoted as charge-averaged, $\sqrt{|\bar{\mathcal{K}}^{\ell^+\ell'^-}|^2 + |\bar{\mathcal{K}}^{\ell^-\ell'^+}|^2}$.

universality tests with $c \rightarrow u \nu \bar{\nu}$, quantitatively

including CKM-corrections $\mathcal{C}_L^U = W^\dagger \mathcal{K}_L^D W + \mathcal{O}(\lambda)$, $\lambda \simeq V_{us} \sim 0.2$

$$R^{\ell\ell'} = |\mathcal{K}_L^{D_{12}\ell\ell'}|^2 + |\mathcal{K}_R^{U_{12}\ell\ell'}|^2, \quad R_{\pm}^{\ell\ell'} = |\mathcal{K}_L^{D_{12}\ell\ell'} \pm \mathcal{K}_R^{U_{12}\ell\ell'}|^2, \quad R_{\pm}^{\ell\ell'} \leq 2 R^{\ell\ell'}$$
$$\delta R^{\ell\ell'} = 2 \lambda \Re(\mathcal{K}_L^{D_{12}\ell\ell'} \mathcal{K}_L^{D_{22}\ell\ell'}{}^* - \mathcal{K}_L^{D_{12}\ell\ell'} \mathcal{K}_L^{D_{11}\ell\ell'}{}^*) < 2 \lambda |\mathcal{K}_L^{D_{12}\ell\ell'}| (|\mathcal{K}_L^{D_{22}\ell\ell'}| + |\mathcal{K}_L^{D_{11}\ell\ell'}|)$$

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) \propto 3 R^{\mu\mu} \lesssim 34, \quad (\text{LU}) \quad (2)$$

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) \propto R^{ee} + R^{\mu\mu} + R^{\tau\tau} \lesssim 196, \quad (\text{cLFC}) \quad (3)$$

$$\mathcal{B}(c \rightarrow u \nu \bar{\nu}) \propto R^{ee} + R^{\mu\mu} + R^{\tau\tau} + 2(R^{e\mu} + R^{e\tau} + R^{\mu\tau}) \lesssim 716. \quad (4)$$

dimuons have the most stringent bounds and provide the LU-limit.

Lepton flavor limit is violated if measured branching ratio is too large!

Upper limits data-driven, and evolve with charged lepton data.

$\mathcal{B}(D^0 \rightarrow \text{inv.}) < 9.4 \cdot 10^{-5}$, at 90 % CL. (Belle '16). Consistency check; constrains operators with light right-handed neutrinos

$$Q_{LR}^{ij} = (\bar{u}_L \gamma_\mu c_L) (\bar{\nu}_{jR} \gamma^\mu \nu_{iR}), \quad Q_{RR}^{ij} = (\bar{u}_R \gamma_\mu c_R) (\bar{\nu}_{jR} \gamma^\mu \nu_{iR}),$$
$$Q_{S(P)}^{ij} = (\bar{u}_L c_R) (\bar{\nu}_j (\gamma_5) \nu_i), \quad Q_{T(T5)}^{ij} = \frac{1}{2} (\bar{u} \sigma_{\mu\nu} c) (\bar{\nu}_j \sigma^{\mu\nu} (\gamma_5) \nu_i),$$

$Q_{S(P)}^{ij}$ would have effect less than $\sim 10\%$ of LU upper limits iff improved limit exists

$$\mathcal{B}(D^0 \rightarrow \text{inv.}) \lesssim 2 \cdot 10^{-6}. \quad (5)$$

would reinforce framework.

Bounding Lepton Number Violation

The final state $D \rightarrow \textit{nothing}$ could be two neutrinos, $D \rightarrow \nu\nu$, allowing to probe LNV in $\Delta L = 2$ transitions.

Lowest order operator $\mathcal{O}_{4a}^{(7)} = L_i^\alpha L_j^\beta \bar{Q}_\alpha^b \bar{U}_a^c H^\rho \epsilon_{\beta\rho}$, de Gouvea

Existing Belle limit on $D \rightarrow \textit{nothing}$ probes LNV effects

$$\Lambda_{\text{LNV}}^{ij} \gtrsim 1.5 \text{ TeV} .$$

New ideas for LU and cLFC tests with dineutrino modes suitable for Belle II, BES III, Z -factory and STCF.

Great opportunity for rare charm decays:

- No exp search on any semileptonic $c \rightarrow u\nu\bar{\nu}$ mode to date!
- many modes: $D \rightarrow \pi$, $D \rightarrow \pi\pi$, $\Lambda_c \rightarrow p$, $\Xi_c \rightarrow \Sigma$... all SM nulltests
- model-independent upper limits sizable $10^{-6} - 10^{-5}$
If in excess of $\sim 1/3$: charged lepton flavor is violated
If in excess of $\sim 1/20$: lepton universality is broken
- Improved $D^0 \rightarrow \text{nothing}$ probes light right-handed neutrinos;
Lepton number violation also testable, presently $\Lambda_{\text{LNV}}^{ij} \gtrsim 1.5 \text{ TeV}$.

BACK UP

h_c	$f(c \rightarrow h_c)$	$N(h_c)$ (a)	$N(h_c)$ (b)
D^0	0.59	$6 \cdot 10^{11}$	$8 \cdot 10^{10}$
D^+	0.24	$3 \cdot 10^{11}$	$3 \cdot 10^{10}$
D_s^+	0.10	$1 \cdot 10^{11}$	$1 \cdot 10^{10}$
Λ_c^+	0.06	$7 \cdot 10^{10}$	$8 \cdot 10^9$

Table 2: Charm fragmentation fractions $f(c \rightarrow h_c)$ and the number of charmed hadrons h_c , $N(h_c)$, expected at benchmarks with $N(c\bar{c}) = 550 \cdot 10^9$ (a, FCC-ee) and $N(c\bar{c}) = 65 \cdot 10^9$ (b, Belle II with 50 ab^{-1}). In absence of further information for the Ξ_c^+ we use $f(c \rightarrow \Xi_c^+) \simeq f(c \rightarrow \Lambda_c^+)$.

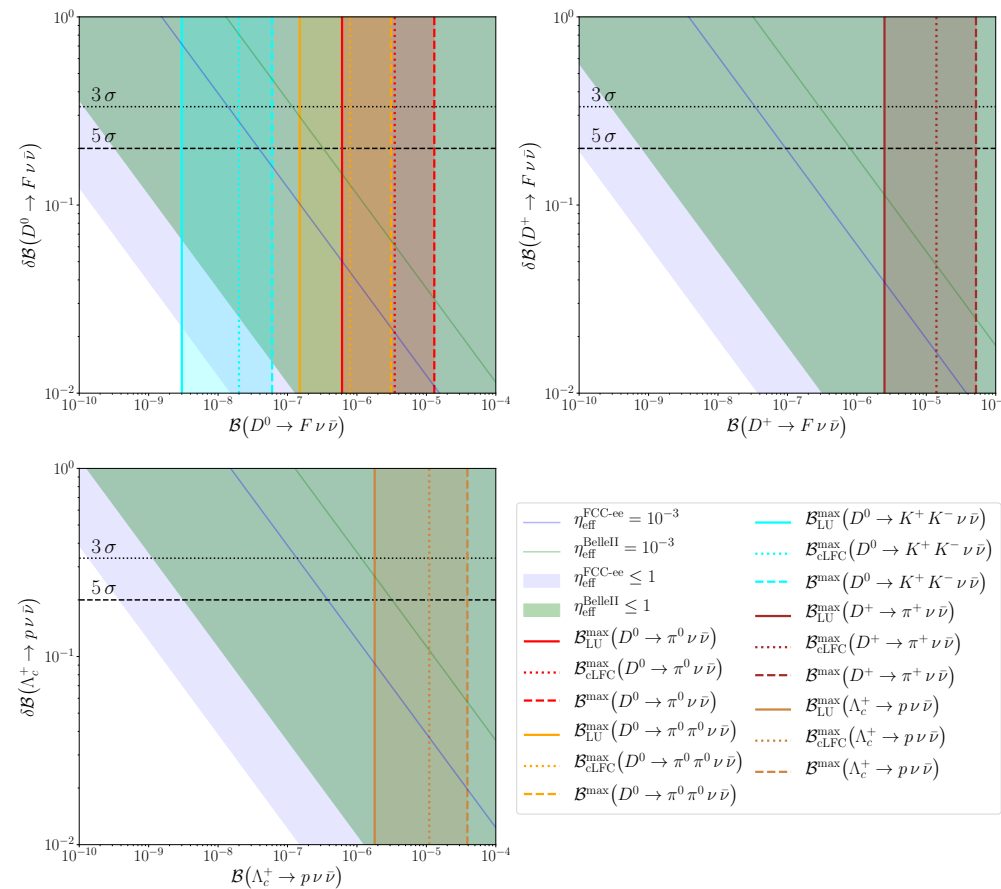


Figure 1: Relative statistical uncertainty $\delta\mathcal{B}$ versus branching ratio $\mathcal{B}$ for decays of the D^0 (upper left), the D^+ (upper right) and the Λ_c^+ (lower plot). Shaded areas correspond to $\eta_{\text{eff}} = 1$, whereas the solid tilted lines illustrate the impact of reconstruction efficiencies $\eta_{\text{eff}} = 10^{-3}$ for the FCC-ee (lilac) and Belle II (green). Horizontal lines indicate 3σ (dotted) and 5σ (dashed) black. Vertical lines represent upper limits assuming LU (solid), cLFC (dotted) and generic lepton flavor (dashed) for different decays.